Surface Plasmon Acceleration Without a Grating

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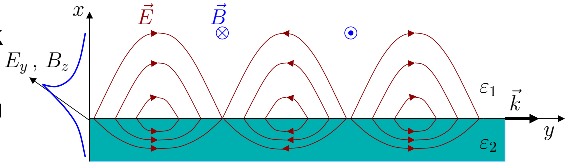
December 15, 2022

Surface Plasmon (Polariton)[1]

SP: a building block of plasmonics (mostly studied in the *linear* regime)



SP excitation $\longrightarrow$ EM field confinement and enhancement
Interface between vacuum and "simple metal" (cold plasma):

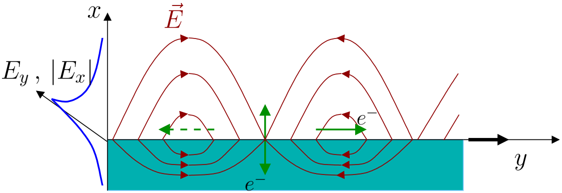
$$\varepsilon_1 = 1 \qquad \varepsilon_2 = 1 - \frac{\omega_p^2}{\omega^2} = 1 - \frac{n_e}{n_c(\omega)} < -1$$

$$k = \frac{\omega}{c}\left(\frac{\omega_p^2 - \omega^2}{\omega_p^2 - 2\omega^2}\right)^{1/2} \qquad \omega < \frac{\omega_p}{\sqrt{2}} \qquad v_{\rm p} = \frac{\omega}{k} < c$$

[1] aka Surface Plasma Wave

Surfin' the Surface Wave

SP can accelerate electrons much like "bulk" plasma waves:
- longitudinal E-component (E_y)
- phase velocity $v_{\mathrm{p}} \lesssim c$ (with $v_{\mathrm{p}} \to c$ when $\omega_p \gg \omega$)



Simple model predicts for maximum energy and emission angle (relative to surface):

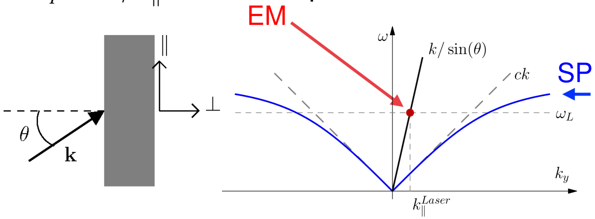
$$\mathcal{E} \simeq m_e c^2 a_{\mathrm{SP}} \frac{\omega_p^2}{\omega^2} \qquad \tan\phi_e = \frac{p_x}{p_y} \simeq \frac{1}{\gamma_{\mathrm{p}}}$$

From: T. Katsouleas, "On the node of a wave", Nature **444** (2006) 688

Exciting Surface Plasmons with Laser Light

Phase matching requirement between SP and plane EM wave ($\omega_{\text{EM}} = |\mathbf{k}_{\text{EM}}|c$):
$\varphi_{\text{EM}} = \varphi_{\text{SP}}$ where $\varphi = \mathbf{k}_{\parallel} \cdot \mathbf{r} - \omega t$
i.e. phase velocities $v_p = \omega / k_{\parallel}$ must be equal $v_{\text{EM}} = v_{\text{SP}}$

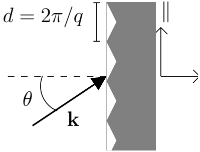
No matching with EM wave at a plane interface:

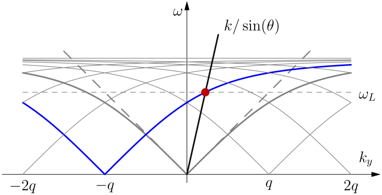


$$v_{\text{EM}} = \frac{c}{\sin\theta} > c \qquad v_{\text{SP}} = c\left(\frac{2 - \omega_p^2/\omega^2}{1 - \omega_p^2/\omega^2}\right)^{1/2} < c$$

Grating Coupling

Periodic grating:
"replica" (*)
of $\omega_{\mathrm{SP}}(k_{\parallel})$
enables matching





$$k_{\mathrm{EM}\parallel} = k_{\mathrm{SP}} \pm nq \quad (n = 1, 2, \ldots)$$

q: grating vector

Matching occurs at "resonant" incidence angles (for $\omega_p/\omega \gg 1$)

$$\sin\theta \simeq n\frac{\lambda}{d} - 1$$

($\equiv$ diffraction order along the surface)

- usually $n = 1$
- actually an infinite grating is not strictly required

(a local surface modulation may suffice)

(*) folding in the Brillouin zone – Floquet-Bloch theorem

Observation of “Surfing” Acceleration

PRL **116,** 015001 (2016) PHYSICAL REVIEW LETTERS week ending 8 JANUARY 2016

Electron Acceleration by Relativistic Surface Plasmons in Laser-Grating Interaction

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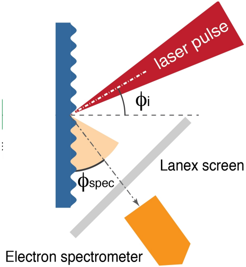
(Received 30 June 2015; published 7 January 2016)

