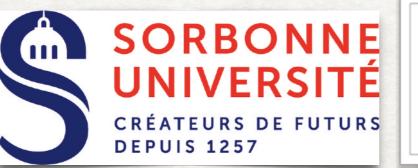
#### ROBERTO MANUZZO, PHD STUDENT

#### "MAGNETOPAUSE STUDY WITH A MULTI-FLUID NUMERICAL CODE SIMULATION"



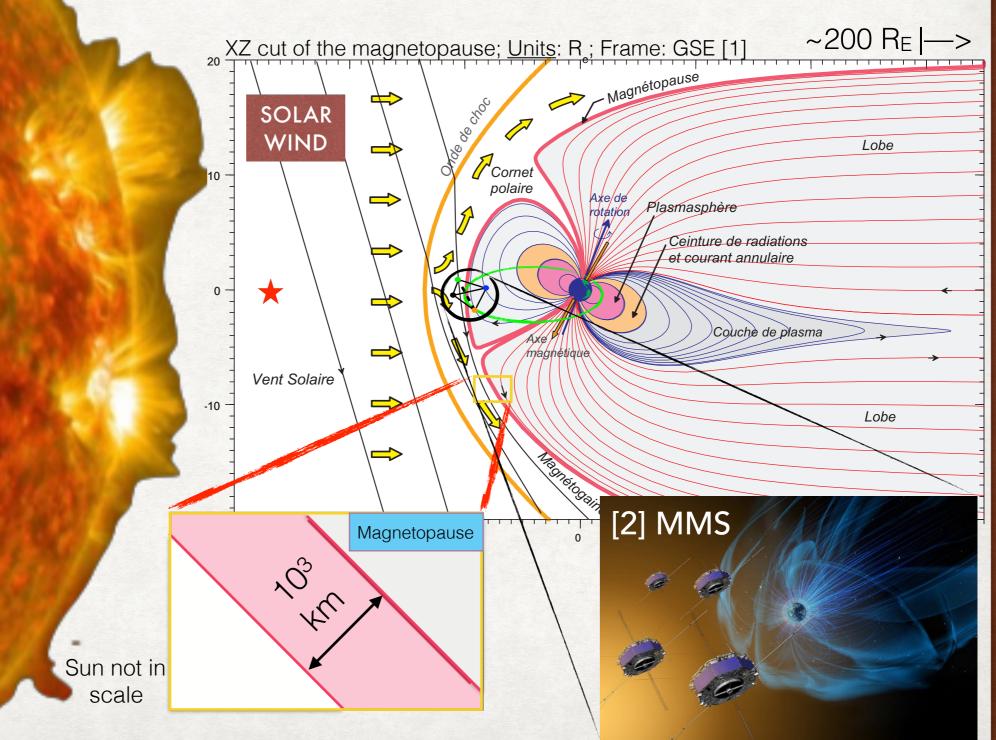




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# "MAGNETOPAUSE STUDY WITH A MULTI-FLUID NUMERICAL CODE SIMULATION"



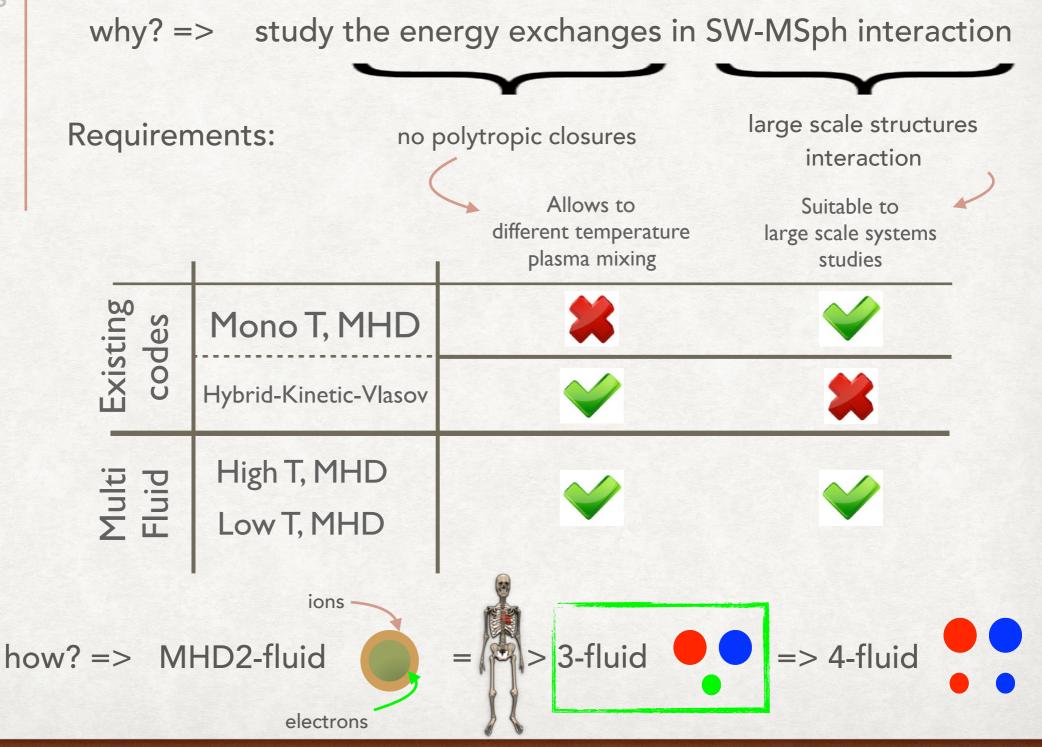
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### "MAGNETOPAUSE STUDY WITH A MULTI-FLUID NUMERICAL CODE SIMULATION"



L. Rezeau

G. Belmont

F. Califano

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R. Denton





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### THE PLAN

- 1. Observations side:
  - mixing of different plasmas at same place and time
    - data location along the s/c path and so across the MP structure
      - determination of a coordinate along the depth of the MP
      - determination of the MP orientation
    - plasma decomposition into its temperature-labeled components
  - initialisation methods for simulations
- 2. Simulations side:
  - ◆ 3-Fluids code
  - ◆ 4-Fluid code
- 3. Merging data analysis and simulations



