

Ground-state configuration of ^{23}O investigated via Coulomb breakup

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Contents

Studies of light neutron-rich nuclei spectroscopy at FRS/LAND:

∅ Electromagnetic excitations as **spectroscopic tool**

- theoretical description
- experimental method

Results on single particle properties of the ground state of neutron-rich nuclei :

∅ Discussion of experimental results and comparison with the models

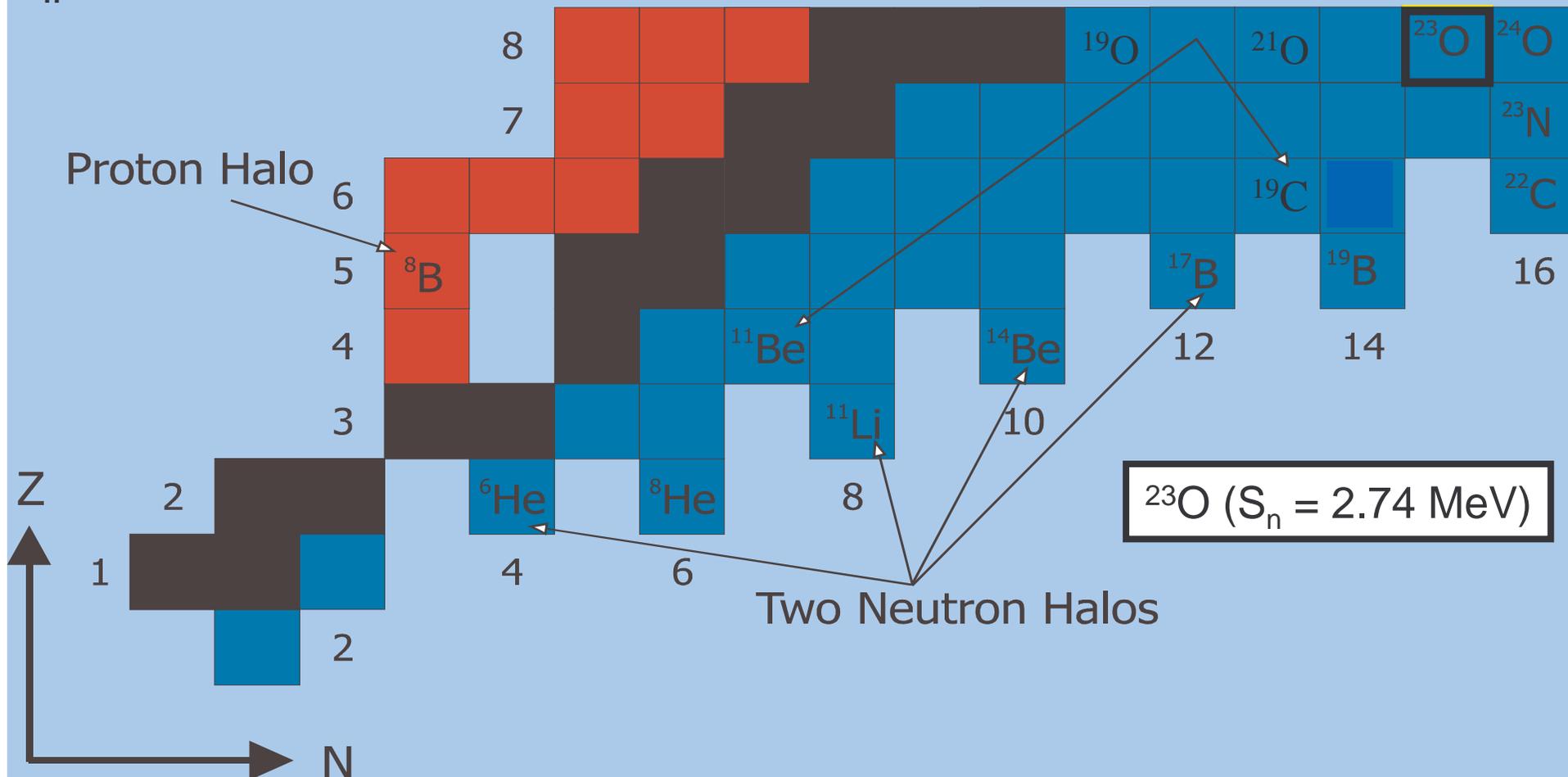
- ^{23}O case

Approaching the neutron drip line

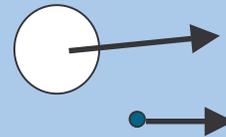
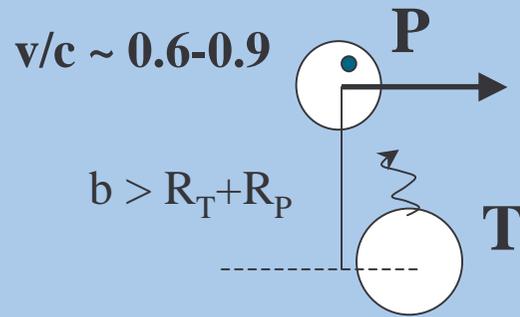
Unusual combination of **large isospin** and **weak binding energy**
→ large diffuseness of nuclear surface (*Halo*)

$S_n < 1$ MeV for halo nuclei

One Neutron Halos



Electromagnetic dissociation



Spectrum of *virtual photons* up to $E_{\text{max}} = \gamma v \hbar / b \approx 25 \text{ MeV}$

∅ **Small nuclear contribution to the cross section** (estimated with C target runs) $\rightarrow \sigma_{\text{e.m.}} = \sigma_{\text{tot}} - \sigma_{\text{nucl}}$