Observation of “Surfing” Acceleration



UHI laser: 25 fs pulse

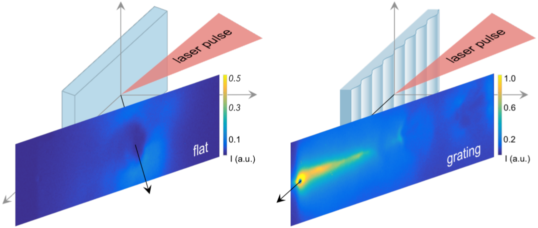
$5 \times 10^{19}\ \mathrm{Wcm}^{-2}$, contrast $\gtrsim 10^{12}$ at 5 ps

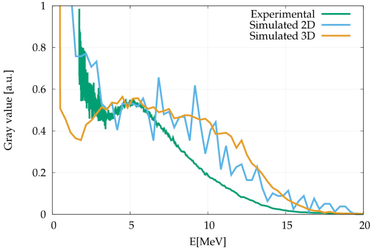
collimated ($\simeq 20°$ cone) electron emission near the surface tangent ($\phi \simeq 2°$)

multi-MeV energy, Total charge $\simeq 100$ pC

(up to $\simeq 650$ pC with blazed gratings:

G. Cantono et al, Phys. Plasmas **25** (2018) 031907)

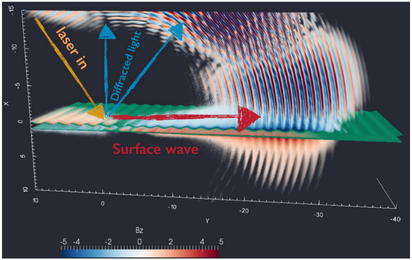


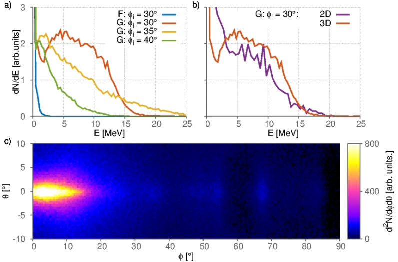


3D Simulations

PICcante code[2]

A.Sgattoni, L.Fedeli, S.Sinigardi et al, arXiv:1503.02464



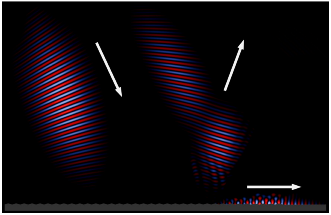


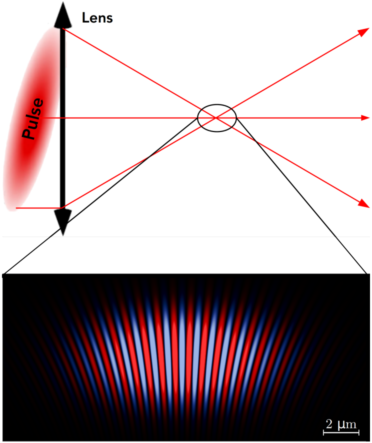
Simulations match experimental observations quantitatively and in detail

[2] Particle-In-Cell Code for AdvaNced simulations on TiEr-0 systems

In-Pulse Rotation of Incidence Angle

Inducing wavefront rotation (WFR) in the laser pulse the effective incidence angle rotates in time $\longrightarrow$ “resonant” condition holds only for an interval shorter than the driving pulse





WFR obtained by focusing a tilted wavefront pulse

Near “Single-Cycle” Surface Plasmon Polariton

ACS Photonics

Cite This: *ACS Photonics* 2018, 5, 1068–1073

Few-Cycle Surface Plasmon Polariton Generation by Rotating Wavefront Pulses

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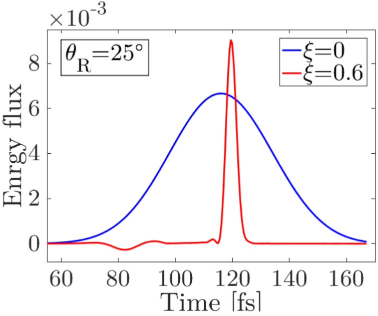
Near "Single-Cycle" Surface Plasmon Polariton

MEEP[3] simulations of
WFR pulse on Ag grating

SPP w/o and with WFR
duration: 3.8 fs (~ 1.4 cycles)

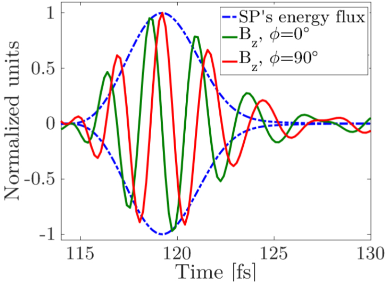


$$E = E(r,z,t)\exp(-i\omega_L t + i r \xi t + \phi)$$

ξ : WFR parameter

laser: 30 fs , $\lambda_L = 0.8\ \mu$m

dependence on
absolute phase ϕ



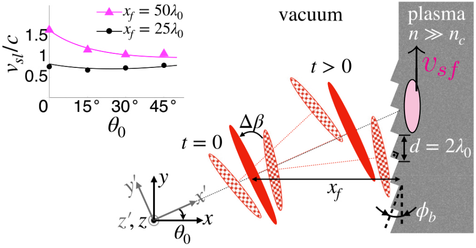
[3] http://ab-initio.mit.edu/wiki/index.php/Meep

