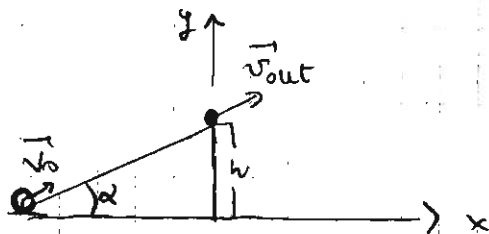


Pb. 1

1) Dalle conservazione dell'energia meccanica

$$\frac{1}{2} m v_0^2 = \frac{1}{2} m v_{out}^2 + mgh \Rightarrow v_{out} = \sqrt{v_0^2 - 2gh}$$

2) Ad $h_{max} \Rightarrow v_y = 0 \Rightarrow \vec{v} = v_x \hat{x}$. Lungo x , il moto è rettilineo uniforme con

$$v_x = v_{out} \cos \alpha = v_{out} \cos(\pi/6)$$

$$3) \quad x(t) = v_x t$$

$$y(t) = h + v_y t - \frac{1}{2} g t^2$$

$$v_x = v_{out} \cos(\pi/6)$$

$$v_y = v_{out} \sin(\pi/6)$$

$$y(t) = 0 \Rightarrow h + v_y t - \frac{1}{2} g t^2 = 0$$

$$g t^2 - 2 v_y t - 2h = 0$$

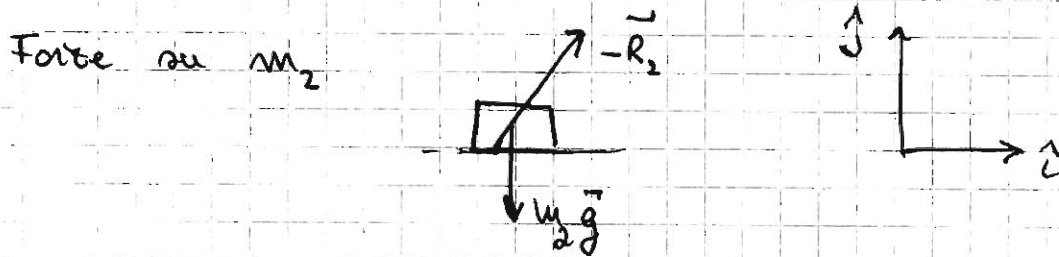
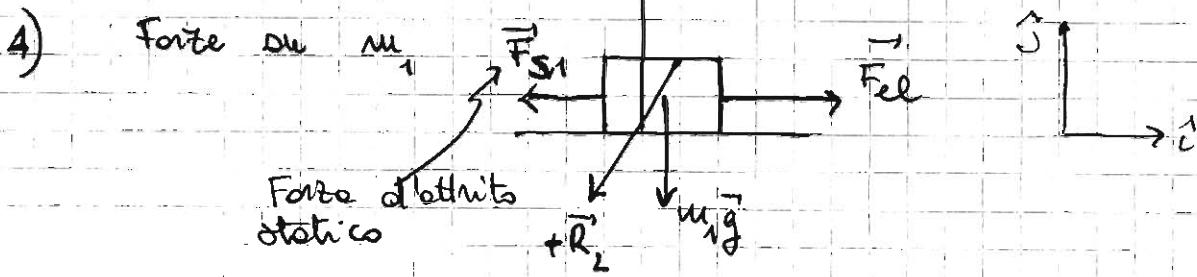
$$t = \frac{v_y \pm \sqrt{v_y^2 + 2hg}}{g}$$

La soluzione accettata è quella positiva $\Rightarrow$

$$d_x = v_x \frac{v_y}{g} \left(1 + \sqrt{1 + \frac{2gh}{v_y^2}} \right)$$

$$= \frac{v_{out}^2}{g} \cos(\pi/6) \sin(\pi/6) \left(1 + \sqrt{1 + \frac{2gh}{v_{out}^2 \sin^2(\pi/6)}} \right)$$