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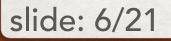
work done

work to be done

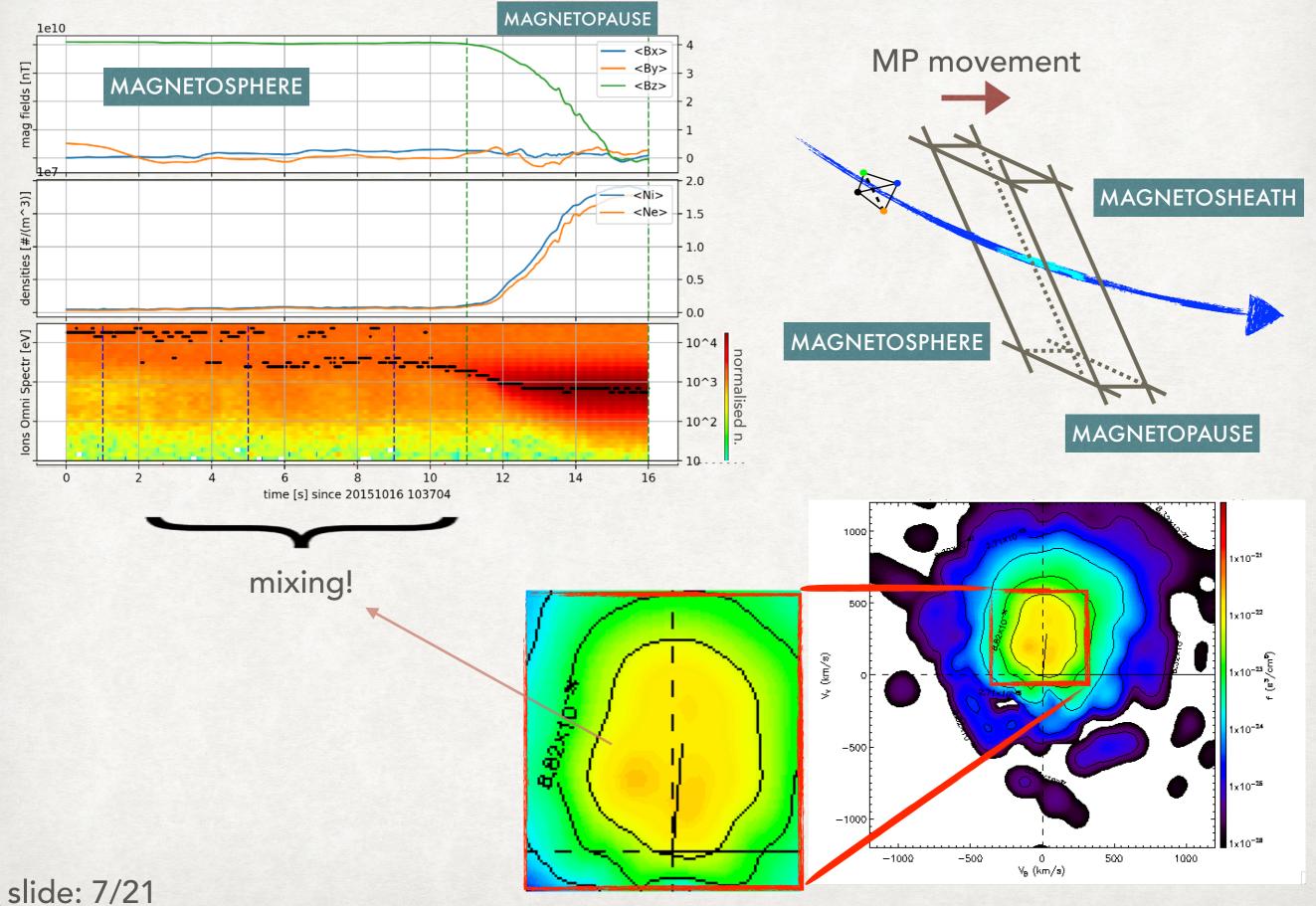
Legend

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### MIXING OF DIFFERENT PLASMAS AT SAME PLACE AND TIME



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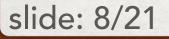
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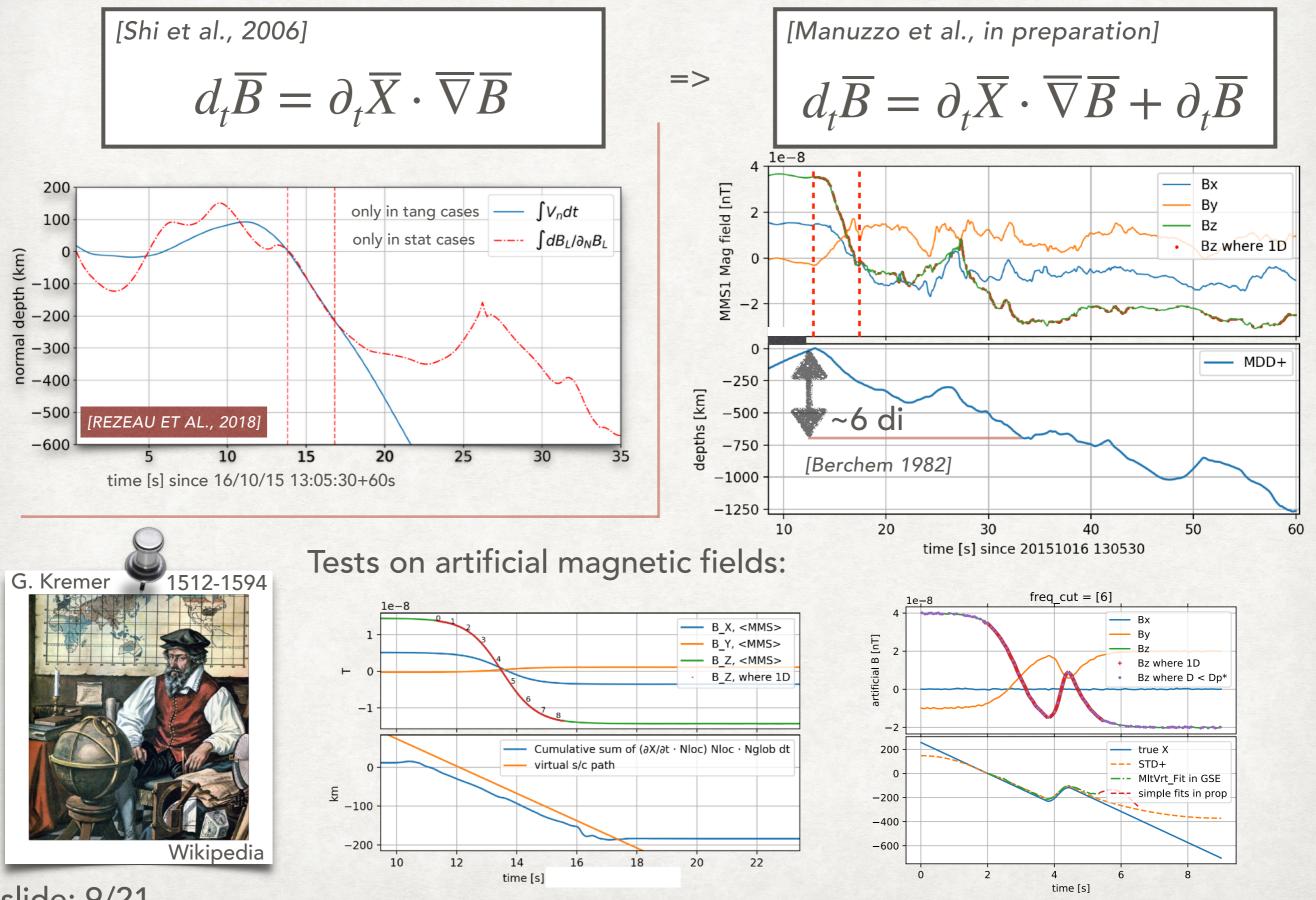
#### Legend

