

High intensity laser-grating interactions: a step towards relativistic plasmonics?

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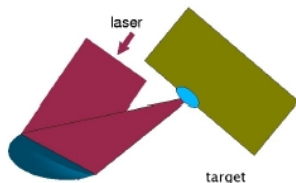
Ultrashort laser-solid interactions: what for?

Ultrashort pulses: < 100 fs duration,
up to $\sim 10^{20}$ W cm $^{-2}$ intensity

Applications:

- ▶ laser-driven particle sources (electrons, **ions**)
- ▶ high harmonic generation for coherent ultrashort X-ray emission
- ▶ isochoric heating of matter

efficient laser-target coupling is a key issue for these processes



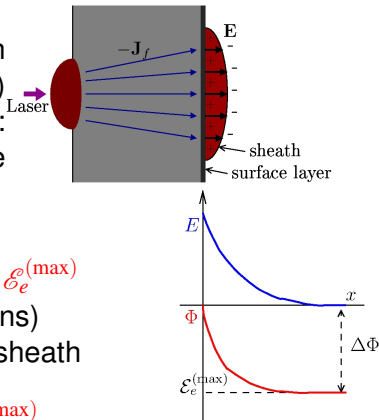
Example: sheath acceleration of protons

Target Normal Sheath Acceleration (TNSA) is driven by “fast” ($\mathcal{E}_e \sim \text{MeV}$) electrons generated in thin targets: protons from surface contaminants are accelerated in the rear sheath

TNSA picture “for dummies”:

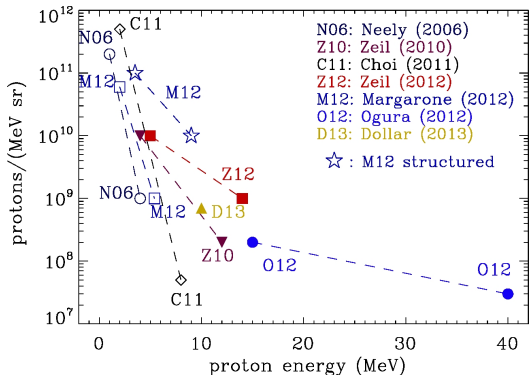
Potential drop for static sheath $e\Delta\Phi = \mathcal{E}_e^{(\text{max})}$
(sheath potential must confine electrons)
Energy gained by “test” proton in the sheath

$$\mathcal{E}_p = e\Delta\Phi = \mathcal{E}_e^{(\text{max})}$$



Efficient electron heating is key to TNSA

Recent TNSA data with “table-top” lasers



Intensity range:

$$I = (1 - 5) \times 10^{19} \text{ W cm}^{-2}$$

(empty symbols)

$$I = (0.8 - 2) \times 10^{21} \text{ W cm}^{-2}$$

(filled symbols)

Pulse duration range:

$$\tau = 25 - 40 \text{ fs}$$

Target thickness range:

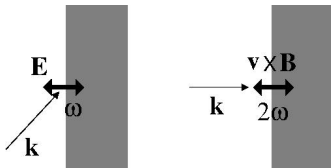
$$\ell = 0.05 - 4.0 \mu\text{m}$$

Need to find strategies to increase absorption and fast electron energy to boost TNSA

A. Macchi et al, Plasma Phys. Contr. Fus. **55**, 124020 (2013)

Fast electron generation: simple (rough) picture

The Lorentz force of the laser wave (amplitude E_L , frequency ω) drives periodic “push-pull” of electrons across the density gradient

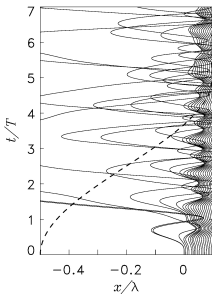


Electrons perform “half-oscillations” on the vacuum side and re-enter in the plasma (where the EM field is screened) and are “absorbed” keeping a net momentum

$$p_e \sim p_{\text{osc}} \sim eE_L/\omega \equiv m_e c a_0$$

$$a_0 = \left(\frac{eE_L}{m_e c \omega} \right) > 1 \longrightarrow \text{relativistic electrons}$$

