

Plasma oscillations

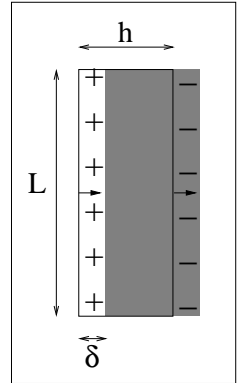
A metallic slab has a square base of side L and a thickness $h \ll L$. The electron and ion density in the slab are $n_e = n$ and $n_i = n_e/Z$, where Z is the ion charge.

Through the appliance of an external field, all the electrons are displaced by $\delta < h$ in the direction perpendicular to the slab surface. We assume that the densities of both electrons and ions are constant, that ions are immobile and that edge effects are negligible.

- a) Find the electrostatic field generated by the displacement of electrons.
- b) Find the electrostatic energy.

The external force is removed, and the “electron slab” starts to oscillate around its equilibrium position.

- c) Find the oscillation frequency, in the small displacement limit ($\delta \ll h$).



Solution

a) Due to the electron rigid displacement, the charge density has the spatial distribution

$$\rho(x) = \begin{cases} +en & (0 < x < \delta), \\ 0 & (\delta < x < h, x < 0, x > h + \delta), \\ -en & (h < x < h + \delta). \end{cases}$$

From $\rho(x)$ we obtain the electrostatic field \mathbf{E} by integrating the equation $\nabla \cdot \mathbf{E} = \partial_x E_x = \rho/\epsilon_0$:

$$E_x(x) = \frac{en}{\epsilon_0} \begin{cases} x & (0 < x < \delta), \\ \delta & (\delta < x < h), \\ h + \delta - x & (h < x < h + \delta), \\ 0 & (x < 0, x > h + \delta). \end{cases}$$

b) To find the electrostatic energy we integrate the “energy density” $u = (\epsilon_0/2)E_x^2$ over the whole volume:

$$\begin{aligned} U_{es} &= \int \frac{\epsilon_0}{2} E_x^2 dV = \frac{\epsilon_0}{2} L^2 \int_0^{h+\delta} E_x^2 dx \\ &= \frac{\epsilon_0}{2} L^2 \left(\frac{en}{\epsilon_0} \right)^2 \left(\int_0^\delta x^2 dx + \int_\delta^h \delta^2 dx + \int_h^{h+\delta} (h + \delta - x)^2 dx \right) \\ &= \frac{(enL)^2}{2\epsilon_0} \left[\frac{\delta^3}{3} + \delta^2(h - \delta) + \frac{\delta^3}{3} \right] = \frac{(enL)^2}{2\epsilon_0} \left(h\delta^2 - \frac{\delta^3}{3} \right). \end{aligned}$$

We used $dV = L^2 dx$.

c) In the limit $\delta \ll h$ we have $U_{es} \simeq (enL)^2 h \delta^2 / 2\epsilon_0$, which is the potential energy of an harmonic oscillator. Note that the expression of U_{es} remains the same also when the electrons move in the $x < 0$ region.

The force on the electron slab is

$$F = -\frac{\partial U_{es}}{\partial \delta} = -\frac{(enL)^2 h \delta}{\epsilon_0}.$$

The equation of motion is

$$M \ddot{\delta} = F \equiv -M \omega^2 \delta,$$

where $M = m_e n L^2 h$ is the total mass of the electron slab. We thus find

$$\omega^2 = \frac{ne^2}{\epsilon_0 m_e} \equiv \omega_p^2,$$

where the *plasma frequency* ω_p is a typical parameter of any conductor, which depends only on the density of (free) electrons.