WFR Enhancement of SPP Amplitude

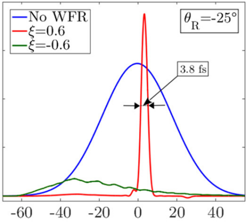
"Sliding focus" effect: WFR makes the laser spot move along the target with velocity

$$v_{sf} \simeq \frac{\Delta\beta\, x_f}{\lambda \cos^2\theta_0} \propto x_f \xi$$

(x_f: waist-to-target distance)



When $v_{sf} \simeq c$ the SPP is "sustained" along its propagation: increase of peak amplitude (Note the effect of the **sign** of ξ)

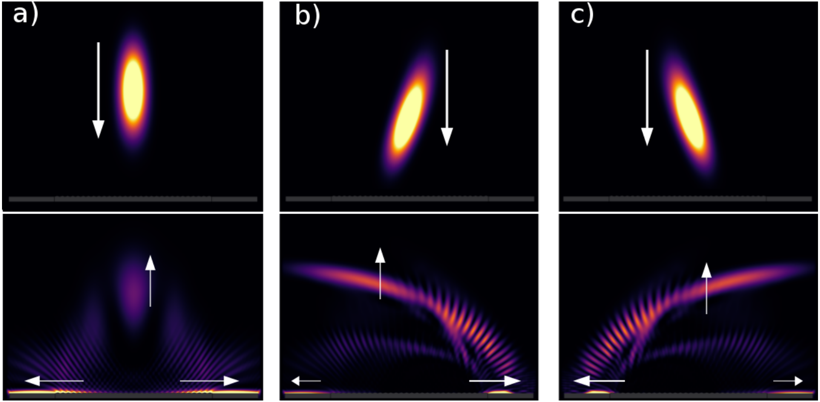


Effect of WFR Direction

Normal incidence: excitation of two symmetric SPs)

a): no rotation b): counterclockwise rotation c): clockwise rotation

Note the scattering from the grating $\longrightarrow$ radiative loss of energy



WFR Effect on Electron Acceleration

PHYSICAL REVIEW E **103**, L021201 (2021)

Letter

Ultrashort high energy electron bunches from tunable surface plasma waves driven with laser wavefront rotation

S. Marini,[1,2] P. S. Kleij,[1,2,3] F. Pisani,[3] F. Amiranoff,[2] M. Grech,[2] A. Macchi,[4,3] M. Raynaud,[1] and C. Riconda[2,*]

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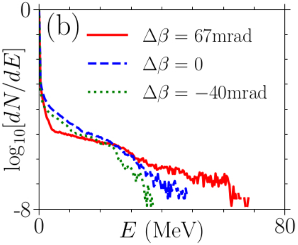
WFR Effect on Electron Acceleration

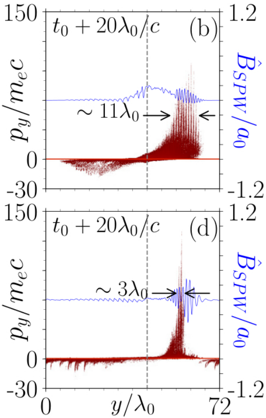
SMILEI open source PIC code

27 fs & 4×10^{19} W cm^{-2} laser pulse

WFR may double the cut-off energy of the electron bunch while shortening its duration down to 8 fs

(simulations by S. Marini and P. Kleij)





Grating Drawbacks

- grating targets are expensive
- need of ultrahigh contrast pulses to preserve the shallow modulation
- strong EM scattering losses of the SP propagating along the grating (inverse to the generation process)

  (might be reduced by having the grating only in the laser spot but would require perfect pointing stability)

$\longrightarrow$ Can we do surfing acceleration *without* a grating?

Acceleration with No Grating Attached

New Journal of Physics

Surface plasmon-driven electron and proton acceleration without grating coupling

J Sarma[1], A McIlvenny[1], N Das[2], M Borghesi[1,*] and A Macchi[3,4,*]

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Dephasing vs Acceleration Lengths

Revisiting (un-)matching between SP and incident laser pulse: phase difference after propagating over a distance L

$$\Delta\phi = (k_{\mathrm{EM}\|} - k_{\mathrm{SP}})L = \left(\frac{1}{v_{\mathrm{EM}}} - \frac{1}{v_{\mathrm{SP}}}\right)\omega L$$

$$\Delta\phi \doteq \pi \quad \longrightarrow \quad L = \frac{\pi}{k_{\mathrm{SP}} - k_{\mathrm{EM}}} \equiv L_{\mathrm{dep}}$$

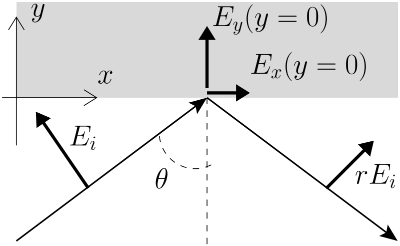
At grazing incidence $(\alpha = \pi/2 - \theta \ll 1)$ $\quad L_{\mathrm{dep}} \simeq \dfrac{\lambda}{\alpha^2 + n_c/n_e}$

