# Fundamental Issues in Fast Ignition Physics: from Relativistic Electron Generation to Proton Driven Ignition

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## Introduction

A key element of **Fast Ignition** of ICF targets is the generation of a **Petawatt power beam** of **"fast" relativistic electrons** and the **transport** of the latter into a dense plasma or a solid target. The electron beam may either **drive ignition directly** or drive the acceleration of a **proton beam** which is in turn used to ignite. Both ignition scenarios involve a number of physical processes which are widely unexplored and challenging for theory and simulation.

We present theoretical and numerical investigations of several fundamental issues, including:

- "filamentation" instabilities of fast electrons in the "coronal" region
- transport modeling in solid-density matter
- requirements for proton-driven ignition.

#### Scenario of the "traditional" approach to Fast Ignition

[M. Tabak et al, 1994; M. H. Key, Nature 412, 775 (2002)]

The success of the scheme relies on the efficiency of several processes involving generation and transport of fast electrons (FE):

- ⇒ laser energy conversion into FE at the interaction surface (dominated by collision-less processes)
- ⇒ transport of FE into the target [affected by either collisionless processes (microinstabilities, self-generated fields) and by collisions, depending on local density and temperature]
  - efficient energy deposition in the core by FE



#### Electron filaments in the corona: bulk or surface effect?

PIC simulations show that fast electrons form **narrow current filaments** as they enter the overdense plasma [2D PIC simulations by H. Ruhl: see Mulser et al, Las. Phys. **10**, 231 (2000)].



The transverse current filamentation (CF) or "Weibel" instability, driven by the repulsion between the "fast" and the background neutralization currents has been often invoked to explain the origin of filaments. However, coupling to longitudinal modes as well as transient and inhomogenity effects should be considered in this context.

An alternative picture is based on the **corrugation of the target surface** which may lead to a local **geometrical collimation** ("funnel effect") of the fast electrons [Macchi, Ruhl & Cornolti, Las. Part. Beams **18**, 375 (2000)]

#### 3D fluid simulation of the "Weibel" current filamentation instability

**Three-dimensional** (3D) relativistic, collisionless fluid simulations of the CF instability have been performed. Initially **two counterstreaming, homogeneous electron beams** with zero net current, i.e.  $n_1v_1 = -n_2v_2$ , and **highly asymmetrical** ( $|v_1| \gg |v_2|$ ) are assumed.

3D Magnetic structures are formed with typical length scales of a few  $d_e = c/\omega_p$  both in the direction of the beams and in the perpendicular plane. No extended filaments in beam direction are observed. These results agree with 3D theory [F. Califano et al., PRE 58, 7837 (1998)]



Simulation example:  $n_1/n_2 = 9$ ,  $v_1 = 0.95c$ ,  $v_2 = -0.10556c$ . If  $n_0 = n_1 + n_2 = 10^{22} \text{ cm}^{-3}$  (typical of PIC simulations), the energy flux of fast electrons (1 MeV energy) is  $4.6 \times 10^{18} \text{ W cm}^{-2}$ . Figure shows isosurfaces of  $A_z$  (vector potential component along beam direction), which is representative of the magnetic field structure because  $B_z \ll (B_x, B_y)$  is found.

#### Imprint effect by surface parametric instabilities

PIC simulations [1] and theory [2] show that pairs of nonlinear **electron surface waves** (ESWs) can be parametrically excited by the laser pulse. For normal incidence and *s*-polarization, the ESWs are driven by the  $\mathbf{v} \times \mathbf{B}$  force and form a **standing wave pattern** on the target surface.

Simulations show that fast electrons are then generated predominantly at **maxima of the standing wave**. This gives an **"imprint" for electron filaments** whose size and spacing is proportional to the laser wavelength as observed in other simulations.



#### Fast electron transport in solid-density matter

Typical transport regime in fast ignitor and laser-solid interaction experiments:

fast  $e^-$ :  $n_f \approx 10^{18} - 10^{20} {\rm cm}^{-3}$ ,  $\varepsilon_f \approx 10^5 - 10^7 {\rm eV}$ ,  $\tau_c > 10^{-12} {\rm s}$ 

slow  $e^-$ :  $n_b \approx 10^{23}$  cm<sup>-3</sup>,  $\tau_c < 10^{-15}$  fs (metal) produce a neutralizing "return" current  $\mathbf{j}_b$ :

$$\mathbf{j}_b \simeq \frac{n_b e^2}{m_e \tau_c} \mathbf{E} = \sigma \mathbf{E}$$



$$|{f j}|_f pprox 10^{11} - 10^{12} \ {
m A \ cm^{-2}}, I pprox 10^5 - 10^6 \ {
m A}$$

 $\rightarrow$  self-generated E and B fields are important

In low  $\sigma$  materials,  $\mathbf{j}_b$  may not compensate well  $\mathbf{j}_f$ 

ightarrow generation of E fields

makes fast electron transport  $\sigma$ -dependent.