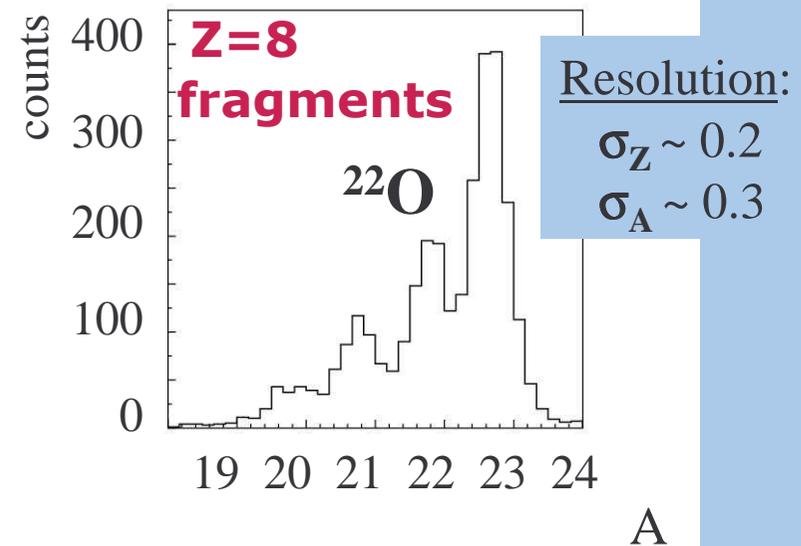
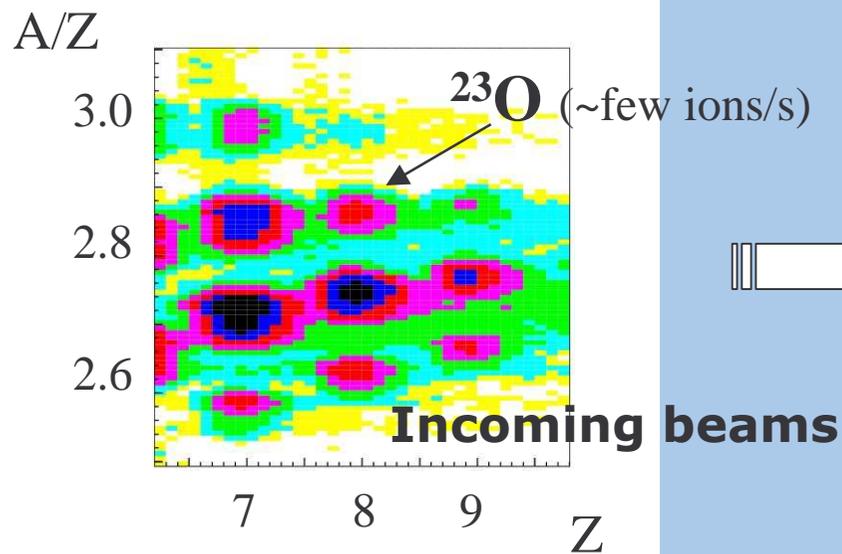
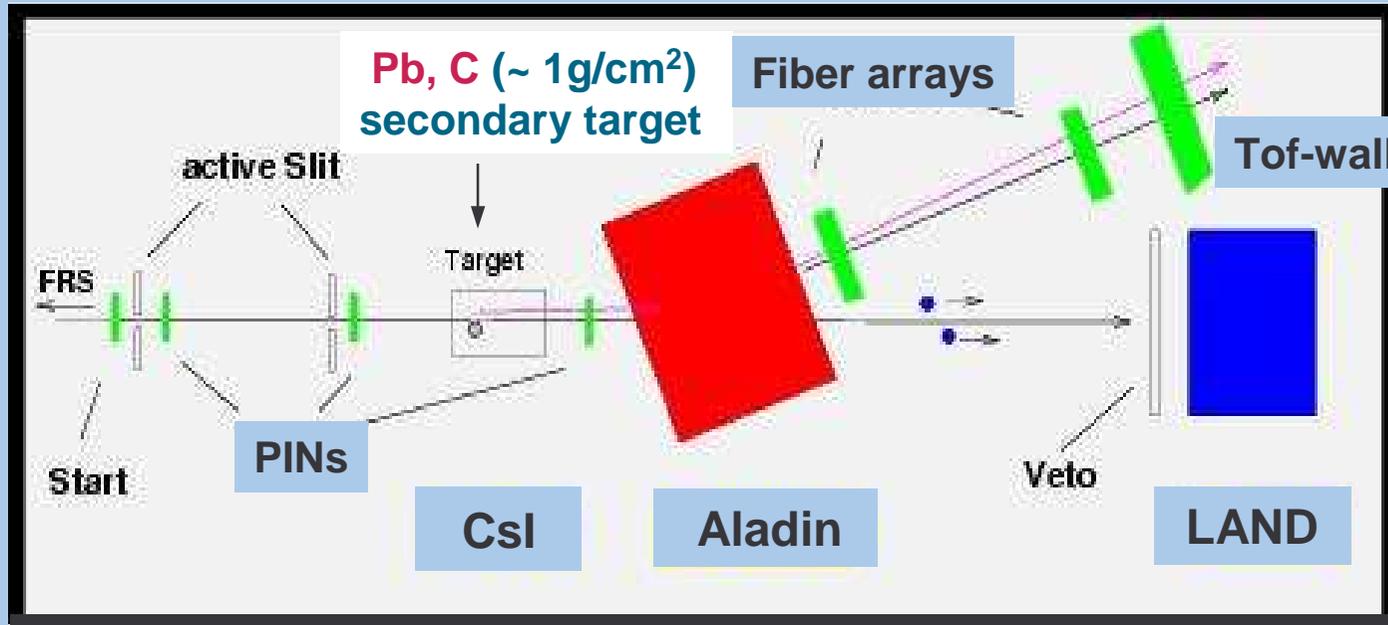
∅ **E1 is dominating over other multipolarities due to the small effective charge**

(S. Typel and G. Baur, PRC 64(2001)024601)

Determination of photon energy (excitation energy) via kinematically complete measurement of the momenta of all outgoing particles (**invariant mass method**)