Acceleration is not limited by dephasing when

$$L_{\mathrm{dep}} > L_{\mathrm{acc}} = \frac{\mathcal{E}_{\max}}{eE_{\mathrm{SP}}} \simeq \frac{\lambda}{\pi}\frac{n_e}{n_c}$$

Coupling without Grating

To drive the SP efficiently
the incident EM field must have:
1) a non-vanishing component parallel to the surface (E_x)
$\longleftarrow$ reflectivity $R = |r|^2 < 1$



2) a E_x/E_y ratio close as possible to the value for the SP:

$$\left.\frac{E_{\mathrm{EM},x}}{E_{\mathrm{EM},y}}\right|_{y=0^+} = -\varepsilon \frac{1-r}{1+r}\frac{\cos\theta}{\sin\theta} \qquad \left.\frac{E_{\mathrm{SP},x}}{E_{\mathrm{SP},y}}\right|_{y=0^+} = -i|\varepsilon|^{1/2}$$

asymptotically equal for $\sin\theta \to 1$

Laser Spot Stretching Effects

At grazing incidence the laser spot is stretched along x
$\longrightarrow$ the intensity on target decreases

$$I(\theta) = I(0)\cos\theta = I(0)\sin\alpha$$

With respect to the case with grating (typical $\theta \lesssim 45°$) the effect on electron acceleration may be compensated by:

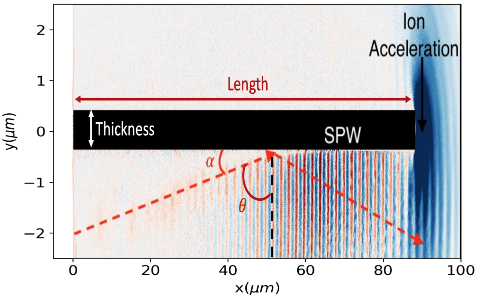
1. slow scaling of electron energy $\mathcal{E}_e \propto E_{\rm SP} \propto \sin^{1/2}\alpha$
2. no loss from radiative scattering
3. no saturation (observed in grating simulations at $a_0 \gtrsim 10$)

Strong (unknown) nonlinear effects on SP are prevented

Simulation Set-Up

PIC code EPOCH
simulations by
J. Sarma & A. McIlvenny

2D Cartesian geometry
Target: fully ionized Au
with CH contaminant layer
electron density
$n_e = 1.7 \times 10^{23}\ \mathrm{cm}^{-3}$
thickness $d = 0.8\ \mu\mathrm{m}$
length $L_T = (90 - 200)\ \mu\mathrm{m}$



Pulse: $\lambda = 0.8\ \mu\mathrm{m}$ ($n_e = 100 n_c$)
Gaussian profiles, width $6.5\ \mu\mathrm{m}$, duration $35\ \mathrm{fs}$ (FWHM)
intensity $I = (0.34 - 7.8) \times 10^{20}\ \mathrm{Wcm}^{-2}$
"relativistic" parameter $a_0 = (5 - 19)$

Electron Spectrum for Different Angles

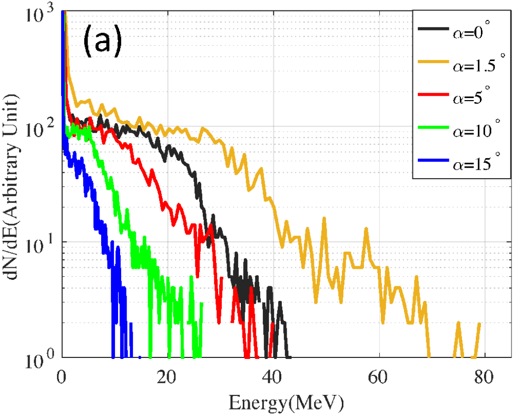
Maximum energy for $\alpha = 1.5°$

Cut–off value doubles with respect to both $\alpha = 5°$ and $\underline{\alpha = 0°}$ (**parallel** incidence, to be discussed later)

($\simeq$ 20 MeV obtained with gratings for similar parameters[4])

$I = 3.4 \times 10^{19}\ \mathrm{Wcm}^{-2}$

$L_T = 90\ \mu\mathrm{m}$



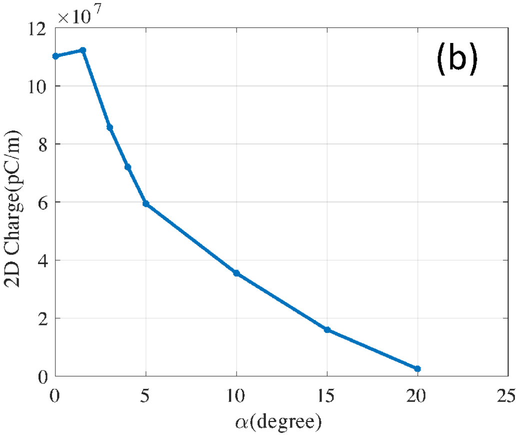
[4] Cantono et al, PoP **25** (2018) 031907