- work done
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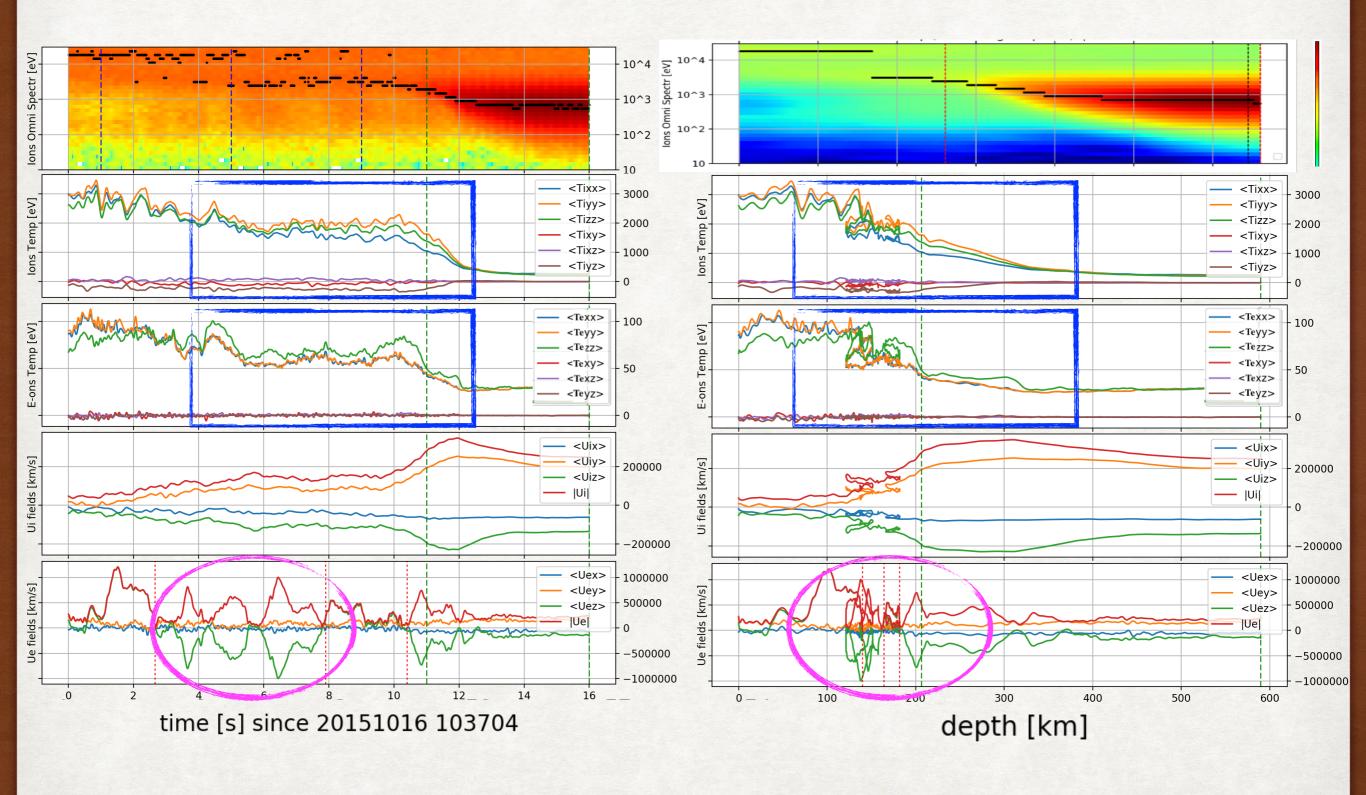


#### DETERMINATION OF A COORDINATE ALONG THE DEPTH OF THE MP



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#### DATA VISUALISATION AS A FUNCTION OF SPACE



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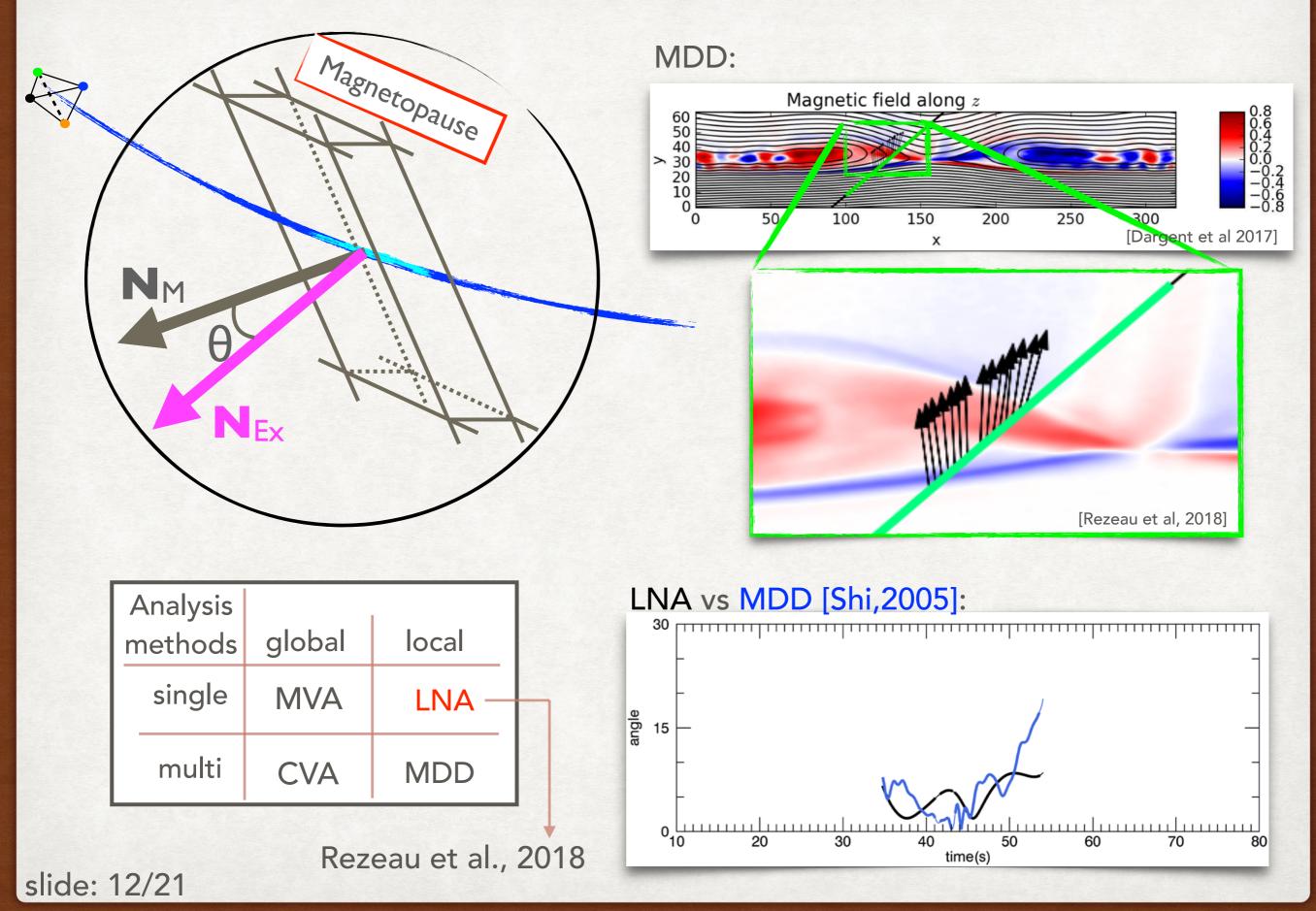
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### DETERMINATION OF THE MAGNETOPAUSE ORIENTATION