LAND experimental setup

At $E_{\text{inc}} = 422 \text{ MeV/u}$



Direct breakup model

After projection on *core* states I_c^π , identified by means of *γ -ray coincidences* :

$$\frac{d\sigma}{dE^*}(I_c^\pi) = \frac{16\pi^3}{9\hbar c} N_{E1}(E^*) \sum_{nlj} C^2 S(I_c^\pi, nlj) \sum_m \left| \left\langle \mathbf{q} \left| \frac{Ze}{A} r Y_m^1 \right| \psi_{nlj}(\mathbf{r}) \right\rangle \right|^2$$

↑ **E1 matrix element**
 $N_{E1}(E^*)$ calculated in semiclassical approx.

↓
spectroscopic factor

The differential cross section for e. m. excitations can provide information on the quantum numbers and spectroscopic factors of ground state configuration

Advantages

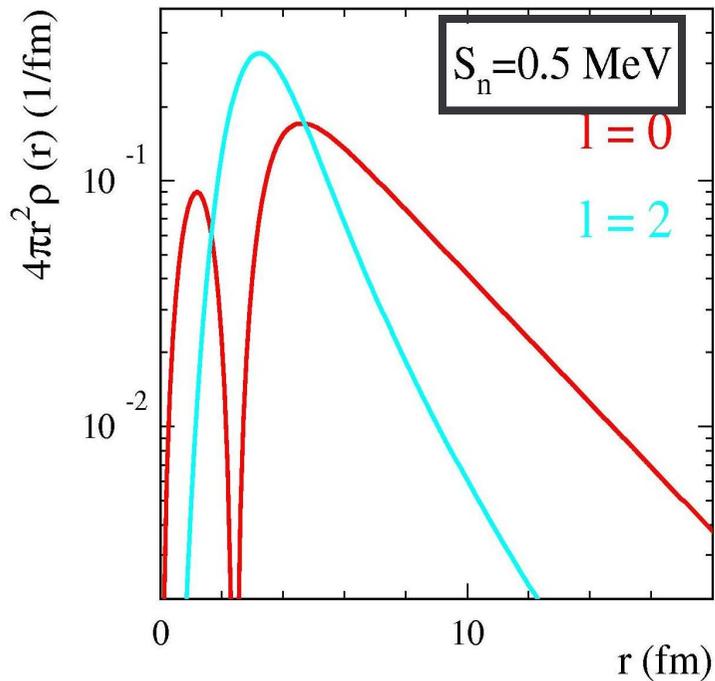
- ∅ interaction is well known
- ∅ high energy approx.
- ∅ sensitivity to low l values

Limitations

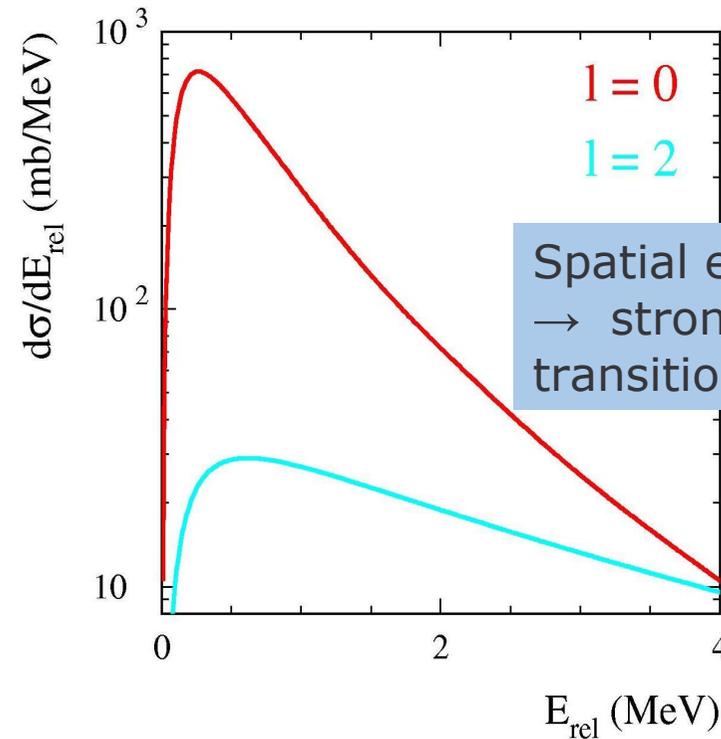
- ∅ core excited states to be identified
- ∅ in case of not so weakly bound breakup systems ?

