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**Laser-Driven Ion Acceleration,
High Field Plasmonics, and
Classical Instabilities**

Andrea Macchi

Event:
INO Annual Symposium 2014

INO Sensor Lab Section
University of Brescia, Ingegneria Area
Via Branze, 38 – Brescia, Italy

Laser-driven ion acceleration, high field plasmonics, and classic instabilities

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INO Annual Symposium, Brescia, Italy
Friday, October 3, 2014



Outlook

- ▶ Laser-plasma ion acceleration
- ▶ High field plasmonics for enhanced acceleration (and beyond)
 - surface wave coupling with grating targets
 - studies on tapered nanoguides
- ▶ Light sail acceleration: high gain and stability issues
 - observation in simulation and experiment
 - plasmonic effects on Rayleigh-Taylor instability

Main coworkers for this talk

L. Fedeli¹, A. Sgattoni, S. Sinigardi², F. Pegoraro¹

CNR/INO, Pisa, Italy

¹Dipartimento di Fisica “Enrico Fermi”, Università di Pisa, Italy

²presently at Dipartimento di Fisica, Università di Bologna, Bologna, Italy

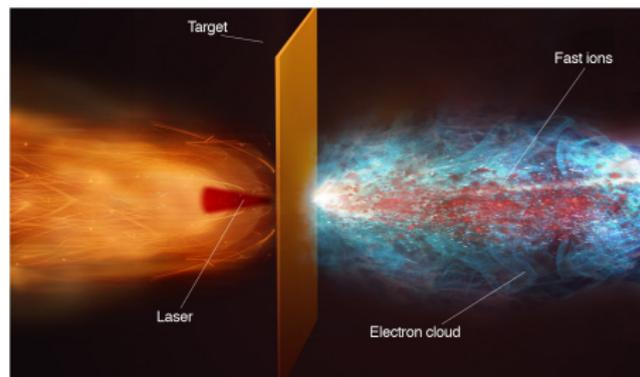
Plus colleagues from several institutes abroad (Queen’s University – Belfast, LULI Ecole Polytechnique – Palaiseau, CEA/LyDyL – Saclay, Technical University –Prague, . . .)

Laser-plasma acceleration of ions (2000–)

A. Macchi, M. Borghesi,
M. Passoni,
Rev. Mod. Phys. **85** (2013) 571

State of the art (2013):

- up to $\simeq 70$ MeV protons observed
(claims of > 100 MeV announced)
- $> 10^{13}$ protons, $> 10^{11}$ C ions accelerated in single shots
(as charge neutralized bunches)
- very low emittance measured ($< 0.1\pi$ mm mrad)
- proofs-of-principle of spectral manipulation and beam focusing



Applications

- Picosecond radiography of plasmas (sensitive to EM fields)
- Warm dense matter production and probing
- Short-lived isotope production
- Hadrontherapy
- High energy physics

Need to increase energy and/or optimize energy spectrum and/or work at high repetition rate . . .

Sheath acceleration of protons

Target Normal Sheath Acceleration (TNSA) is driven by “fast” ($\mathcal{E}_e \sim \text{MeV}$) electrons generated in thin targets: protons from surface contaminants are accelerated in the rear sheath

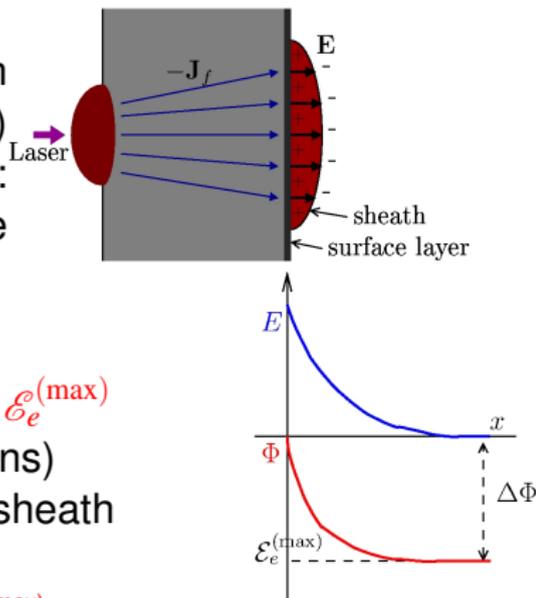
TNSA picture “for dummies”:

Potential drop for static sheath $e\Delta\Phi = \mathcal{E}_e^{(\text{max})}$
(sheath potential must confine electrons)

Energy gained by “test” proton in the sheath

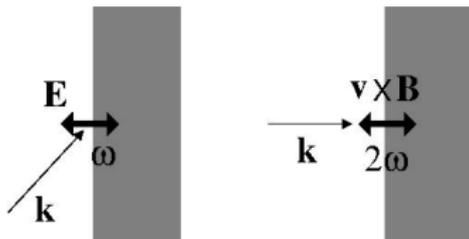
$$\mathcal{E}_p = e\Delta\Phi = \mathcal{E}_e^{(\text{max})}$$

Increasing electrons' energy \rightarrow increasing protons' energy



Fast electron generation: simple (rough) picture

The Lorentz force of the laser wave (amplitude E_L , frequency ω) drives periodic “push-pull” of electrons across the density gradient

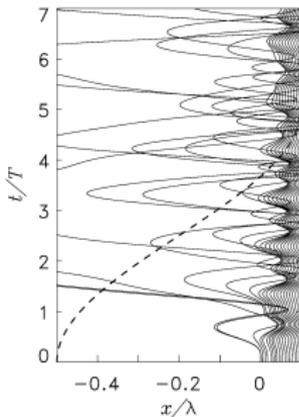


Electrons perform “half-oscillations” on the vacuum side and re-enter in the plasma (where the EM field is screened) and are “absorbed” keeping a net momentum

$$p_e \sim p_{\text{osc}} \sim eE_L/\omega \equiv m_e c a_0$$

$$a_0 = \left(\frac{eE_L}{m_e c \omega} \right) > 1 \longrightarrow \text{relativistic electrons}$$

A. Macchi, *A Superintense Laser-Plasma Interaction Primer* (Springer, 2013)



Plasmonics at high fields for enhanced coupling

Concept: exploit **collective electron excitation** in **sub-wavelength structures** to concentrate light near the surface

Strategy: excitation of **surface waves** in **periodic structures** (“Catching light by a **grating**”)

Phase matching possible at a **plasma-vacuum** interface if the shallow grating is preserved!

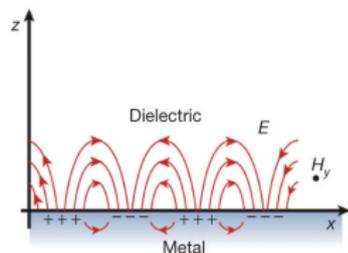
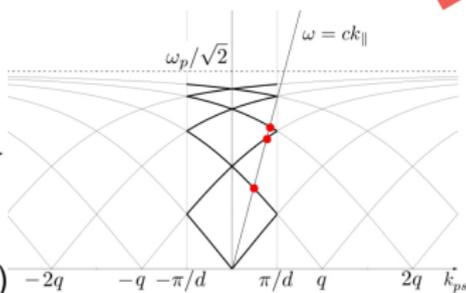
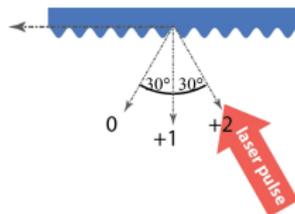


Fig.: O. Benson,
Nature **480**, 193 (2011)



Question: is this possible in the **high field, relativistic** regime?