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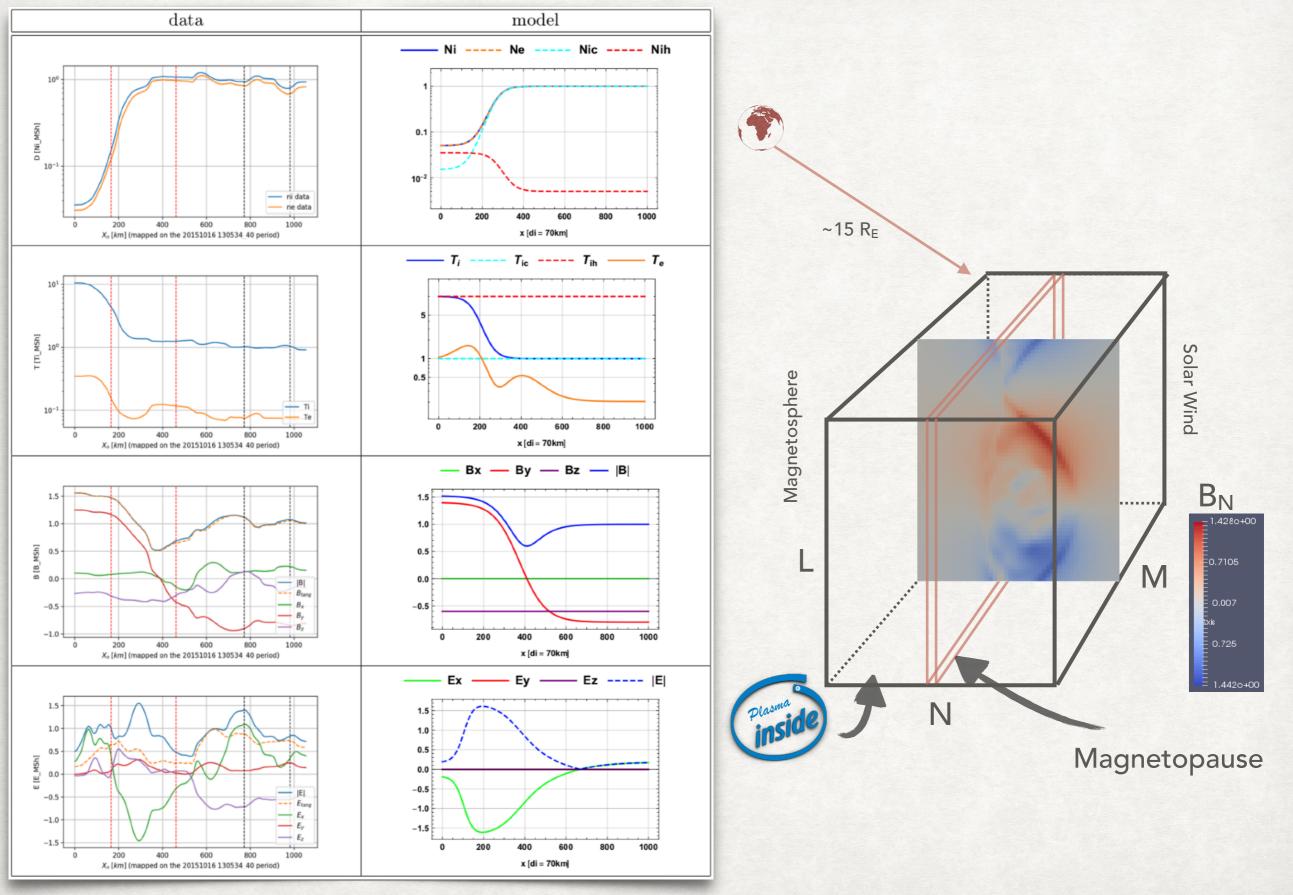
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#### INITIALISATION METHODS FOR SIMULATIONS: A 3FLUID APPROACH



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#### THE 3-FLUID CODE

#### THE ALGORITHM

1. Transport of the density of ions and electrons entropies  $(nS \equiv Pn^{1-\gamma})$ ,  $\alpha = \{ic, ih, e\}$ :

$$\frac{\partial \left(n_{\alpha}S_{\alpha}\right)}{\partial t} + \overline{\nabla} \cdot \left[\mathbf{U}_{\alpha}\left(n_{\alpha}S_{\alpha}\right)\right] = 0$$

2. Ions densities,  $\alpha = \{ic, ih\}$ :

$$\frac{\partial n_{\alpha}}{\partial t} + \overline{\nabla} \cdot (n_{\alpha} \mathbf{U}_{\alpha}) = 0$$

3. Electron density from updated ions densities<sup>1</sup>,  $\alpha = \{ic, ih\}$ :

$$n_e^* = \sum_{\alpha} n_{\alpha}^* \tag{3}$$

(1)

(2)

(5)

(6)

(7)

(8)

(9)

4. Ions momenta,  $\alpha = \{ic, ih\}$ :

$$\frac{\partial \left(n_{\alpha} \mathbf{U}_{\alpha}\right)}{\partial t} + \frac{1}{m_{\alpha}} \overline{\nabla} P_{\alpha} + \overline{\nabla} \cdot \left(n_{\alpha} \mathbf{U}_{\alpha} \mathbf{U}_{\alpha}\right) - \frac{q_{\alpha} n_{\alpha}}{m_{\alpha}} \left(\mathbf{E} + \mathbf{U}_{\alpha} \times \mathbf{B}\right) = 0 \quad (4)$$

5. Ions and electrons pressures,  $\alpha = \{ic, ih, e\}$ :

$$P_{\alpha}^* = \left[n_{\alpha}S_{\alpha}\right]^* \left(n_{\alpha}^*\right)^{\gamma-1}$$