A. Macchi, *A Superintense Laser-Plasma Interaction Primer* (Springer, 2013)



Looking for resonant coupling

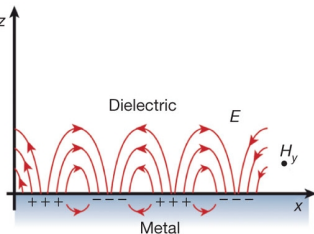
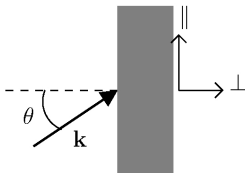
Idea: enhancement of the surface field and of absorption by exciting a **normal mode** of the target plasma

Resonant coupling requires matching of the phase $\varphi = \mathbf{k}_{\parallel} \cdot \mathbf{r} - \omega t$ between the laser and the resonant mode characterized by (ω_m, \mathbf{k}_m) :

$$\omega \doteq \omega_m \quad k_{\parallel} = k \cos \theta \doteq (k_m)_{\parallel}$$

Normal modes of step boundary metal/plasma: **surface waves**
figure from:

O.Benson, Nature **480**, 193 (2011)



Surface wave coupling: the matching problem

Plasma-vacuum interface

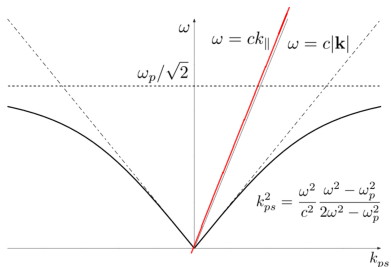
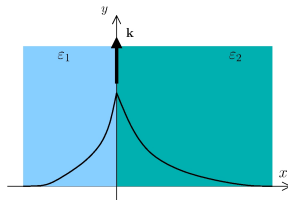
$$\varepsilon_1 = 1 \quad \varepsilon_2 = 1 - \frac{\omega_p^2}{\omega^2} = 1 - \frac{n_e}{n_c}$$

Dispersion relation $\omega = \omega(k)$

$$k = \frac{\omega}{c} \left(\frac{\omega_p^2 - \omega^2}{\omega_p^2 - 2\omega^2} \right)^{1/2}$$

$$\omega < \omega_p/\sqrt{2} \quad v_{ph} < c$$

No matching with $\omega = ck_{\parallel}$



Surface wave matching in periodic structures

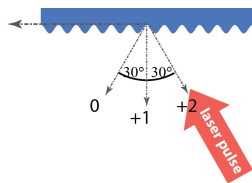
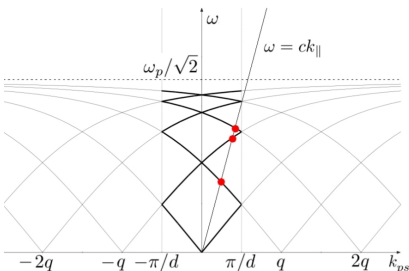
In a **spatially periodic** medium with period d , folding of $\omega_{SW}(k)$ in the Brillouin zone $|k| < \pi/d$ (Floquet-Bloch theorem) allows phase matching

Figure: M.Lupetti, M.Sc. Thesis, 2011

Resonant coupling with EM wave is possible in a **grating** at an angle of incidence

$$\sin \theta_{\text{res}} + \lambda/d = \left(\frac{1 - \omega_p^2/\omega^2}{2 - \omega_p^2/\omega^2} \right)^{1/2}$$

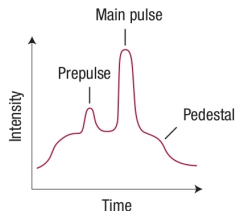
(provided $\omega_{SW}(k)$ does not change much)



Earlier work on laser-grating interaction

Coupling of laser field with surface waves is a “building block” of **plasmonics**: the art of light concentration and manipulation on the sub-wavelength scale

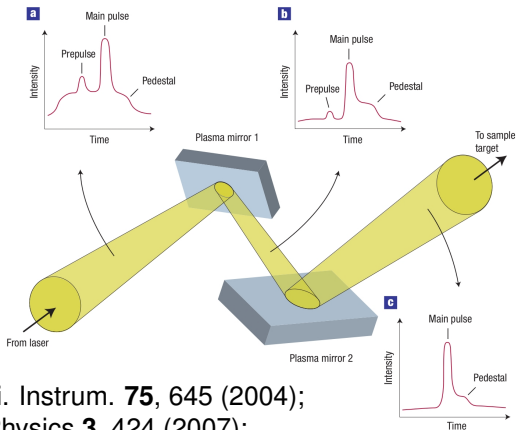
Experiments performed since 1990's at high intensity have been limited by **prepulse** effects causing destruction of shallow grating structure before short pulse interaction



A general issue: do surface waves exist in a **high-field, nonlinear** and **relativistic** regime?