Evidence of Resonant Surface-Wave Excitation in the Relativistic Regime through Measurements of Proton Acceleration from Grating Targets

T. Ceccotti,^{1,*} V. Floquet,¹ A. Sgattoni,^{2,3} A. Bigongiari,⁴ O. Klimo,^{5,6} M. Raynaud,⁷ C. Riconda,⁴ A. Heron,⁸ F. Baffigi,² L. Labate,² L. A. Gizzi,² L. Vassura,^{9,10} J. Fuchs,⁹ M. Passoni,³ M. Květon,⁵ F. Novotny,⁵ M. Possolt,⁵ J. Prokūpek,^{5,6} J. Proška,⁵ J. Pšikal,^{5,6} L. Štolcová,^{5,6} A. Velyhan,⁶ M. Bougeard,¹ P. D'Oliveira,¹ O. Tcherbakoff,¹ F. Réau,¹ P. Martin,¹ and A. Macchi^{2,11,†}

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⁹LULI, UMR7605, CNRS-CEA-Ecole Polytechnique-Paris 6, 91128 Palaiseau, France

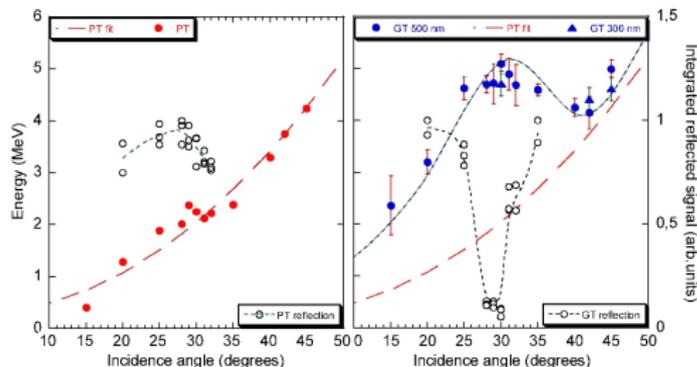
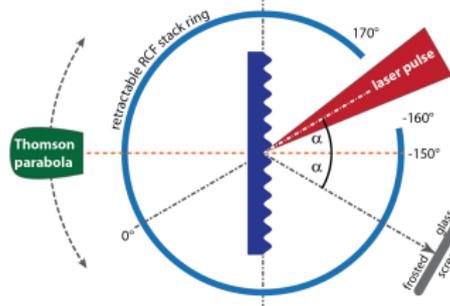
¹⁰Dipartimento SBAI, Università di Roma "La Sapienza," Via A. Scarpa 14, 00161 Roma, Italy

¹¹Dipartimento di Fisica "Enrico Fermi," Università di Pisa, Largo Bruno Pontecorvo 3, I-56127 Pisa, Italy

Experimental evidence: grating-enhanced TNSA

LaserLAB experiment at SLIC, CEA Saclay
laser UHI, 28 fs, $5 \times 10^{19} \text{ W cm}^{-2}$