Low-lying E1 strength as spectroscopic tool

Density distribution



Differential cross section



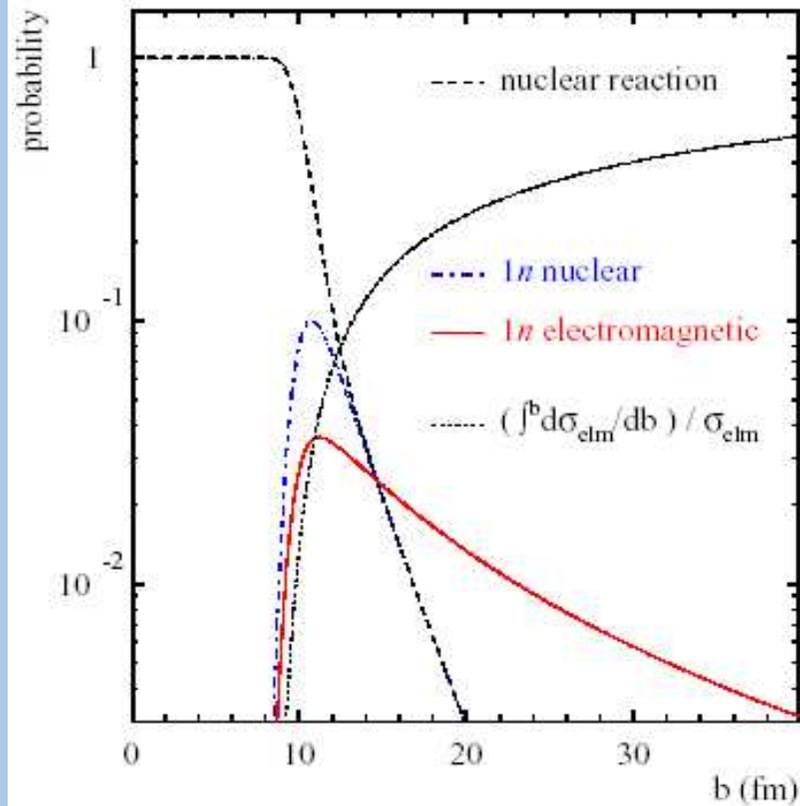
Spatial extension (*Halo*)
→ strong non-resonant transition ($\sim 100 \text{ mb/MeV}$)