6. Ions velocities,  $\alpha = \{ic, ih\}$ :

$$\mathbf{U}_{\alpha}^{*}=\left(n_{\alpha}\mathbf{U}_{\alpha}\right)^{*}/n_{\alpha}^{*}$$

7. Magnetic field

$$\frac{\partial \mathbf{B}}{\partial t} + \overline{\nabla} \times \mathbf{E} = 0$$

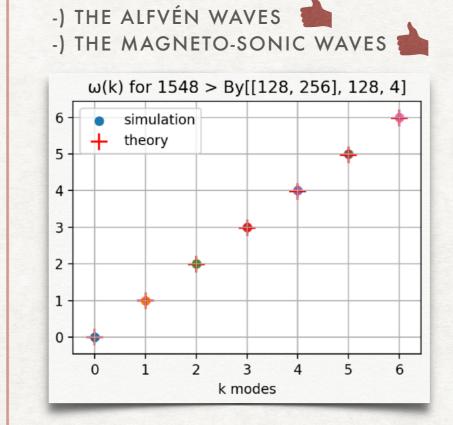
8. Electronic velocities,  $\alpha = \{ic, ih\}$ :

$$\mathbf{U}_{e}^{*} = \left[\sum_{\alpha} \left(n_{\alpha}^{*}\mathbf{U}_{\alpha}^{*}\right) + \frac{1}{q_{e}\mu_{0}}\overline{\nabla}\times\mathbf{B}^{*}\right]/n_{e}^{*}$$

9. Electric field

$$\mathbf{E}^* = -\mathbf{U}_e^* \times \mathbf{B}^* + \frac{1}{q_e n_e^*} \overline{\nabla} P_e^* + \frac{\eta}{\mu_0} \overline{\nabla} \times \mathbf{B}^*$$

-) TIME INDEPENDENCE OF EQUILIBRIUM STATES



THE RESULTS =>

1e-3

20

40

0.0 --0.5

ă <sup>−1.0</sup>

-1.5

-2.0

x profile

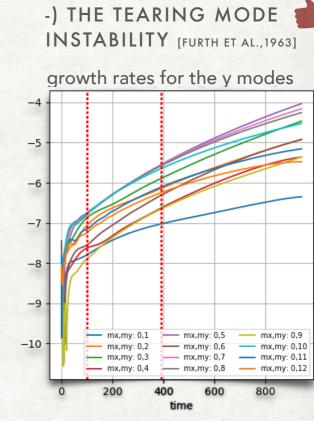
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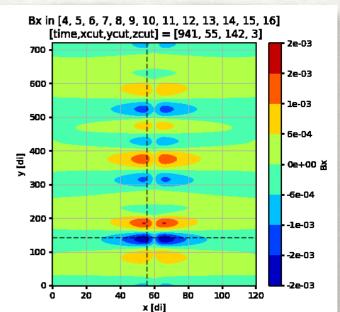
x [di]

RO

100

120







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#### THE 4-FLUID CODE

1) On avance les ions :

#### ... some algorithms are under study:

0. Electric field  

$$\mathbf{E} = -\frac{(n_{e,h}\mathbf{U}_{e,h} + n_{e,c}\mathbf{U}_{e,c})}{n_{e,h} + n_{e,c}} \times \mathbf{B} - \frac{1}{e(n_{e,h} + n_{e,c})}\overline{\nabla} \cdot (P_{e,h} + P_{e,c})\overline{\overline{I}} \qquad (1)$$

1. Ions densities

$$\frac{\partial n_{i,\alpha}}{\partial t} + \overline{\nabla} \cdot (n_{i,\alpha} \mathbf{U}_{i,\alpha}) = 0 \tag{2}$$

2. Ions Momenta

$$\frac{\partial (n_{i,\alpha} \mathbf{U}_{i,\alpha})}{\partial t} + \overline{\nabla} \cdot \left(\frac{1}{m_i} P_{i,\alpha} \mathbf{I} + n_{i,\alpha} \mathbf{U}_{i,\alpha} \mathbf{U}_{i,\alpha}\right) - \frac{e n_{i,\alpha}}{m_i} \left(\mathbf{E} + \mathbf{U}_{i,\alpha} \times \mathbf{B}\right) = 0$$
(3)

3. Ions pressures (trough the conservation of  $S_{i,\alpha} \equiv P_{i,\alpha} n_{i,\alpha}^{1-\gamma}$ )

$$\begin{cases}
4.1: compute \quad \frac{\partial S_{i,\alpha}}{\partial t} + \mathbf{U}_{i,\alpha} \cdot \overline{\nabla} S_{i,\alpha} = 0 \\
4.2: finding \quad S^*_{i,\alpha} \\
4.3: obtain \quad P^*_{i,\alpha} = S^*_{i,\alpha} (n^*_{i,\alpha})^{\gamma - 1}
\end{cases}$$
(4)

4. Magnetic field

$$\frac{\partial \mathbf{B}}{\partial t} + \overline{\nabla} \times \mathbf{E} = 0 \tag{5}$$

5. Electrons densities

$$\begin{cases} 6.1: define & n_{e,+}^* = n_{i,h}^* + n_{i,c}^* \\ 6.2: define & n_{e,-} = n_{e,h} - n_{e,c} \\ 6.3: compute & \frac{\partial n_{e,-}}{\partial t} + \overline{\nabla} \cdot (n_{e,h} \mathbf{U}_{e,h} - n_{e,c} \mathbf{U}_{e,c}) = 0 \\ 6.4: finding & n_{e,-}^* = n_{e,h}^* - n_{e,c}^* \\ 6.5: obtain & n_{e,h}^* = \frac{n_{e,+}^* + n_{e,-}^*}{2} \\ 6.6: obtain & n_{e,c}^* = \frac{n_{e,+}^* - n_{e,-}^*}{2} \end{cases}$$
(6)

6. Electron pressures (trough the conservation of  $S_{e,\alpha}\equiv P_{e,\alpha}n_{e,\alpha}^{1-\gamma}$  )

$$\begin{cases} 7.1: \ compute \quad \frac{\partial S_{e,\alpha}}{\partial t} + \mathbf{U}_{e,\alpha} \cdot \overline{\nabla} S_{e,\alpha} = 0\\ 7.2: \ finding \quad S^*_{e,\alpha}\\ 7.3: \ obtain \qquad P^*_{e,\alpha} = S^*_{e,\alpha} (n^*_{e,\alpha})^{\gamma-1} \end{cases}$$
(7)

7. Evolve E:

slide: 18/21

$$\begin{cases} 8.1: with \qquad \mathbf{J} = e\left(n_{i,h}\mathbf{U}_{i,h} + n_{i,c}\mathbf{U}_{i,c} - n_{e,h}\mathbf{U}_{e,h} + n_{e,c}\mathbf{U}_{e,c}\right)\\ 8.1: compute \quad \frac{\partial \mathbf{E}}{\partial t} - c^2\left(\overline{\nabla} \times \mathbf{B} - \mu_0 \mathbf{J}\right) = 0 \end{cases}$$
(8)

8. Electron moments

Problems arise here... how to find  $n_{e,h}^* \mathbf{U}_{e,h}^*$  and  $n_{e,c}^* \mathbf{U}_{e,c}^*$ ?

