

UNIVERSITY OF PISA Ph.D. graduate school in Physics

Mutual Information in molecular liquids

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Outline of the presentation

• Context

- Glass transition
- Investigated system and method

Displacement-Displacement correlation

- Average number of MI-correlated particles
- Standard deviation
- Correlation with structure
- Conclusion and future work



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Liquids cooled fast enough avoid crystallization and fall out of equilibrium





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Liquids cooled fast enough avoid crystallization and fall out of equilibrium



Structure does not change undergoing the glass transition



Coarse-grained polymeric model offer a good framework for numerical simulation

Bond interaction

 $U_{bond}(l) = k(l - l_0)^2$

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Lennard-Jones Interaction

$$U_{LJ}(r) = 4\epsilon \left[\left(\frac{\sigma}{r} \right)^{12} - \left(\frac{\sigma}{r} \right)^{6} \right] \xrightarrow{0.75}_{0.50}_{0.00}_{-0.25}_{-0.25}_{-0.50}_{-0.75}_{-$$







Readapted from: S. Bernini, F. Puosi, and D. Leporini, J. Phys. Condens.Matter 29, 135101 (2017).



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Fast mobility (<u²>) and relaxation (τ_{α}) are correlated

Molecular dynamics simulations

Experiments



$$\log \tau_{\alpha} = \alpha + \beta \left(\frac{1}{\langle u^2 \rangle}\right) + \gamma \left(\frac{1}{\langle u^2 \rangle}\right)^2$$



$$\log \tau_{\alpha} = \alpha + \beta \left(\frac{\langle u_g^2 \rangle}{\langle u^2 \rangle} \right) + \gamma \left(\frac{\langle u_g^2 \rangle}{\langle u^2 \rangle} \right)^2$$

Mutual information gives the general degree of dependence between two random variables

Shannon Entropy

$$H(X) = -\int dx \ p(x) \log p(x)$$

it quantify the randomness of a random variable



X and Y, random variables with joint probability distribution p(x, y)H(X, Y)if X and Y are independent p(x, y) = p(x)p(y)H(X, Y) = H(X) + H(y)if not I(X, Y) = H(X) + H(Y) - H(X, Y)

Mutual Information

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X and Y, random variables with joint probability distribution p(x, y)H(X, Y)if X and Y are independent p(x, y) = p(x)p(y)H(X, Y) = H(X) + H(y)if not $I(X,Y) = \left[\int dx dy \ p(x,y) \log \left| \frac{p(x,y)}{p(x)p(y)} \right| \right]$

Mutual Information



Pearson correlation coefficient

$$C(X, Y) = \frac{\langle (x - \langle x \rangle)(y - \langle y \rangle) \rangle}{\sigma_X \sigma_Y} = ?$$



Pearson correlation coefficient

$$C(X,Y) = \frac{\langle (x - \langle x \rangle)(y - \langle y \rangle) \rangle}{\sigma_X \sigma_Y} = 0$$

X and Y are independent !?



Pearson correlation coefficient

$$C(X,Y) = \frac{\langle (x - \langle x \rangle)(y - \langle y \rangle) \rangle}{\sigma_X \sigma_Y} = 0$$

X and Y are independent !?



 $I(X, Y) \neq 0$

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We built three set of iso-relaxing states



We investigate displacement correlation between pairs of particle through mutual information



Iso-configurational ensemble

- Take an equilibrated configuration
- Erase all the velocities
- Re-assign velocities according to a M-B at the same temperature

Displacement distribution in the iso-configurational ensemble



 $I_{ij}(t) = I(\delta \vec{r}_i(t), \delta \vec{r}_j(t))$

Particles above threshold are said to be significantly correlated



At each time we can compute the distribution of the number of particle correlated to a tagged one







The maximum value of $\overline{n}(t)$ do not increase with τ_{α}





Peak of correlation

Loss of correlation



The maximum value of $\overline{n}(t)$ do not increase with τ_{α}



Loss of correlation



The maximum value of $\overline{n}(t)$ do not increase with τ_{α}



Build up correlation







Loss of correlation

Mutual information correlation length increases weakly approaching the glass transition



Conventional displacement correlation functions

$$C_{\bar{u}}(r,t) = \left\langle \hat{\mathbf{u}}_{i}(t_{0},t) \cdot \hat{\mathbf{u}}_{j}(t_{0},t) \right\rangle \longrightarrow \xi_{\bar{u}}$$

$$C_{\delta u}(r,t) = \frac{\left\langle \delta u_{i}(t_{0},t) \delta u_{j}(t_{0},t) \right\rangle}{\left\langle [\delta u(t_{0},t)]^{2} \right\rangle} \longrightarrow \xi_{\delta u}$$

Standard deviation reveals two length scales



we need to characterize monomers on the basis of $n_i(t)$

High $n_i(\tau_{early})$ implies high mobility High $n_i(\tau_{late})$ implies low mobility





τ_{early} and τ_{late} mark two modes of relaxation







Build up correlation

Peak of correlation

Loss of correlation



Isolate these two modes

 $\frac{Threshold}{\langle n_i \rangle + 2\sigma_{n_i}}$

au_{early} and au_{late} are the structural relaxation times of the two fractions



Scaling has already proven to hold for subset of particles of a bulk system



M. Becchi, A. Giuntoli, and D. Leporini, SoftMatter 14, 8814 (2018)

Scaling holds for the early and late-relaxing fractions



Local structure correlates with the two populations

• Local density



$$c(n_i(t), \delta r_i^2(t)) = \frac{\langle (n_i(t) - \langle n_i(t) \rangle)(\rho_i(t) - \langle \rho_i(t) \rangle) \rangle}{\sigma_{n_i} \sigma_{\rho_i}}$$

Average over populations for all the states

Local structure correlates with the two populations

• Topological characterization of local structure



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Conclusion and future work

- Investigated displacement-displacement correlation of pair of particle through mutual information
- Average number of correlated particle \rightarrow MI length scale

A. Tripodo, A. Giuntoli, M. Malvaldi, and D. Leporini, SoftMatter, (2019)

• Standard deviation \rightarrow Two modes of relaxation and scaling

Soon-to-be-published

In the immediate future

- Correlation of the modulus of the displacement
- Connection of τ_{late} with τ_{ee}