shape of differential cross section
→ γ -ray coincidences
cross section

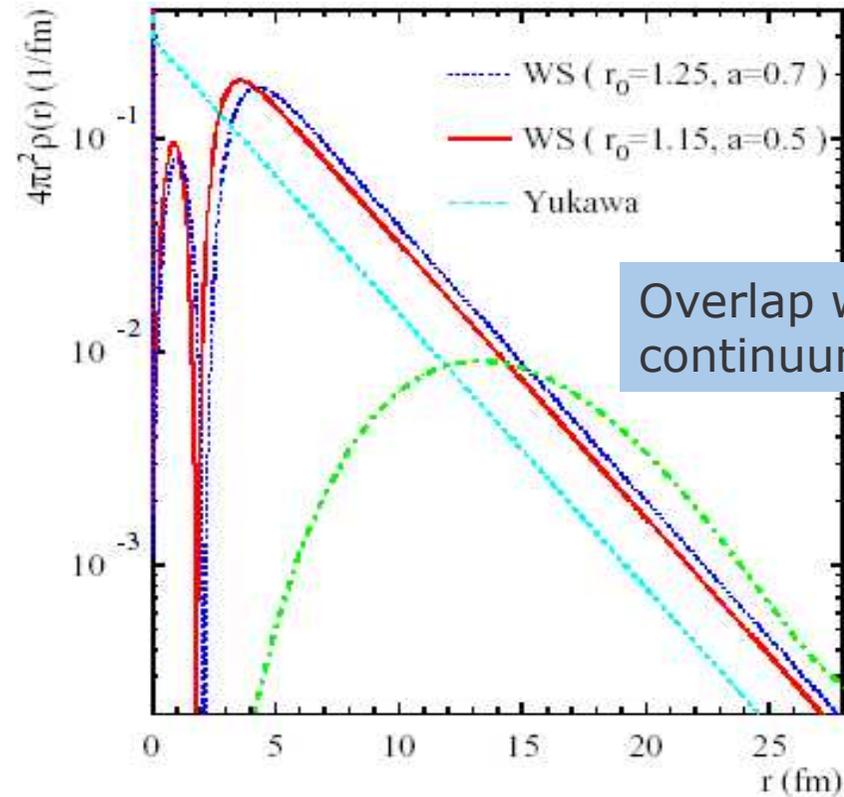
→ orbital angular momentum /
→ identification of core state
→ spectroscopic factor

Sensitivity of Coulomb breakup

Reaction probabilities



Halo-neutron densities



Overlap with continuum w.f.

R. Palit et al., PRC 68(2003)034318

Sensitivity to low-density tail of the wave function

Information available on ^{23}O nucleus

s-wave g.s. configuration

Narrow $p_{//}({}^{23}\text{O} \rightarrow {}^{22}\text{O})$ distribution \rightarrow halo-like tail of the w. f.
(**s**-wave neutron)

Experiments of ${}^{22}\text{O}$ γ -spectroscopy (*P. Thirolf et al., PLB 485(2000)*, *O. Sorlin et al., NPA 685(2001)186c*) \rightarrow **2^+** state in ${}^{22}\text{O}$ at **3.2 MeV** \rightarrow **N=14 sub-shell closure**

d-wave g.s. configuration

$p_{//}({}^{23}\text{O} \rightarrow {}^{21}\text{O})$ distribution \rightarrow $(s_{1/2})^2 d_{5/2}^{-1}$ configuration (in conflict with shell model prediction)

Large **${}^{23}\text{O}$ interaction cross section** \rightarrow modification of the ${}^{22}\text{O}$ core structure (*R. Kanungo et al., PRL 88(2002)142502*)

Theory :