Pb. 2



Equilibrio $\Rightarrow$

$$\begin{cases} 0 = m_1 \vec{g} + \vec{F}_{el} + \vec{N} + \vec{R}_2 + \vec{F}_{s,1} \\ 0 = m_2 \vec{g} - \vec{R}_2 \end{cases}$$

$$\Rightarrow \begin{cases} 0 = -m_1 g + N - R_{2y} \\ 0 = -F_{s,1} + F_{el} - R_{2x} \\ 0 = -m_2 g + R_{2y} \\ 0 = +R_{2x} \end{cases} \Rightarrow \begin{cases} N = m_1 g + m_2 g \\ F_{el} = F_{s,1} \\ \left. \begin{aligned} R_{2y} &= m_2 g \\ R_{2x} &= 0 \end{aligned} \right\} \uparrow \end{cases}$$

Poiché $F_{s,1} \leq \mu_s N = \mu_s (m_1 + m_2) g$

$$F_{el} \leq \mu_s (m_1 + m_2) g$$

||

$$k d \leq \mu_s (m_1 + m_2) g \Rightarrow \mu_s \geq \frac{k d}{(m_1 + m_2) g}$$

5) Con lo stesso diagramma delle forze

6)
$$\begin{cases} m_1 \vec{a}_1 = m_1 \vec{g} + \vec{F}_{el} + \vec{N} + \vec{R}_2 + \vec{F}_{d,1} \\ m_2 \vec{a}_2 = m_2 \vec{g} - \vec{R}_2 \end{cases}$$
 ← forze d'attrito dinamico

$$\begin{cases} 0 = -m_1 g + N - R_{2y} \\ m_1 a_1 = F_{el} - R_{2x} - \mu_d N \\ 0 = -m_2 g + R_{2y} \\ m_2 a_2 = +R_{2x} \end{cases}$$

forze d'attrito dinamico su 2

$$R_{2x} \equiv \vec{F}_{d,2} = \mu_d R_{2y}$$

$$R_{2y} = m_2 g \quad m_2 a_2 = +\mu_d R_{2y} = +\mu_d m_2 g$$

$$\Rightarrow a_2 = \mu_d g$$

Per 1:

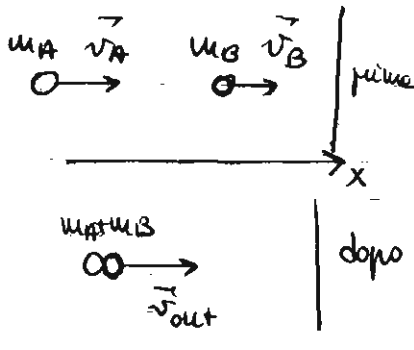
$$\begin{cases} 0 = -m_1 g + N - m_2 g \\ m_1 a_1 = kd - \mu_d m_2 g - \mu_d N \end{cases}$$

$$\begin{cases} N = (m_1 + m_2) g \\ m_1 a_1 = kd - \mu_d (2m_2 + m_1) g \end{cases}$$