9. Go to point 1 and repeat the loop.

$$\begin{array}{cccc} -\operatorname{continuit\acute{e}s}: & n_{i1}^{1} = n_{i1} - \nabla.(\boldsymbol{\varphi}_{l1}) & (1) \\ & n_{l2}^{2} = n_{l2} - \nabla.(\boldsymbol{\varphi}_{l2}) & (2) \\ -\operatorname{momentums}: & n_{l}\boldsymbol{\varphi}_{l1}^{*} = m_{l}\boldsymbol{\varphi}_{l1} - \nabla.(p_{l1}\mathbf{I} + m_{l}\frac{\boldsymbol{\varphi}_{l1} \boldsymbol{\varphi}_{l1}}{n_{l1}}) + q[n_{l1}\mathbf{E} + \boldsymbol{\varphi}_{l1} \times \mathbf{B}] & (3) \\ & n_{l}\boldsymbol{\varphi}_{l2}^{*} = m_{l}\boldsymbol{\varphi}_{l2} - \nabla.(p_{l2}\mathbf{I} + m_{l}\frac{\boldsymbol{\varphi}_{l2} \boldsymbol{\varphi}_{l2}}{n_{l2}}) + q[n_{l2}\mathbf{E} + \boldsymbol{\varphi}_{l2} \times \mathbf{B}] & (4) \\ -\operatorname{fermetures}: & p_{l1}^{*} = K_{l1}(n_{l1}^{*})^{\gamma_{l1}} & (5) \\ & p_{l2}^{*} = K_{l2}(n_{l2}^{*})^{\gamma_{l2}} & (6) \\ \end{array}$$

$$\begin{array}{c} \mathbf{2) \ On \ avance \ B: \\ -\operatorname{Faraday:} & \mathbf{B}^{*} = \mathbf{B} - \nabla \times (\mathbf{E}) & (7) \\ 3) \ On \ avance \ les \ densités \ électroniques \ : \\ & - \ différence \ des \ continuités: & n_{e2}^{*} - n_{e1}^{*} = n_{e2} - n_{e1} - \nabla.(\boldsymbol{\varphi}_{e2} - \boldsymbol{\varphi}_{e1}) & (8) \\ & - \ quasi-neutralité: & n_{e2}^{*} + n_{e1}^{*} = n_{i2}^{*} + n_{i1}^{*} & (9) \\ On \ en \ déduit \ n_{e1}^{*} \ et \ n_{e2}^{*} \ séparément. \\ - \ fermetures: & p_{e1}^{*} = K_{e1}(n_{e1}^{*})^{\gamma_{e1}} & (10) \\ & p_{e2}^{*} = K_{e2}(n_{e2}^{*})^{\gamma_{e2}} & (11) \\ \end{array}$$

 $\boldsymbol{\varphi}_{e2}^* + \boldsymbol{\varphi}_{e1}^* = \boldsymbol{\varphi}_{i2}^* + \boldsymbol{\varphi}_{i1}^* - \frac{1}{q} \nabla \times \mathbf{B}^* / \mu_0$ (12)

- différence des momentums :

$$m_{e}(\boldsymbol{\varphi}_{e2}^{*} - \boldsymbol{\varphi}_{e1}^{*}) = m_{e}(\boldsymbol{\varphi}_{e2} - \boldsymbol{\varphi}_{e1}) - \nabla \left[ (p_{e2} - p_{e1})\mathbf{I} + m_{e}\left(\frac{\boldsymbol{\varphi}_{e2}\boldsymbol{\varphi}_{e2}}{n_{e2}} - \frac{\boldsymbol{\varphi}_{e1}\boldsymbol{\varphi}_{e1}}{n_{e1}}\right) \right] -q\left[ (n_{e2} - n_{e1})\mathbf{E} + (\boldsymbol{\varphi}_{e2} - \boldsymbol{\varphi}_{e1}) \times \mathbf{B} \right]$$
(13)

De (12) et (13), on déduit  $\boldsymbol{\varphi}_{e2}^*$  et  $\boldsymbol{\varphi}_{e1}^*$  séparément

- somme des momentums (au moment *t*+1):

$$m_{e}(\boldsymbol{\varphi}_{e2}^{**} + \boldsymbol{\varphi}_{e1}^{**}) = m_{e}(\boldsymbol{\varphi}_{e2}^{*} + \boldsymbol{\varphi}_{e1}^{*}) - \nabla \left[ (p_{e2}^{*} + p_{e1}^{*})\mathbf{I} + m_{e}\left(\frac{\boldsymbol{\varphi}_{e2}^{*}\boldsymbol{\varphi}_{e2}^{*}}{n_{e2}^{*}} + \frac{\boldsymbol{\varphi}_{e1}^{*}\boldsymbol{\varphi}_{e1}^{*}}{n_{e1}^{*}}\right) \right] - q\left[ (n_{e2}^{*} + n_{e1}^{*})\mathbf{E}^{*} + (\boldsymbol{\varphi}_{e2}^{*} + \boldsymbol{\varphi}_{e1}^{*}) \times \mathbf{B}^{*} \right]$$
(14)

#### work in progress...

#### SUMMARY

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1.Already published:

1. Rezeau, Belmont, Manuzzo, Aunai, Dargent, 2018, "Analysing the magnetopause internal structure: new possibilities offered by MMS tested in a case study", J. Geophys. Res., 122.

2.Previewed:

- 1. A study about the Magnetopause kinematics and thickness and about the localisation of spacecraft/data with respect to the magnetopause itself
- 2. A paper on the new 3-fluid and 4-fluid codes
- 3. A paper as co-author in magnetopause frame determination (Denton, work in progress)

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### AGU 2017 IN NEW ORLEANS

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11-15 Dec. 2017 What will 404 discover?