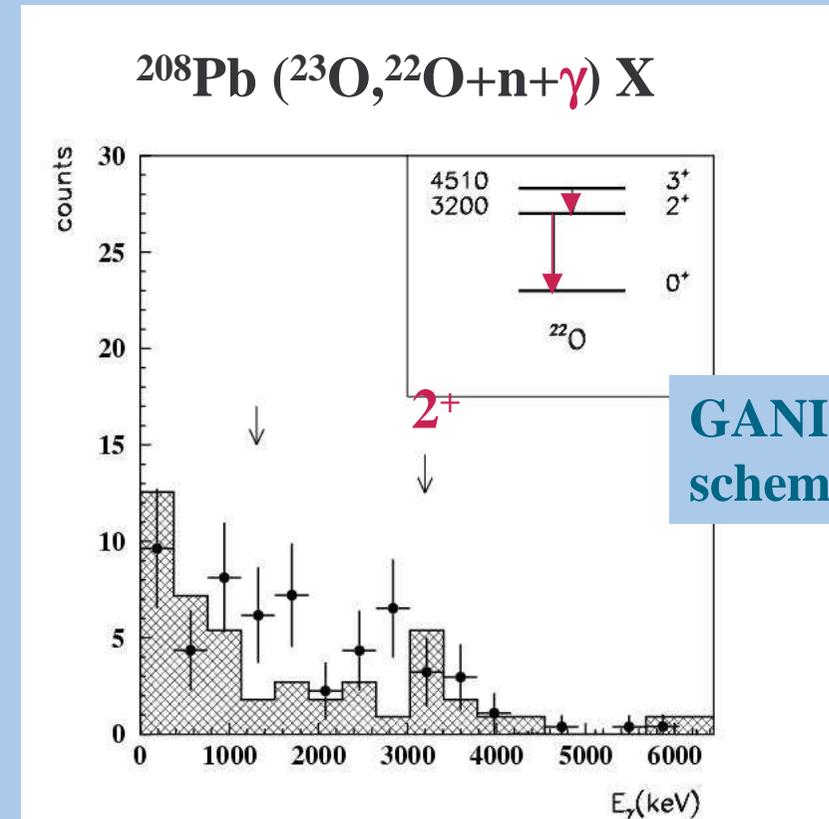
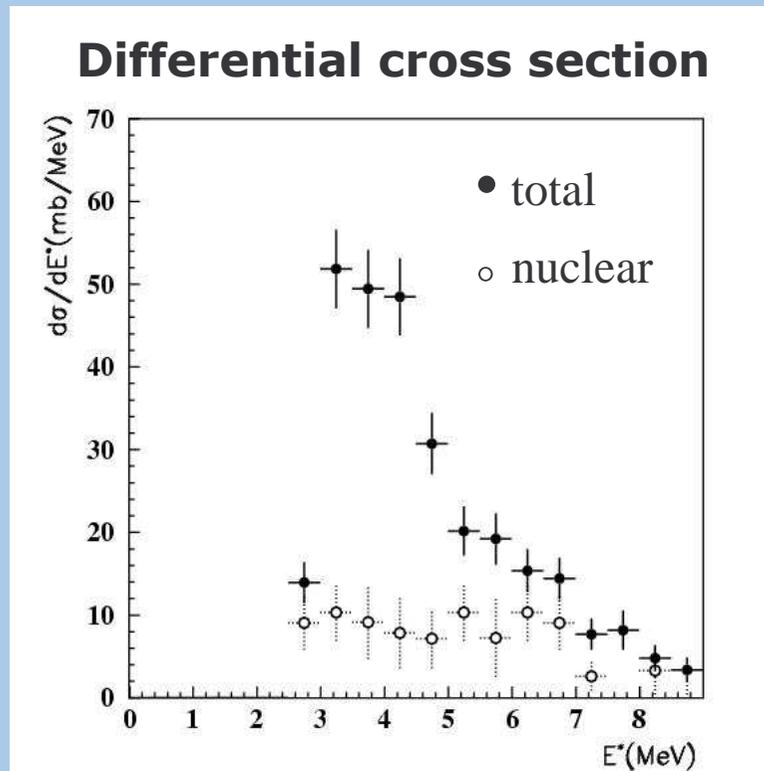
Prediction of shell closure at N=16 rather than at N=20 for O isotopes

(*T. Otsuka, present meeting*)

Coulomb breakup of ^{23}O

At $E_{\text{inc}} = 422 \text{ MeV/u}$

| | σ_{tot} (mb) | σ_{nucl} (mb) | $\sigma_{\text{e.m.}}$ (mb) |
|-----------|----------------------------|-----------------------------|-----------------------------|
| Pb | 131±8 | 37±6 | 94±10 |
| C | 12±2 | 12±2 | - |



At most 30% ^{22}O in excited states

According to the γ -analysis, the quantum numbers of the valence neutron **uniquely** determine spin and parity of $^{23}\text{O}_{\text{g.s.}}$