Accelerated Charge for Different Angles

Peak value $\simeq 10^8\ \mathrm{pCm}^{-1}$ in 2D corresponds to $\simeq 780\ \mathrm{pC}$ in 3D ($\simeq 660\ \mathrm{pC}$ obtained with blazed gratings for similar parameters[5])

$I = 3.4 \times 10^{19}\ \mathrm{Wcm}^{-2}$

$L_T = 90\ \mu\mathrm{m}$



[5] Cantono et al, PoP **25** (2018) 031907

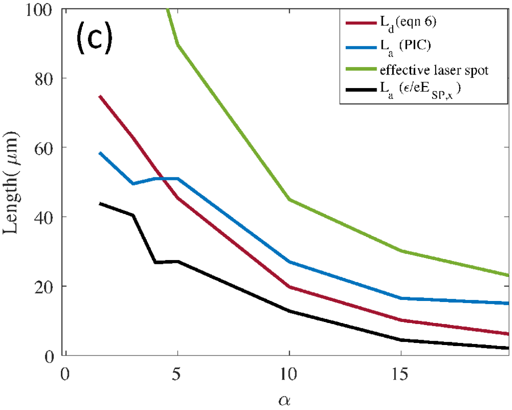
Acceleration Length for Different Angles

Two different estimates of the acceleration length (L_{acc}) are fairly close to the dephasing length L_{dep} as a function of α

Spot width on target $L_S > L_{\mathrm{dep}}$ (could be further optimized)

$I = 3.4 \times 10^{19}\ \mathrm{Wcm}^{-2}$

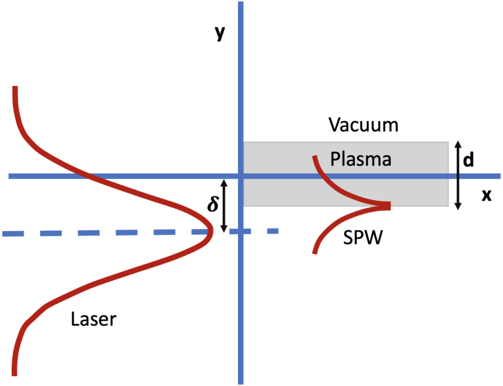
$L_T = 200\ \mu\mathrm{m}$



Parallel Incidence

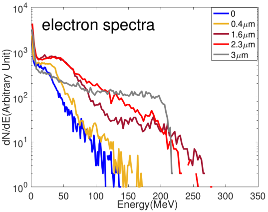
For **parallel** incidence ($\alpha = 0°$) the laser pulse will interact with the (left) short edge of the target

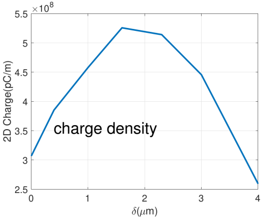
Additional parametric dependence on the "focal spot shift" δ $\equiv$ distance between the laser propagation axis and the target midplane (surface at $y = -\delta/2$)

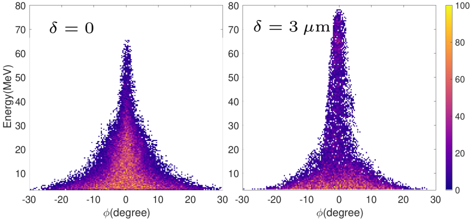


Energy and Charge Maxima for $\delta \neq 0$

$I = 6 \times 10^{20}\ \mathrm{Wcm}^{-2}$
$\mathcal{E}_{\max} \simeq 250\ \mathrm{MeV}$
for $\delta \simeq 2\ \mu\mathrm{m}$
Total charge $\simeq 3.4\ \mathrm{nC}$
(3D estimate)
for $\delta = 1.6\ \mu\mathrm{m}$

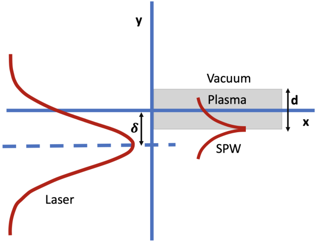






Energy-angle distributions ($I = 3.4 \times 10^{19}\ \mathrm{Wcm}^{-2}$)
Electrons are strongly collimated with almost symmetrical distribution even for "asymmetrical" interaction with $\delta \neq 0$

On the effect of "shift" δ



Similar coupling conditions as for grazing incidence:

Laser field components (in vacuum)

$$E_{\mathrm{EM},y} \simeq E_0 \exp\left(-(y+\delta)^2/w^2\right) \exp\left(ik_{\mathrm{EM},x}\right)$$

$$E_{\mathrm{EM},x} \simeq 2y/(ik_{\mathrm{EM},x}w^2) \exp\left(-(y+\delta)^2/w^2\right) \exp\left(ik_{\mathrm{EM},x}\right)$$

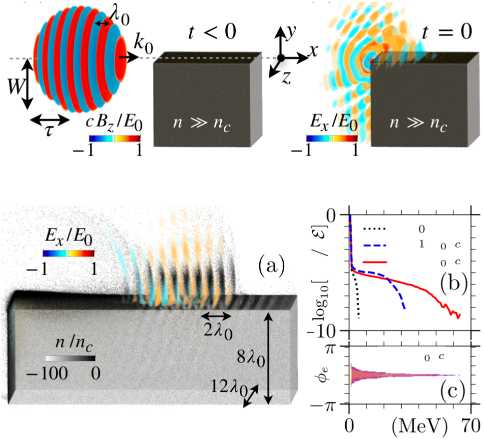
$$\left|\frac{E_{\mathrm{EM},x}}{E_{\mathrm{EM},y}}\right|_{y=-d/2} \simeq \frac{(\delta - d/2)\lambda}{\pi w^2}$$

For $w = 3.9\mu\mathrm{m}$, the $|E_x|/|E_y|$ ratio at the surface ($y = -d/2$) for $\delta = 2.3\ \mu$m is the same as for $\alpha = 1.7°$.