#### **CAGU FALL MEETING** Determining the Thickness and the Sub-Structure Details of the Magnetopause from MMS Data R. Manuzzo<sup>1,2</sup>, G. Belmont<sup>1</sup>, L. Rezeau<sup>1</sup>, F. Califano<sup>2</sup>

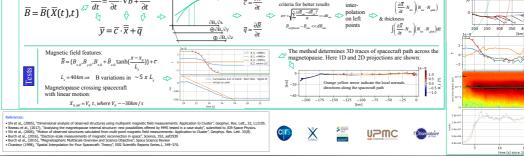
I: Université Pierre et Marie Curie, Paris: 2:Università di Pisa, Ita

1. FUDULOU. The magnetopause thickness, like its mean location, is a notion that can have different meanings depending which parameters are considered (magnetic field or plasma properties). In any case, all the determinations have been done, up to now, considering the magnetopause boundary as a structure strictly stationary and ID (or with a simple curvature). These determinations have show not be wery sensitive to the accuracy of the normal direction, because it affects the projection of the quantities of spaceraft path across the magnetopause in order to compute the dimensionality, the local normals and spaceraft path across the magnetopause in order to compute the dimensionality, the local normals and spaceraft path across the magnetopause in order to compute the dimensionality, the local normals and spaceraft path across the magnetopause in order to compute the dimensionality, the local normals and spaceraft path across the magnetopause in order to compute the dimensionality, the local normals and spaceraft path across the magnetopause in order to compute the dimensionality. II. The problem: The observational studies in space plasma physics deal with time series of data measured by spacecraft crossing plasma structures. Scientists often need to deduce the geometrical features (orientations, dimensions) of these structures being key features of phenomena happening inside them. The time-to-space translation process often involves strong hypotheses never fully satisfied in real conditions. Furthermore, the help of data-initialised simulations can be sometimes misleading because the evolution of the real plasma can differ substantially from its virtual counterpart due to the over-simplified boundary conditions (kinetic and Vlasov simulations) or the over-simplified governing equations (fluid simulations) used in simulations.

III. The solution consists in recovering the space plasma structure geometry from the magnetic field by means of a time integration of the aX/at term contained within the Lagrangian derivative of  $\mathbf{B}=\mathbf{B}(\mathbf{x}(t); t)$ :  $\mathbf{B}/dt = \partial X/dt + \nabla \mathbf{B} + \partial \mathbf{B}/dt$ . As a matter of facts, thanks to the multi-spacecraft NASA's Magnetometer MultiScale mission (MMS) [Burch et al., 2015], we have the capabilities to recover the  $\nabla \mathbf{B}$  term [Chanteur, 1998].

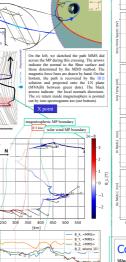
III.a: We implemented the Spatio-Temporal difference method [Shi et al. 2006], by means of the inversion of dB/dt =  $\partial X/\partial t \cdot \nabla B$  (neglecting the  $\partial B/\partial t$ term due to the hypothesis of stationarity). Care is needed during the inversion procedure because the singularities of  $\nabla B$  (due to 1D structures and experimental fluctuations which causes the zeros of the adjoint of  $\nabla B$  to be elsewhere with respect to the zeros of det( $\nabla B$ ).

 $\left(\frac{\partial \overline{X}}{\partial t} \cdot N_{hc}\right) \left(N_{hc} \cdot N_{pb0}\right)$  $\overline{B} = \overline{B}(\overline{X}(t), t) \iff \frac{d\overline{B}}{dt} = \frac{\partial \overline{X}}{\partial t} \cdot \overline{\nabla}\overline{B} + \frac{\partial \overline{B}}{\partial t}$ added to modify the t which  $(\nabla B)^*=0$  to be & thicknes  $\frac{\partial \overline{X}}{\partial t} \cdot \overline{N}_{hc} = \left\{ \left\{ \left[ \frac{d\overline{B}}{dt} \cdot (\overline{\nabla}\overline{B})^{-1} \det(\overline{\nabla}\overline{B}) \right] \cdot \overline{N}_{hc} \right\} + \epsilon \right\} / \det(\overline{\nabla}\overline{B})$  $\int \left(\frac{\partial \overline{X}}{\partial t} \cdot N_{bc}\right) \left(N_{bc} \cdot N_{gbb}\right)$ Magnetic field model  $\overline{B} = \{\sim 0, \sim 0, \widetilde{B}_{x,0} \tanh(\frac{x - x_0}{r})\}$ ... to be compared with the real data coming from 16/10/15, 13:05:30  $L_1 = 120 km \Rightarrow B$  variations in  $\sim 5 x L_1$ agnetopause crossing spacecraft th standing & back and forth motio =>  $X_{NMP} = V_N(t) t$ , where  $V_{NMP} \sim -90 km/s$ III.b: Under the hypothesis of constancy of the  $\partial X/\partial t$  and the  $\partial B/\partial t$  terms during periods of a few data points (instead of the hypothesis of stationarity using begins to the use importance of antenney of the output and use owner than used on the importance of antenney of the owner and use importance of antenney of the owner owner that the importance of antenney of the owner owne path (work in progress)  $\frac{d\overline{B}}{d\overline{B}} = \frac{\partial \overline{X}}{\partial \overline{X}} \cdot \overline{\nabla}\overline{B} + \frac{\partial \overline{B}}{\partial \overline{B}}$ normal velocity criteria for better results dt 7 dt dt



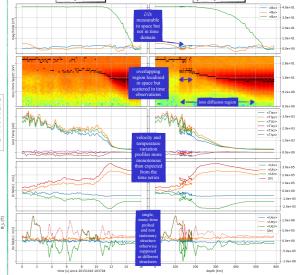
We concentrated on a very famous case study (the crossing used by Burch et al. (2016) to compute the LMN frame for the close reconnection event) in order to test our method. We found a more complex MP back and forth motion than expected and we got a determination of the normal(s) lirection(s) along the s/c path which seems more n-line than the kine atics sketched by Burch et al. (2016) and that show how labelling an entire crossing with only one global normal can be a rough procedure in view of data projection

New insights onto the 16/10/15, 13:05:30 + 60s event



The space representation of data for another case study (16/10/15, 10:37:04 + 16s event) unlocked observations such as The space representation to data for another case study LOV(15, LD-21/04 + Los even) functioned observations such as the drift terms measurements, the precise pinpointing of mining sub-byers where the magnetospheric and solar wind pinbiliations and the revealing of a many-time probed and nonstation of pinpoint studences in there is develored as different studences and the revealing of a many-time probed and non stationary structure otherwise simply believed as different structures. time representation space representation

New insights in the magnetopause sub-laver structure: the 16/10/15, 10:37:04 + 16s event



Conclusions and perspectives:

V. Some results: Here are collected two from the main results obtained in computing the MMS path trough the Magnetopause during two crossings

We presented two methods to analyse the internal structure of the magnetopause, the two differing on the hypothesis about the importance of the  $\partial/\partial t$  term and in the number of data point to be processed each step. No hypothesis has been done a priori on the dimensionality (but for 1D or 2D projections of the resulting 3D trajectories), allowing these procedures to be applied to real cases with a very few resulting 3D trajectones), allowing these procedures to be applied to real cases with a very tew restrictions. We performed tests to validate the methods on artificial magnetic fields affected by noise and probed as if the magnetopause overcame MMS with constant velocities, sudden standings and/or back and forth motions. We applied the methods onto two real cases study during 16/10/15: for the 13:05:30+60s case we obtained a new perpective on the MP kinematics; for the 10:37:04+16s case we got new understandings about the sub-layers structures, dimensions and respective locations. We aims to ameliorate the methods, by cross-checks, error computations and applications to fluid simulation data probed by virtual spacecraft. Finally, the methods will allow to measure the plasma flux flowing across the magnetonewise. magnetonause

...BUT ALSO EGU 2018, AGU 2018...