Need for ultraclean pulses: plasma mirrors

Plasma mirrors yielding $\sim 10^{12}$ pulse-to-prepulse contrast allow to preserve target structuring until the short pulse interaction



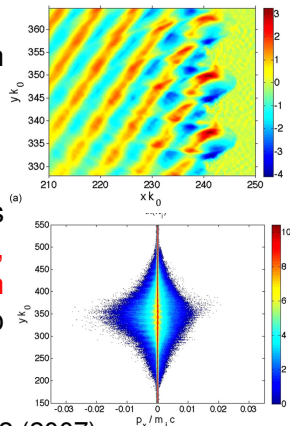
B. Dromey et al, Rev. Sci. Instrum. **75**, 645 (2004);
C. Thaury et al, Nature Physics **3**, 424 (2007);
figure from P. Gibbon, *ibid.*, 369.

Surface wave coupling in the “relativistic” regime

No theory for SW in the nonlinear, high field, **relativistic** electrons regime

$$a_0 > 1 \longrightarrow I\lambda^2 > 1.4 \times 10^{18} \text{ Wcm}^{-2}\mu\text{m}^2$$

However, particle-in-cell simulations show **enhancement of absorption, electron heating and ion acceleration** for laser-grating interactions up to $I\lambda^2 \sim 10^{20} \text{ Wcm}^{-2}\mu\text{m}^2$



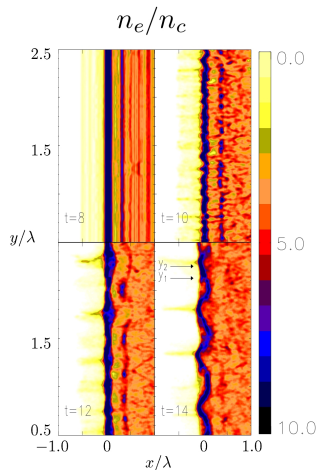
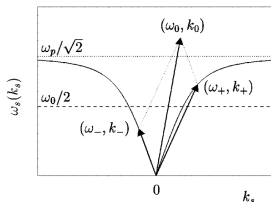
M. Raynaud et al, Phys. Plasmas **14**, 092702 (2007);

A. Bigongiari et al, *ibid.* **18**, 102701 (2011); **20**, 052701 (2013)

Another early evidence for relativistic SW

Laser-driven periodic surface oscillations decay in two surface waves via a period-doubling process; similarity to **Faraday Waves (Ripples)** in a liquid

No grating necessary for nonlinear phase matching

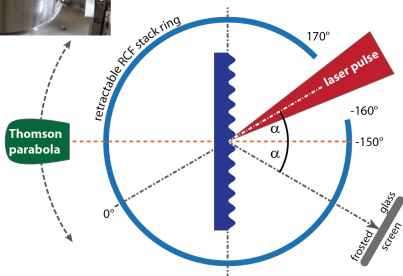
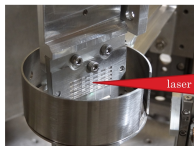


A. Macchi et al, Phys. Rev. Lett. **87**, 205004 (2001);
Phys. Plasmas **9**, 1704 (2002)

Experimental set-up



LaserLAB experiment at SLIC, CEA Saclay
laser UHI, 28 fs, $5 \times 10^{19} \text{ W cm}^{-2}$, **contrast** $\sim 10^{12}$



Grating:

- $d = 2\lambda \rightarrow \theta_{\text{res}} = 30^\circ$
- depth $\delta = 0.3 - 0.5 \mu\text{m}$