A Possibly Similar Observation?

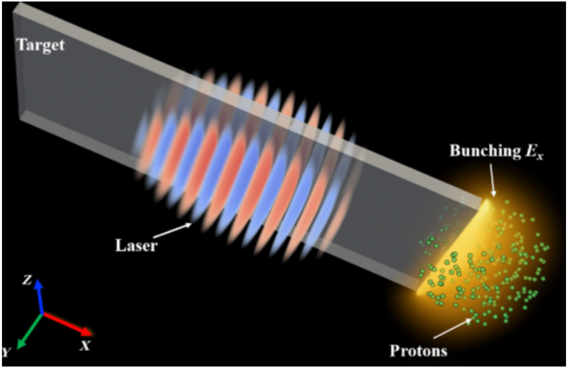
3D SMILEI simulations

S.Marini, P.Kleij, M.Grech, M.Raynaud, C.Riconda, "Electron acceleration by laser plasma wedge interaction" arXiv:2202.08226



Proposed "Peeler" Proton Acceleration

"A fs laser (red and blue) is incident on the edge of a micron-thick tape (grey) [. . .] Abundant electrons are accelerated forward by the intense laser."

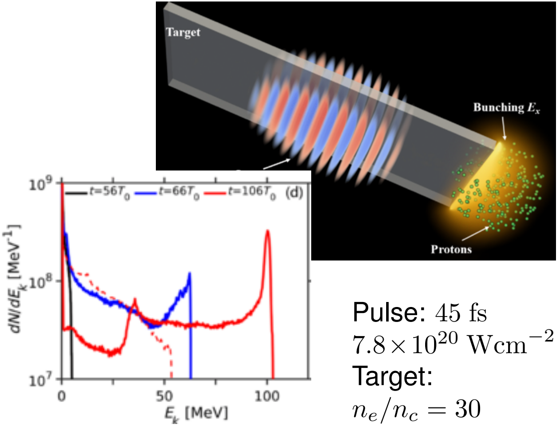


3D simulations with VLPL and EPOCH codes

X.Shen, A.Pukhov, B.Qiao,
Phys. Rev. X **11** (2021) 041002

Proposed “Peeler” Proton Acceleration

“[. . .] at the rear edge a longitudinal bunching field is established (yellow). Protons (green dots) are simultaneously accelerated and leading to a highly monoenergetic beam.”



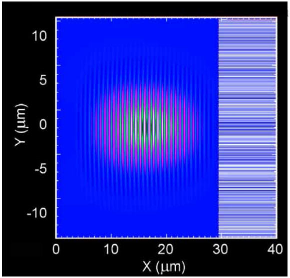
Pulse: 45 fs
7.8×10^{20} Wcm^{-2}
Target:
$n_e/n_c = 30$
$d = 50$ nm

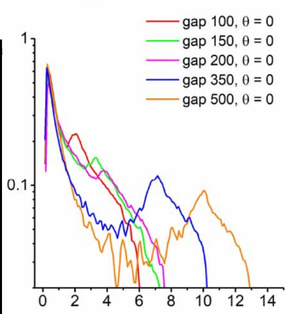
X.Shen, A.Pukhov, B.Qiao,
Phys. Rev. X **11** (2021) 041002

Earlier Numerical Observation?

2D simulations of an array of parallel foils:
electron acceleration attributed to SP
peaked proton spectra appear for large gaps

Pulse: 30 fs
2.8×10^{20} Wcm^{-2}
Target: $n_e/n_c = 60$
$d = 0.6\ \mu$m



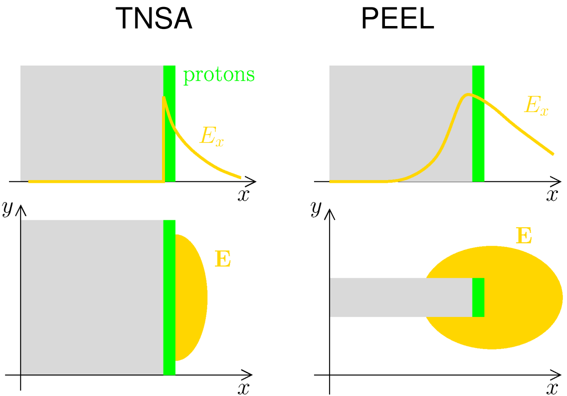


G. Cristoforetti et al, Plasma Phys. Control. Fusion **62** (2020) 114001

Origin of Monoenergetic Proton Spectra

TNSA: fast electrons are less than protons in the layer which screen the $\mathbf{E}$-field producing a sharp gradient

PEEL: protons are less than fast electrons and the space charge $\mathbf{E}$-field on the proton layer is spatially "smooth".