#### Experimental study of transport inhibition by self-generated fields

A first evidence of **material-dependent penetration** of fast electrons has been obtained by comparing results (from **K-shell emission** diagnostics) in **AI targets** and **AI-coated plastic targets**, so that **the spectrum of fast electrons is identical** in the two cases [1].



From K-shell data it is inferred that the **penetration range is much shorter in plastic than in AI** and **cannot be explained by collisional stopping** of fast electrons alone.

[1] F. Pisani, A. Bernardinello, D. Batani et al, Phys. Rev. E 62, R5927 (2000)

#### Fast electron penetration in low-density foam targets

The K- $\alpha$  yield  $I_{K\alpha}$  has been measured as a function of the density  $\rho$  of **plastic foam targets** with **constant mass thickness**  $\rho d$ . From  $I_{K\alpha}$  the **penetration range** R of fast electrons is obtained.

The deviations from the **purely collisional be**havior ( $R \sim \rho^{-1}$ ,  $I_{K\alpha} = \text{cost.}$ ) indicate the effect of **self-generated fields** on fast electron transport [2].



[2] D. Batani, A. Antonicci, F. Pisani et al, Phys. Rev. E 65 066404 (2002)

#### Transport simulations with the PENELOPE code

The 3D Monte-Carlo electron-photon transport code PENELOPE [Beró et al, NIM B **100**, 31 (1995)] has been adapted by including **field generation** and **K**- $\alpha$  **diagnostics** [Antonicci et al. (2001)] and used to simulate the experiment by Pisani et al. Local **charge neutrality** is assumed.



The **penetration range** of fast electrons is affected by **self-generated fields**; the effect is **stronger in CH** (low conductivity).

The comparison of simulations for Al targets with (left) and without (right) fields included show self-collimation of the fast electron beam in the first case.

#### Simulations of transport inhibition by electrostatic fields

To address the importance of **electrostatic (ES) fields** in low- $\sigma$  materials, a new hybrid code with **charge separation effects included** is under development. Preliminary results in 1D are reported for an Al isothermal plasma ( $T_e = 100 \text{ eV}$ ), fast electrons with 200 keV energy, beam intensity  $3.6 \times 10^{17} \text{ W cm}^{-2}$ . Profiles of fast electron density and ES field and the (x- $p_x$ ) phase space plot are shown at t = 100 fs.



Surface charges generate a strong back-holding ES "capacitor" field (thin line). The density profile (thick) is very different from that obtained from Bell's formula (dotted) [Bell et al., PPCF **39**, 653 (1997)].

Phase space plots show that most of the electrons are reflected back by the ES field, with only a small number penetrating deep into the target. Broadening in momentum space is due to Coulomb collisions.

#### Proton acceleration by fast electrons in solid targets

The electric field  $E_b \approx -j_f/\sigma$  generated into the target bulk slows down fast electrons and accelerates ions, in particular hydrogen ions which are present also in metallic targets as impurities.

In addition, hydrogen ions can be accelerated efficiently at the target back surface by fast electrons escaping the target and forming a "Debye" sheath with density  $n_d$ , with an electric field

$$E_s \approx 2\sqrt{\pi m_e c^2 (\gamma_f - 1) n_d}$$

The energy of protons from the back surface is about one half of the fast electron energy (e.g. a few MeV) with a laser energy conversion efficiency into protons up to a few percent.

These features make these laser-accelerated protons **extremely interesting for applications**. In the following a study of their application in the **fast ignition of ICF targets** is investigated.

#### Study of fast ignition by laser-generated proton beams



The requirements for fast ignition by protons are studied by means of **analytical modeling** and **2D radiation-hydro-nuclear simulations** to obtain the **minimum proton beam energy**  $E_{ig}$  for ignition as a function of **fuel density**  $\rho$ , **proton beam temperature**  $T_p$  and **source-to-fuel distance** d.

For homogeneous precompressed fuel spheres,

d=1-4mm,  $T_p\simeq 5$  MeV,

 $E_{ign} = 90 (d/1 \text{mm})^{0.7} / (\rho / [100 \text{g/cm}^3])^{1.3} \text{kJ}.$ 

 $E_{ign}$  is larger than earlier estimates due to effects related to beam velocity dispersion and reduced plasma stopping power at high temperatures.

S. Atzeni, M. Temporal and J. J. Honrubia, Nucl. Fusion 42, L1 (2002)

### 2D simulations of ignition by protons and fuel burn



M. Temporal, J. J. Honrubia and S. Atzeni, Phys. Plasmas 9, 3098 (2002).

# Perspectives for future work

- Understanding current filamentation: study transient and inhomogeneity effects in Weibel-like instabilities, go deeper into the imprint effect of surface waves.
- Improving models of fast electron transport: account for electrostatic effects and conductivity changes due to ionization.
- Integration of transport and fusion codes: evaluating the energy transport flux by electrons and protons to give as an input to burn simulations.
- **Different schemes of fast ignition**: simulation of alternative target design, test of "coronal" ignition [S. Hain and P. Mulser, PRL **86**, 1015 (2001)].

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