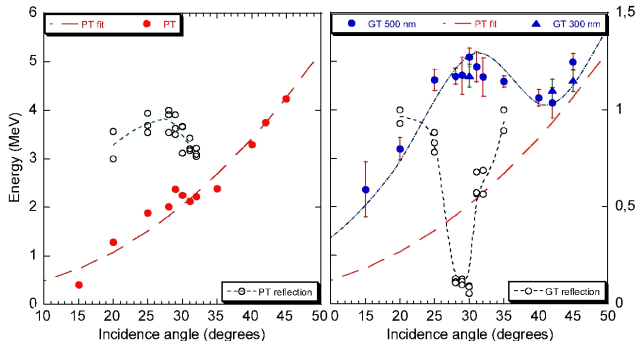
Diagnostics:

- Thomson Parabola for proton detection
- Radio-Chromic Film (RCF) "ring" for radiation emission at any angle
- Reflected light

Plane target vs grating

Proton energy cut-off \mathcal{E}_{\max} and reflected light vs incidence angle:

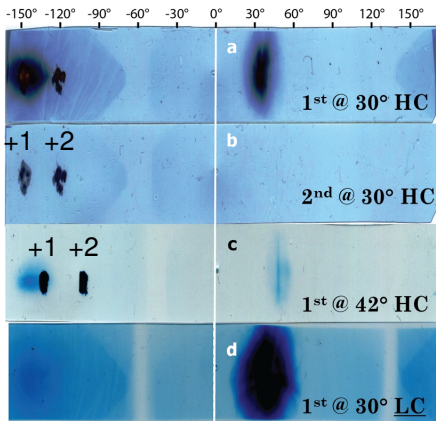
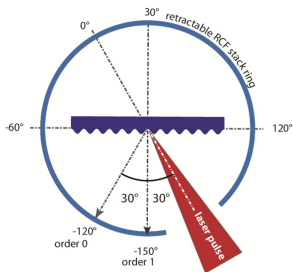
- broad maximum (minimum) around $\theta_{\text{res}} = 30^\circ$
- $\sim 2.5X$ enhancement in \mathcal{E}_{\max} at θ_{res} , ~ 2 at small angles



PT: plane target
GT: grating target
● \mathcal{E}_{\max} , GT, $\delta = 0.5 \mu\text{m}$
▲ \mathcal{E}_{\max} , GT, $\delta = 0.3 \mu\text{m}$
● \mathcal{E}_{\max} , PT
○ reflected light

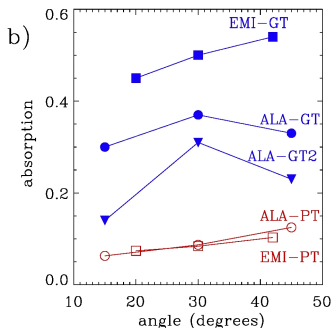
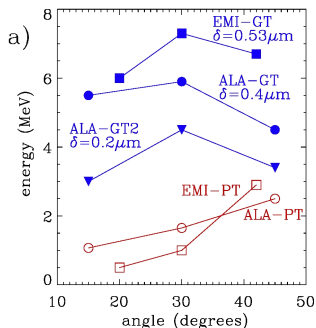
Grating signatures on RCF

Diffraction orders produce angle-dependent “burn spots” for High Contrast (HC), not observed with Low Contrast (LC)



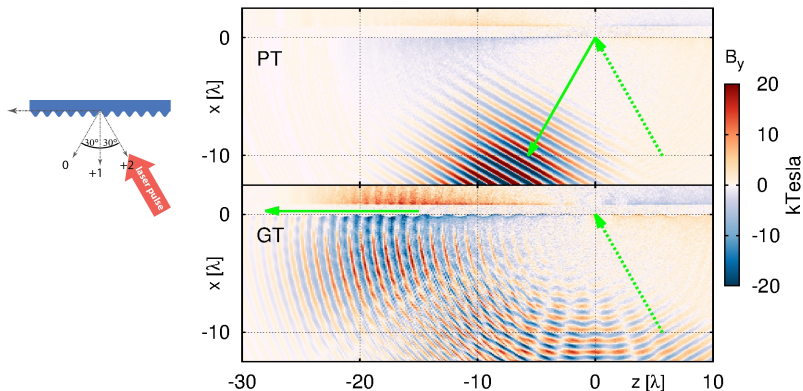
Comparison with PIC simulations

Two simulations campaigns with Particle-In-Cell codes **EMI2D** (CPHT, École Polytechnique) and **ALADYN** (Italy) fairly reproduce experimental trend (2D simulations, different set-up for the two codes)