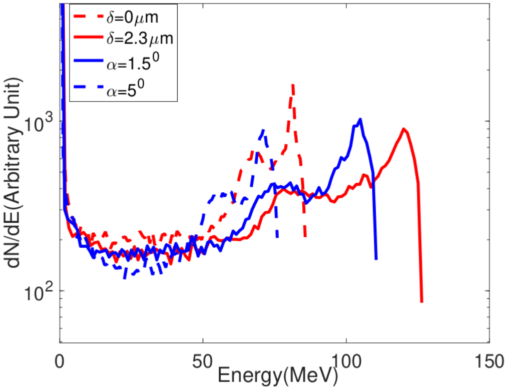
All protons experience almost the same field $\longrightarrow$ monoenergetic acceleration

Proton Spectra: 2D simulation (High I)

Highest cut-off energy is reached for parallel incidence with "shifted" pulse ($\delta = 2.3\ \mu\text{m}$)

Slightly lower energy at grazing incidence ($\alpha = 1.5°$)

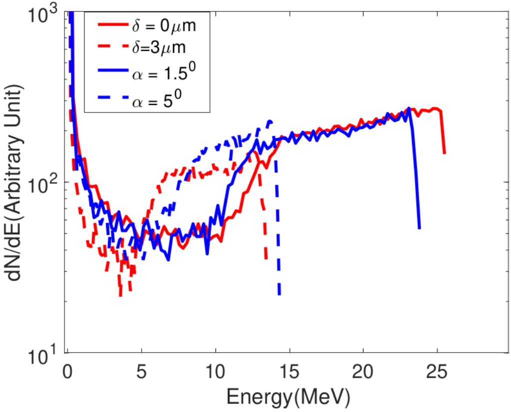
$I = 6 \times 10^{20}\ \text{Wcm}^{-2}$



Proton Spectra: 2D simulation (Low I)

Sharply peaked spectra are not apparent for lower intensity

$I = 3.4 \times 10^{19}\ \mathrm{Wcm}^{-2}$



Proton Spectra: 3D simulation (preliminary)

- electron spectra similar to 2D
- Shen et al.'s proton spectra reproduced

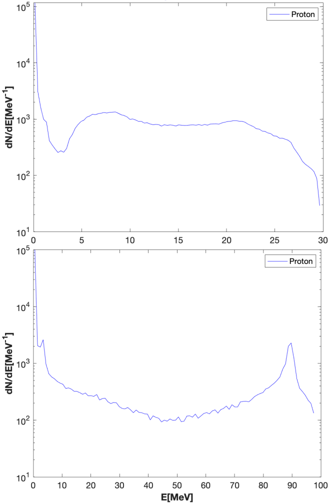
For a lower energy pulse (GEMINI class) the proton spectrum is deeply affected

- Geometrical effect on proton acceleration?
- Insufficient electron production?

$I = 6 \times 10^{20}\ \mathrm{Wcm}^{-2}$

pulse width $3\ \mu$m

(4 times narrower than Shen's)



Conclusions

- Surface Plasmons may be efficiently excited without a grating (no phase matching)
  - at grazing incidence
  - at parallel incidence with enhancement from “shifting” the laser pulse direction
- the SP is sustained along a distance sufficient to accelerate electrons to high energies (possibly beyond previous results with grating targets)
- The “peeler” schemes exploits the high charge SP-driven electrons for mono-energetic acceleration of protons
  - Grazing incidence and “shift” effect enhance proton acceleration
  - preliminary 3D simulations suggest that large pulse energy is needed to sustain the acceleration

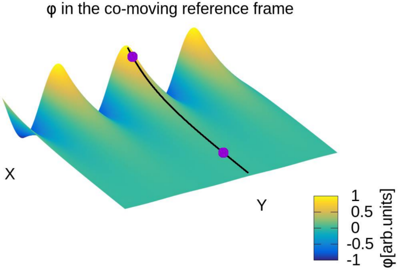
EXTRA SLIDES

Simple Model of SP Acceleration - I

SP field on the vacuum side is electrostatic in the wave frame S' moving with phase velocity $\beta_{\rm p} = v_{\rm p}/c$ with respect to S (lab)
Electrostatic potential in S':

$$\Phi' = -\left(\frac{\gamma_{\rm p} E_{\rm SP}}{k}\right) e^{k'x} \sin k'y' \qquad k' = \frac{k}{\gamma_{\rm p}} \qquad \gamma_{\rm p} = (1-\beta_{\rm p}^2)^{-1/2}$$

The motion is 2D: the energy gain depends on the “kick angle” from the top of the potential hill



Simple Model of SP Acceleration - II

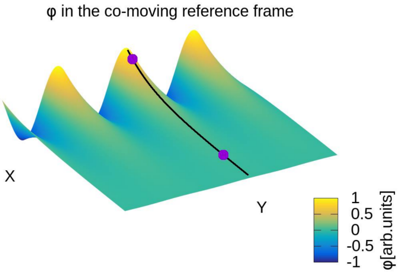
Assume as the most likely case an electron going downhill along the x-direction and acquiring an energy $W' = eE_{\mathrm{SP}}/k'$

$$W \simeq \gamma_{\mathrm{p}} W' \simeq m_e c^2 a_{\mathrm{SP}} \frac{\omega_p^2}{\omega^2} \qquad (a_{\mathrm{SP}} = eE_{\mathrm{SP}}/m_e \omega c)$$

