Radiation Reaction and Laser Polarization Effects on Ultraintense Laser Acceleration of Thin Foils

Andrea Macchi

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Coworkers

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Motivations

► Radiation Pressure Dominant Acceleration (RPDA) of thin $(\ell \lesssim 1 \ \mu m)$ solid foils by ultra-relativistic laser pulses $(a_0 > (m_p/m_e)/2\pi, I\lambda^2 > 10^{23} \text{ W cm}^{-2}\mu m^2)$ has been suggested as a route to "unlimited" acceleration towards the GeV/nucleon limit (relativistic ions)

[Esirkepov 04, Bulanov 10]

 The simple, idealized "Light Sail" model predicts 100% efficiency and monoenergeticity in the asymptotic, relativistic limit

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 $E_{\max} \simeq m_p c^2 \mathscr{F}/2$

 $\mathscr{F} = 2(\rho \ell)^{-1} \int_0^\infty I(t') dt'$

Favorable scaling with laser pulse fluence \mathscr{F} 100% efficiency in the relativistic limit"Perfect" monoenergeticity for "rigid", coherent motion of the foilLimits: "slow" energy gain, foil transparency



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Favorable scaling with laser pulse fluence \mathscr{F} Image: Comparison of the relativistic limit100% efficiency in the relativistic limit \mathcal{I} "Perfect" monoenergeticity for "rigid", coherent motion of the foilLimits: "slow" energy gain, foil transparency



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Physics background 2: "unlimited" acceleration

Transverse expansion of the target reduces surface density $\rho \ell$ \Rightarrow "unlimited" acceleration possible at the expense of the number of accelerated io [Bulanov 10]

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- Early numerical demonstration of RPDA [Esirkepov 04] by 3D simulations suggests polarization is inessential
- Unlimited acceleration first demonstrated by 2D simulation and circular polarization (CP) [Bulanov 10]
- Several studies after [Macchi 05] showed that CP enforces RPDA also at low intensities, although more recent work suggests a regime where also Linear Polarization (LP) works well [Qiao 12]
- Need to clarify polarization and 3D effects in "unlimited" RPDA

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- Radiation Friction (or Radiation Reaction RR) effects are expected to play an important role in ultrarelativistic laser-plasma dynamics
- \Rightarrow address role of RR on RPDA
- Effect of RR on ion acceleration (both in thin and thick targets) investigated by several authors [Schlegel 09, Chen 11, Capdessus 12]
- Collective laser-plasma acceleration also of interest as the context for first direct observation of RR [Di Piazza 09, Hadad 10]

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⇒ use RPDA as a test and benchmark case to develop an adequate modeling of RR in laser-plasma interactions

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- "Feasible" 3D Particle-In-Cell simulations are at the limit of present-day computational resources: hard to check numerical convergence
- RPDA simulations further challenging because of high density and long times to reach asymptotic regime
- when including RR the high-energy, low-density tail of the electron distribution needs to be resolved properly
- ⇒ check effects of limited resolution with simulations in lower dimensionality (2D)

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⇒ upgrade the code for more efficiency, use larger supercomputers ... and do what you can

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"Reduced" Landau-Lifshitz equation for electrons

$$\begin{aligned} \frac{d\mathbf{p}}{dt} &= -e\left(\mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B}\right) + \mathbf{F}_{rad} \\ \mathbf{F}_{rad} &= -\left(\frac{2r_c^2}{3}\right) \times \left\{\gamma^2 \left[\left(\mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B}\right)^2 - \left(\frac{\mathbf{v}}{c} \cdot \mathbf{E}\right)^2\right] \frac{\mathbf{v}}{c} \\ &- \left[\left(\mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B}\right) \times \mathbf{B} + \left(\frac{\mathbf{v}}{c} \cdot \mathbf{E}\right) \mathbf{E}\right]\right\} \end{aligned}$$

 \blacktriangleright Spin force and smaller terms in ${\rm I\!F}_{\rm rad}$ are neglected

• Dominant term ($\sim -\gamma^2 \mathbf{v}$) acts as nonlinear friction

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Kinetic equation for electrons and some properties

$$\partial_t f + \nabla_{\mathbf{r}} \cdot (\mathbf{v}f) + \nabla_{\mathbf{p}} \cdot (\mathbf{F}f) = 0, \qquad f = f(\mathbf{r}, \mathbf{p}, t)$$

$$\mathbf{F} = -e(\mathbf{E} + \mathbf{v} imes \mathbf{B}/c) + \mathbf{F}_{\mathsf{rad}}, \qquad
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$$\frac{d}{dt}\int f\ln f d^3\mathbf{p} d^3\mathbf{q} = \int f d^3\mathbf{p} d^3\mathbf{q} \nabla_{\mathbf{p}} \cdot (\mathbf{F}_{\mathsf{rad}}) < 0$$

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- Simple, modular and computationally cheap algorithm for inclusion in standard PIC codes has been developed [Tamburini et al. NJP 10, 123005 (2010)]
- Numerical benchmark of single particle motion with exact solution in a plane wave [Di Piazza 08]
- High-frequency radiation is assumed consistently to be incoherent and to escape from the plasma; it appears as energy dissipation
- Dominant term in reduced LL force identical to other models and approximations [Schlegel 09, Chen 11]
- Classical approach to RR estimated to be valid up to intensities < 10²⁴ W cm⁻² (open issue)

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Radiation Reaction misunderstanding

Quoting an anonymous referee:

- "The standard PIC algorithm does include the RR when it appears as N² effect. For example, the RPDA of plasma is the result of radiation reaction acting on the plasma as a whole [...] The RPDA itself is an example of the RR effect acting on the plasma as a whole. [...] at higher resolution the RR effects become better approximated by the standard PIC algorithm"
- Evidence of confusion between Radiation Reaction and Radiation Pressure, and of ignorance of the physical meaning of Radiation Reaction (reading of textbooks of classical electrodynamics is recommended)

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- ► Laser pulse: $(9T) \times (10\lambda)^2$ (FWHM) $[T = \lambda/c]$ $\sin^2 \times$ Gaussian shape, $a_0 = 280$ (198) for LP (CP), $\lambda = 0.8 \ \mu$ m
- ► Plasma: $\ell = 1\lambda$, $n_0 = 64n_c$, Z = A = 1Note: $a_0 \simeq \zeta = \pi (n_e/n_c)(\ell/\lambda)$
- ▶ Numerical: $1320 \times 896 \times 896$ grid, $\Delta x = \Delta y = \Delta z = \lambda/44$, $\Delta t = T/80 = \lambda/80c$, 216 particles per cell (for both *e* and *p*), 1.526×10^{10} in total

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Runs performed on 1024 processors (1.7 GBytes each) of IBM-SP6 at CINECA (Italy)

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▶ Numerical: $1320 \times 896 \times 896$ grid, $\Delta x = \Delta y = \Delta z = \lambda/44$, $\Delta t = T/80 = \lambda/80c$, 216 particles per cell (for both *e* and *p*), 1.526×10^{10} in total

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Runs performed on 1024 processors (1.7 GBytes each) of IBM-SP6 at CINECA (Italy)

Andrea Macchi

- ► Laser pulse: $(9T) \times (10\lambda)^2$ (FWHM) $[T = \lambda/c]$ $\sin^2 \times \text{Gaussian shape}, a_0 = 280$ (198) for LP (CP), $\lambda = 0.8 \ \mu\text{m}$
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Space-energy distribution: CP, no RR

lons

Electrons

→ E → < E →</p>

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Symmetric, collimated distribution of ions Cut-off energy of ~ 1.6 GeV at t = 20T

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Space-energy distribution: CP, with RR



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Ion distribution unchanged by RR

Cooling of electrons in pulse tail due to radiative losses

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Space-energy distribution: LP, no RR

lons

Electrons

→ E → < E →</p>

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Asymmetric distribution with highest energy ions off-axis Cut-off energy of $\sim 0.9~{\rm GeV}$ much lower than for CP

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Space-energy distribution: LP, with RR

lons

LP RR (b) LP RR (b) 1100 500 t=20T t=20T 4 2 800 300 600 260 400 220 200 300 200 180 100 160 (MeV) (MeV) 10 18 18 16 18 20 80

Electrons

Image: A matrix

<- E> < E>

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Strong cooling of electrons by radiative losses Cut-off energy is *increased* up to ~ 1.1 GeV by RR

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Sections of 3D fields [a1)-d3)] vs 2D simulations [e)-f)]



Focusing of the pulse down to $\sim \lambda^3$ volume for CP [see series -3)] "Wrapping" and focusing effects are weaker in 2D vs 3D [see e)-f)] Breakthrough of the foil occurs for LP [see series -1)-2)]

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Comparison of 3D ion spectra with 2D results (both S and P for LP) for both the same and higher resolution



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► For Circular Polarization, the energy cut-off (corresponding to ions on axis) is *higher* in 3D than in 2D, and also *higher* than 1D "Light Sail" scaling for the same value of $n_0 \ell$

Explanations for "3D increase":

1: more efficient rarefaction by transverse expansion

increase of energy density on axis by pulse self-wrapping

- Circular Polarization optimizes ion acceleration (collimated distribution, negliglible RR effects) with respect to Linear Polarization
- ⇒ Need to push simulations to longer durations to evaluate the ultimate energy gain

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- ► For Circular Polarization, the foil moves coherently at $v_x \sim c$ and wave reflection is small $(R \sim (1 - v_x/c)/(1 + v_x/c))$
- ⇒ electrons move with $v_x \simeq 1$ in a "quasi" propagating wave with $E \simeq B$ ⇒ $\mathbf{F}_{rad} = 0$: weak RR effects
- ► For Linear Polarization, the J × B oscillating term causes v_x ~ -c periodically ("colliding" geometry) maximizing F_{rad}: strong RR effects
- Nevertheless, RR reduces the pulse penetration through the foil so that the cut-off energy of ions is *increased* by RR despite the higher radiative losses

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 RR effects are quite sensitive to numerical effects: the contribution of the highest energy electrons is not well resolved with few particles per cell

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Understanding Radiation Reaction effects

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- Proposal: anomalies in Thomson scattering from laser-plasma accelerated electron bunches
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- PIC simulations may support the modeling of such experiment
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Work in progress in Italy: γ -RESIST experiment

Study of all-optical. laser-plasma source of γ -rays with FLAME laser at INFN laboratories, Frascati (Italy)



Possible test bed for experiments on detecting RR signatures in a context oriented to high-energy, fundamental physics

Image: A matrix

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- Work sponsored by the FIRB-MIUR (Italy) project SULDIS ("Superintense Ultrashort Laser-Driven Ion Sources")

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 Use of supercomputing facilities at CINECA (Italy) sponsored by the ISCRA project TOFUSEX ("TOwards FUII-Scale simulations of laser-plasma EXperiments") award N.HP10A25JKT-2010

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