$$\Rightarrow a_1 = \frac{kd - \mu_d (2m_2 + m_1) g}{m_1}$$

7) $m_A \vec{v}_A + m_B \vec{v}_B = (m_A + m_B) \vec{v}_{out}$

$\Rightarrow \vec{v}_{out} = \frac{m_A \vec{v}_A + m_B \vec{v}_B}{m_A + m_B}$

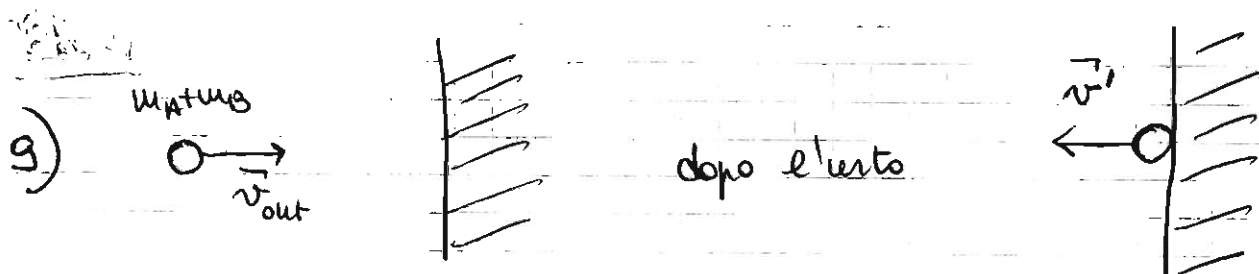


8) $E_{diss} = E_{cin}(\text{prima dell'urto}) - E_{cin}(\text{dopo l'urto})$

$$= \frac{1}{2} m_A v_A^2 + \frac{1}{2} m_B v_B^2 - \frac{1}{2} (m_A + m_B) \frac{(m_A v_A + m_B v_B)^2}{(m_A + m_B)^2}$$

$$= \frac{1}{2} \frac{1}{m_A + m_B} \left[\cancel{m_A^2 v_A^2} + m_A m_B v_A^2 + m_A m_B v_B^2 + \cancel{m_B^2 v_B^2} - \cancel{m_A^2 v_A^2} - \cancel{m_B^2 v_B^2} - 2 m_A m_B v_A v_B \right]$$

$$= \frac{1}{2} \frac{m_A m_B}{m_A + m_B} [v_A - v_B]^2$$



Troviamo v' :

$$\frac{1}{2} (m_A + m_B) v_{out}^2 = \frac{1}{2} (m_A + m_B) v'^2$$

↑
urto elastico

$\Rightarrow v' = v_{out}$

Per la scelta delle direzioni e verso, si utilizza il teorema dell'impulso. L'unica forza impulsiva che agisce su $m_A + m_B$ è $\vec{R}$ (reazione del muro, $\perp$ al muro, perché è urto elastico $\Rightarrow$ non ci sono forze di attrito dissipative)

$$\Rightarrow \vec{I}_{\vec{R}} = \int_0^{\Delta t = \text{tempo dell'urto}} \vec{R} dt \quad \text{dove } \text{essen} \perp \text{ al muro}$$

$$\Rightarrow \text{perch\`e} \quad \vec{I} = \vec{Q}_{\text{fin}} - \vec{Q}_{\text{in}}$$

↑
quantit\`e di moto delle palline A+B

$$I_{\parallel} = 0 = m_{AB} v'_{\parallel} - 0 \Rightarrow v'_{\parallel} = 0 \Rightarrow \vec{v}' \perp \text{ muro}$$

mostra

$$I_{\perp} = I_R = (m_A + m_B) v' - (m_A + m_B)(-v')$$

$$= 2(m_A + m_B) v' = 2(m_A + m_B) v_{\text{out}}$$

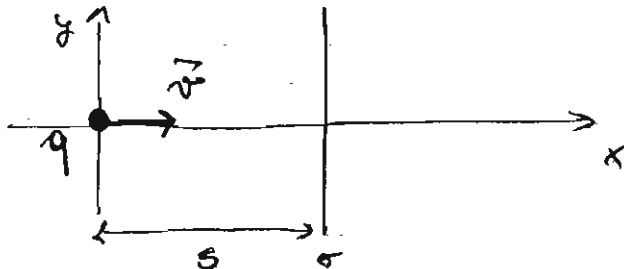
Pb. 4

10) Il campo elettrico di un piano infinito con densit\`a sup.  $+\sigma$  \`e uniforme, $\perp$ al piano (vd. figura) e di modulo $E = \frac{\sigma}{2\epsilon_0}$



(lo si pu\`o eventualmente dimostrare facilmente col Teorema di Gauss)

$\Rightarrow$



$$\vec{F} = q\vec{E} = -q \frac{\sigma}{2\epsilon_0} \hat{x} \Rightarrow \vec{a} = \frac{\vec{F}}{m} = -\frac{q}{m} \frac{\sigma}{2\epsilon_0} \hat{x}$$