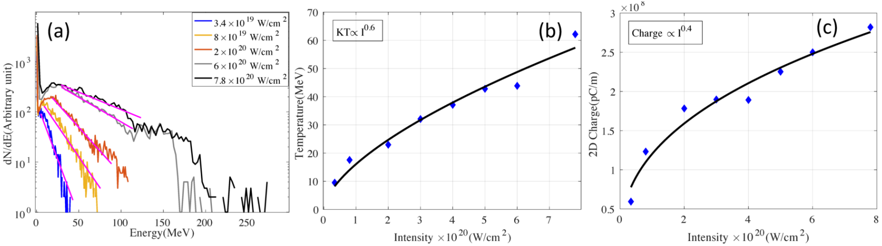
with ejection angle in L (for $W' \gg m_e c^2$)

$$\tan\phi_e = \frac{p_x}{p_y} \simeq \frac{1}{\gamma_{\mathrm{p}}}$$

$\rightarrow$ high energy electrons are beamed near the surface ($\tan\phi_e \ll 1$)



Scaling with Laser Intensity

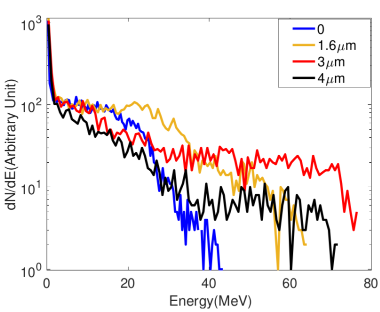


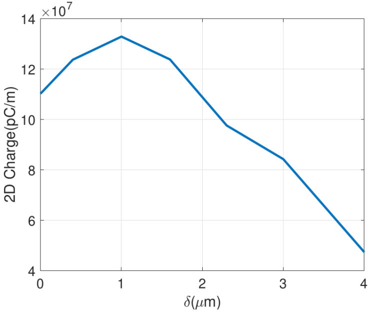
a) electron spectra b) electron “Temperature” c) charge density

Temperature exceeds “ponderomotive” values ($T_{\rm p} = m_e c^2(\sqrt{(1 + a_0^2/2)} - 1)$) by one order of magnitude

Peak charge density value $\simeq 3 \times 10^8\ \mathrm{pCm}^{-1}$ in 2D corresponds to an estimate $\simeq 1.9\ \mathrm{nC}$ in 3D

Effect of δ on Spectra and Charge (Low I)

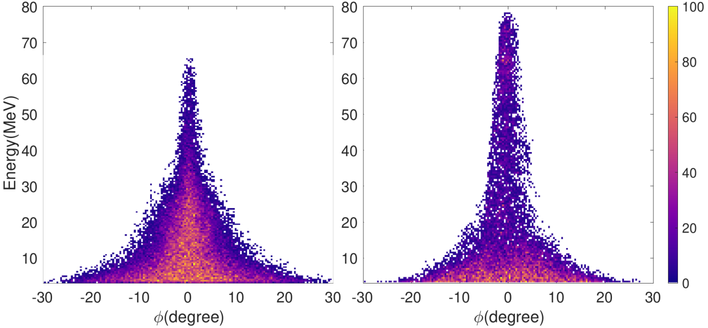




a) electron spectra b) charge density

$I = 3.4 \times 10^{19}\ \mathrm{Wcm}^{-2}$
Maximum values are *not* for $|\delta| \leq d/2$

Electron Collimation

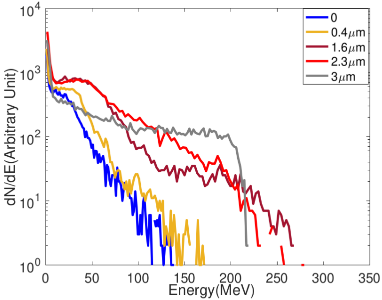


a) $\delta = 0$ b) $\delta = 3\ \mu\mathrm{m}$

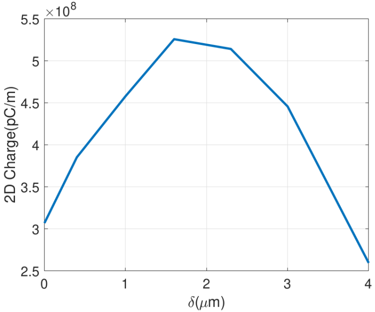
$I = 3.4 \times 10^{19}\ \mathrm{Wcm}^{-2}$

Electrons are strongly collimated with almost symmetrical distribution even for "asymmetrical" interaction with $\delta \neq 0$

Effect of δ on Spectra and Charge (High I)



a) electron spectra



b) charge density

$I = 6 \times 10^{20}\ \mathrm{Wcm}^{-2}$

Total 3D charge would be $\simeq 3.4\ \mathrm{nC}$ for $\delta = 1.6\ \mu\mathrm{m}$