Surface wave in simulations

Snapshots of EM fields show **localized wave** propagating along the surface at resonant angle of incidence (plus reflection at various diffraction orders)

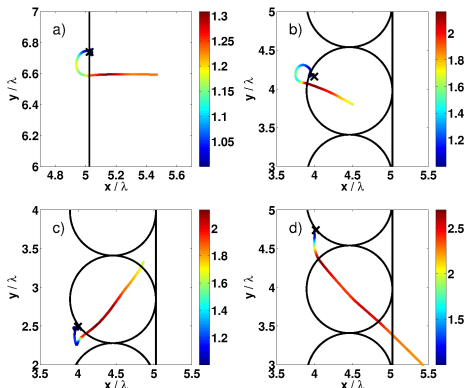


Enhanced electron heating out of resonance

Stochastic heating at a modulated interface is more efficient than in plane targets: electrons have more “re-entering trajectories” available

Effect observed in **microsphere-covered targets**

(PIC sim. by O. Klimo et al)

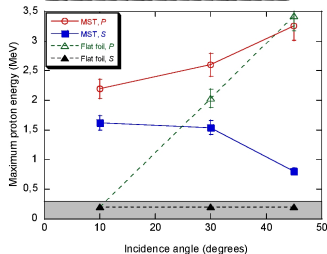
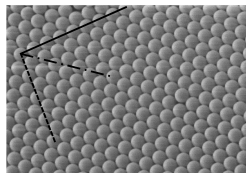


V. Floquet et al, J. Appl. Phys. **114**, 083305 (2013)

Experiment on microsphere-covered targets

Measurements taken in same campaign at SLIC

Some enhancement of proton energy observed only at small angles of incidence
(compare with Margarone et al, Phys. Rev. Lett. **109**, 234801 (2012))



V. Floquet et al, J. Appl. Phys. **114**, 083305 (2013)

Conclusions: open issues and future work - I

From the point of view of using grating targets for “enhanced” proton acceleration:

- ▶ energy increase at resonant angle does not exceed maximum energy in plane targets at grazing incidence: need to test different (larger) angles
- ▶ grating structure (and also microsphere covering . . .) did not work with very thin foils: need to use smarter materials to guarantee target integrity
- ▶ efficiency of grating effect for longer pulses (delivering more energy) uncertain because of hydrodynamics on ~ 100 fs temporal scale
- ▶ is it feasible to embed gratings in a complex target design (e.g. microcones, . . .)?

Conclusions: open issues and future work - II

- ▶ From a point of view of “general interest”:
 - ▶ First experimental indication of **surface waves** in the regime of **relativistic** electrons
- possible next steps:
- ▶ work out theory of nonlinear, relativistic SW
 - ▶ investigate detailed mechanism of energy absorption and electron acceleration by SW
 - ▶ design **plasmonics** applications in the **high fields** regime, exploiting the SW resonance

Main reference

PRL **111**, 185001 (2013)

PHYSICAL REVIEW LETTERS

week ending
1 NOVEMBER 2013

Evidence of Resonant Surface-Wave Excitation in the Relativistic Regime through Measurements of Proton Acceleration from Grating Targets

T. Ceccotti,^{1,*} V. Floquet,¹ A. Sgattoni,^{2,3} A. Bigongiari,⁴ O. Klimo,^{5,6} M. Raynaud,⁷ C. Riconda,⁴ A. Heron,⁸ F. Baffigi,² L. Labate,² L. A. Gizzi,² L. Vassura,^{9,10} J. Fuchs,⁹ M. Passoni,³ M. Květon,⁵ F. Novotny,⁵ M. Possolt,⁵ J. Prokūpek,^{5,6} J. Proška,⁵ J. Pšikal,^{5,6} L. Štolcová,^{5,6} A. Velyhan,⁶ M. Bougeard,¹ P. D'Oliveira,¹ O. Tcherbakoff,¹ F. Réau,¹ P. Martin,¹ and A. Macchi^{2,11,†}

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T. Ceccotti et al, Phys. Rev. Lett. **111**, 185001 (2013)



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- ▶ ...

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