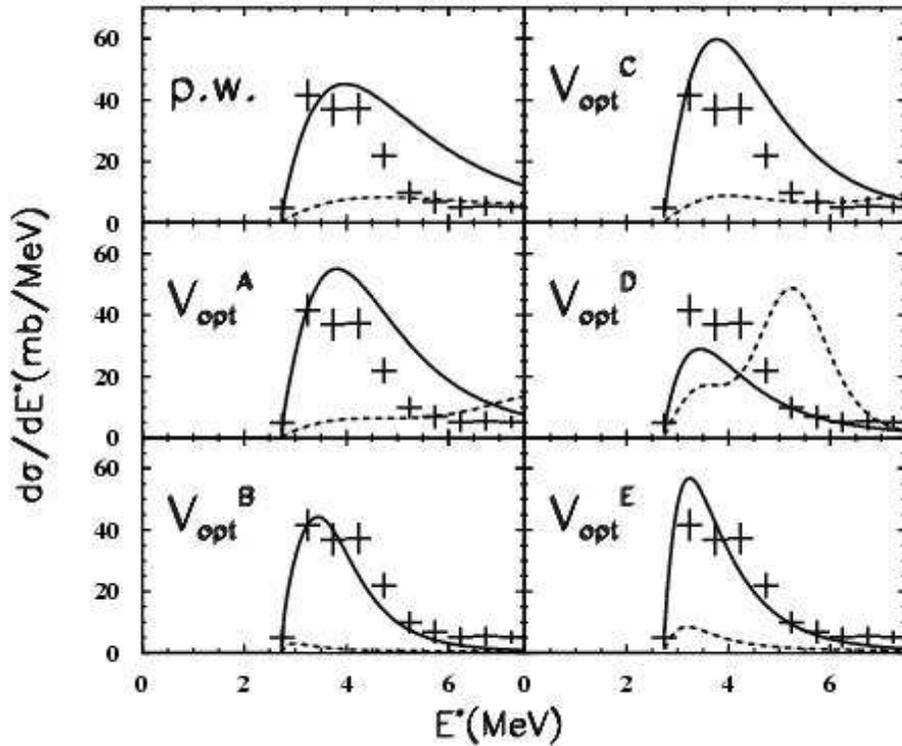
Comparison with models

Differential e. m. cross section

$$\sigma_{\text{e.m.}} = 94 \pm 10 \text{ mb}$$

— $|2s_{1/2} \nu \otimes {}^{22}\text{O}(0^+) \rangle, C^2S=1$

..... $|1d_{5/2} \nu \otimes {}^{22}\text{O}(0^+) \rangle, C^2S=1$



Using distorted waves :

| | | V_0 (MeV) | $\sigma^{\text{th}}(0^+)\text{mb}$ | C^2S |
|---------------------|----------|-----------------|------------------------------------|-----------------------------------|
| S.Type1 calc. | A | 42 | 149 | 0.52 ± 0.10 |
| | B | 55 | 73 | 1.05 ± 0.21 |
| Perey & Perey | C | 49 [§] | 154 | 0.50 ± 0.10 |
| | D | 47 [§] | 63 | 1.22 ± 0.24 |
| Chadwick & Young | E | 49 | 89 | 0.87 ± 0.17 |

§ values quoted for $E_n=0$

V_0 depth of the real part
of optical potentials

corrected for ${}^{22}\text{O}$ core excited states

Proposed $J^\pi=1/2^+$ for ${}^{23}\text{O}_{\text{g.s.}} \rightarrow (2s_{1/2})^2 1d_{5/2}^{-1}$ configuration not supported

Summary

Spectroscopic results for $^{23}\text{O}_{\text{g.s.}}$ via exclusive measurements of one-neutron e. m. breakup :

∅ Contribution of core excited states at most 30%

∅ Proposed $\mathbf{J^\pi = 1/2^+}$ for $^{23}\text{O}_{\text{g.s.}}$ → $(2s_{1/2})^2 1d_{5/2}^{-1}$ configuration not supported

$C^2S \sim 1$ for $\mathbf{|2s_{1/2} \nu \otimes ^{22}\text{O}(0^+) \rangle}$ → support proposed N=14,16 shell closure for O isotopes

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