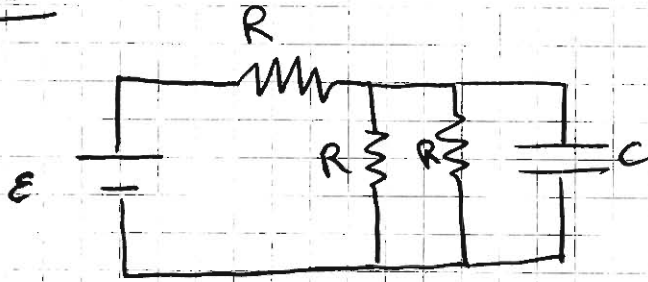
$$11) \quad v_x = v_0 - at \quad v_x = 0 \Rightarrow t = \frac{v_0}{a}$$

$$x(t) = v_0 t - \frac{1}{2} at^2$$

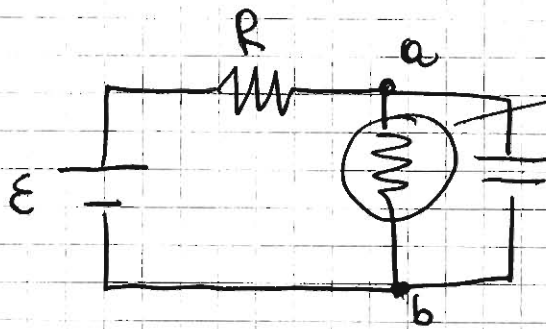
$$\Rightarrow x = \frac{v_0^2}{2a} \Rightarrow d = s - x = s - \frac{v_0^2}{2a} = s - \frac{v_0^2 m \epsilon_0}{\sigma q}$$

Pb. 5

12)

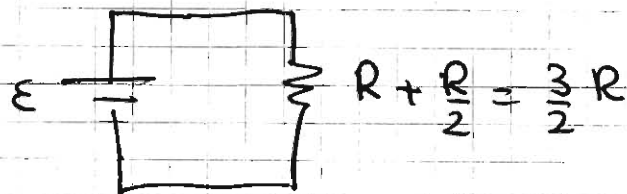


↓ equivalente a



$$\frac{1}{R_{eq}} = \frac{1}{R} + \frac{1}{R} \Rightarrow R_{eq} = \frac{R}{2}$$

In condizioni stazionarie, il cond. è carico e quindi con ci passa corrente $\Rightarrow$ i passano solo dalle resistenze $\Rightarrow$



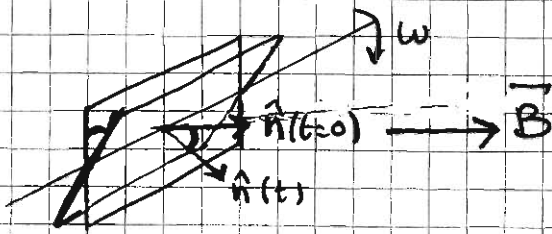
$$\Rightarrow \frac{3}{2} R i = E \quad i = \frac{2}{3} \frac{E}{R}$$

$$13) \quad \Delta V_{ab} = i \frac{R}{2} = \frac{2}{3} \frac{E}{R} \frac{R}{2} = \frac{E}{3}$$

Poichè $\frac{Q}{\Delta V_{ab}} = C \Rightarrow Q = C \cdot \frac{E}{3}$

Pb. 6

14)



$$\phi_B(t) = B(a \cdot b) \cos(\omega t)$$

$$\Rightarrow \frac{d\phi_B}{dt} = -Bab \omega \sin \omega t \quad \Rightarrow$$

$$\mathcal{E} = Bab \omega \sin \omega t$$

$$15) \quad P_{\max} = \frac{\mathcal{E}_{\max}^2}{R} = \frac{(Bab\omega)^2}{R}$$