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### FORMATIONS

- 1. [30h] Python coding (online certified course)
- 2. [1 week] "Description fluide et cinétique des plasmas" 2017, Observatoire de Meudon, Paris
- 3. [2 weeks] École de Physique des Houches 2017
- 4. [1 week] Cineca Summer School on High Performance Computing, 2017, Bologna, Italy
- 5. [2.5 h] Ethique et intégrité scientifique



### End

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- 3. Stone & al., "The Advanced Composition Explorer", Space Science Reviews 86-1, 1, (1998)
- 4. Rezeau et al., 2018, "Analyzing the magnetopause internal structure: new possibilities offered by MMS tested in a case study", J. Geophys. Res., 122.
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How to discover several species in one distribution function? -) Reminder: the physical meaning of a distribution

$$f_{MB}(v) = n(\frac{m}{2\pi kT})^{3/2} Exp(-\frac{m(v-u)^2}{2kT})$$

-) 1V Hermite transform representation<sup>[\*]</sup>:

$$\begin{cases} f_v = \sum_m f_m \psi_m(v) \\ \int_{\infty}^{\infty} \psi_m \psi_l = \delta_{ml} \\ \psi_m(v) = \frac{H(\frac{v-u}{v_{th}})}{\sqrt{2^m m!} \sqrt{\pi} v_{th}} exp(-\frac{m(v-u)^2}{2v_{th}^2}) \\ H_m(v) = (-1)^m e^{v^2} \frac{d^m}{dv^m} e^{-v^2} \end{cases}$$

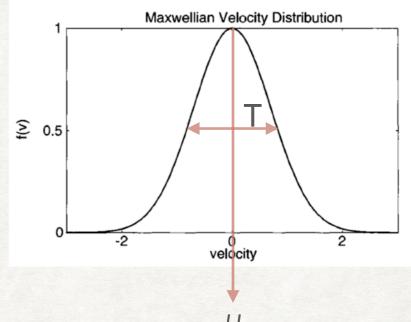
$$N.B.: \psi_{m=0}(v) \propto F_{MB}$$

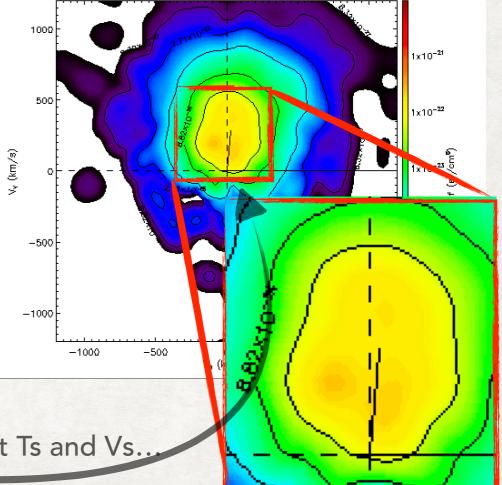
-) idea: apply the HTR to the 3V distr. func.:

$$\psi_m(\mathbf{v}) = \psi(m_x, v_x)\psi(m_y, v_y)\psi(m_z, v_z)$$

allowing for the presence of different species having different Ts and Vs...

[\*] Servidio et al., "Magnetospheric MultiScale (MMS) observation of plasma velocity-space cascade: Hermite representation and theory", Phys. Plasmas, arXiv:1707.08180v1 (2017)





return

The Minimum Directional Derivative method (MDD, Shi [2005])...

$$\overline{\overline{G}} = \nabla \overline{B} \qquad \overline{\overline{R}} = \overline{D} = \frac{\partial \overline{B}}{\partial \overline{n}}$$

If **n** = **N** was the invariant direction along which all the parameters remain constant =>  $D^2 = 0$ 

$$\begin{array}{lll} diagonalization(\overline{\overline{G}}\,\overline{\overline{G}}^{T}) & \left\{ \begin{matrix} \lambda_{1} & \overline{v}_{1} & \text{maximum} \\ \lambda_{2} & \overline{v}_{2} & \text{intermediate} \\ \lambda_{3} & \overline{v}_{3} & \text{minimum} \end{matrix} \right\} & \text{of } \mathsf{D} & \begin{array}{ll} \lambda_{1} \gg \lambda_{2} \wedge \lambda_{3} & => 1\mathsf{D} \ \text{str.} \\ \lambda_{1} \sim \lambda_{2} \gg \lambda_{3} & => 2\mathsf{D} \ \text{str.} \\ \lambda_{1} \sim \lambda_{2} \sim \lambda_{3} & => 3\mathsf{D} \ \text{str.} \end{matrix}$$

... our Local Normal Aanlysis method (LNA, [Rezeau, 2018])

Hp.: stationarity and 1D

R

$$\overline{J} = \overline{N} \times \partial_N \overline{B} \perp \partial_N \overline{B} = -V_N \cdot \partial_N \overline{B}$$

### Methods used to find the normal to the magnetopause

Common Hypotheses: magnetopause =1D and stationary layer  $\Rightarrow$ 

$$\mathbf{B}_n \neq \mathbf{B}_n(N,t)$$

Single spacecraft Multi spacecraft BV Minimum Variance **Constant Velocity** Approach method Analysis method method 1) Find **N** in order to minimize: 1) Other hypotheses needed: lagnetopause -) no flow through the  $\sum \left\| \left( \mathbf{B}_{i} - ar{\mathbf{B}} 
ight) \cdot \mathbf{N} 
ight\|^{2}$ magnetopause -) B behaves like:  $egin{aligned} B_L &= B_{0L} cos(lpha) \ B_N &= B_{0N} \ B_M &= B_{0M} sin(lpha) \end{aligned}$ 2) How to? Diagonalize  $M_{\mu
u}=raket{B_{\mu}B_{
u}}-raket{B_{\mu}}raket{B_{
u}}{B_{
u}}$  $\mu,
u=x,y,z$ where:  $\triangle t_{\alpha\beta}$  $lpha = lpha_1 + (lpha_2 - lpha_1) \, rac{{\sf N}}{y_{max}}$  $\mathbf{r}_{lphaeta}$ and

$$(\mathbf{V} \triangle t_{\alpha\beta}) \cdot \mathbf{N} = \mathbf{r}_{\alpha\beta} \cdot \mathbf{N}$$

Pro: Less hypoth, Simple, Mean Vs: Problem with MMS data

return

3) Results:

- -) eigenvalues:  $\lambda_1$ ,  $\lambda_2$ ,  $\lambda_3$
- -) eigenvectors: L, M, N

Pro: simple hypotheses Vs: bad determination if not 1D ſ

$$\mathsf{N} = \int_{crossing} \mathbf{V}_{BF}(t) \cdot \mathbf{N} \, dt$$

Pro: particle data involved too Vs: Hypotheses