contrast $\sim 10^{12}$

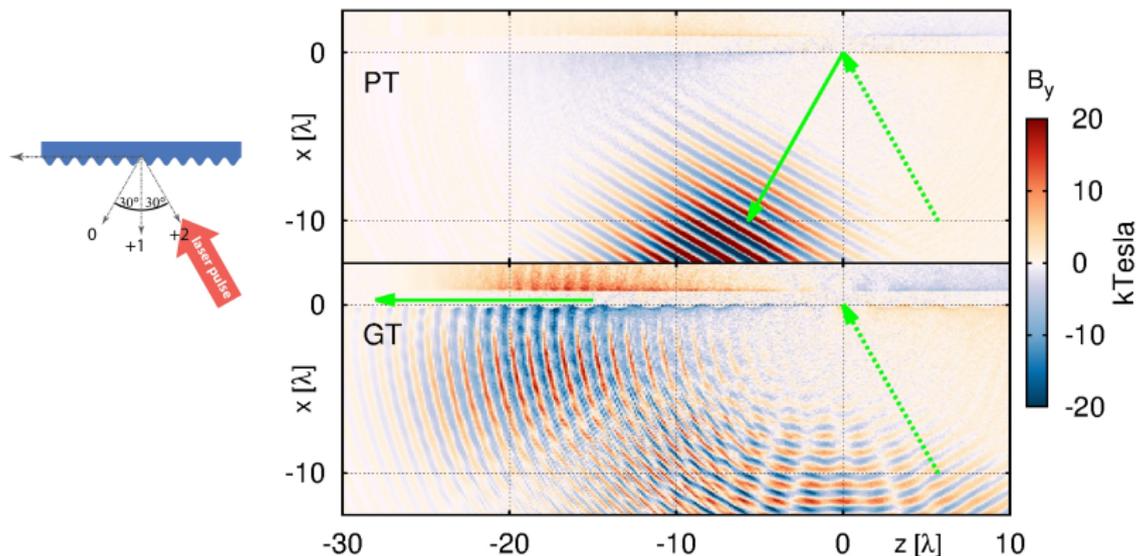


3X increase observed in proton energy at resonant angle

T.Ceccotti et al, PRL **111** (2013) 185001

Surface wave in PIC simulations

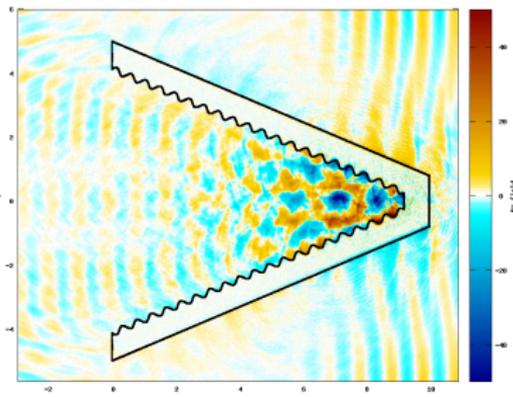
Snapshots of EM fields show **localized wave** propagating along the surface at resonant angle of incidence (plus reflection at various diffraction orders)



Exploring high field plasmonics with simulations

First PIC simulations of a tapered waveguide for light nano-focusing and amplification

If 100X field amplification confirmed for ultraintense lasers, $10^{25} \text{ W cm}^{-2}$ intensity is reached!



Simulations performed on FERMI super-computer with the Open Source Particle-In-Cell code PICcante originally developed by L.Fedeli, A.Sgattoni, S.Sinigardi
github.com/ALaDyn/piccante



Radiation Pressure Acceleration: “Light Sail”

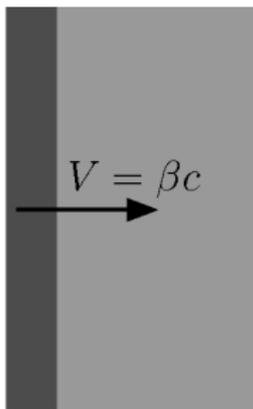
Model:

Perfect mirror boosted
by a plane wave:

mechanical efficiency η and
momentum transfer to mirror
derived by Doppler shift and
photon number conservation

I, ω

I_r, ω_r



$$\frac{dp}{dt} = \frac{2I}{c} \frac{1-\beta}{1+\beta} \quad \eta = \frac{2\beta}{1+\beta}$$

High efficiency ($\eta \rightarrow 1$) but slow gain ($dp/dt \rightarrow 0$) as $\beta \rightarrow 1$

Light Sail formulas and scaling

$$\begin{aligned} \mathcal{E}_{\max} &= m_p c^2 \mathcal{F}^2 / (2(\mathcal{F} + 1)) \\ &\simeq m_p c^2 \mathcal{F}^2 / 2 \quad (\mathcal{F} \ll 1) \end{aligned}$$

$$\mathcal{F} = 2(\rho \ell)^{-1} \int_0^\infty I(t') dt' \simeq 2I\tau_p / \rho \ell$$

$$\mathcal{E}_{\text{ion}}(t) \propto (2It / \rho \ell c^2)^{1/3}$$

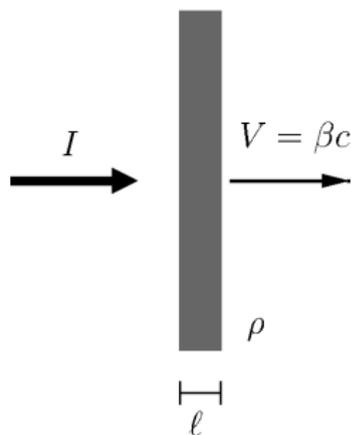
(for $t \gg \rho \ell c^2 / I$, $\mathcal{E}_{\text{ion}} > m_p c^2$)

Favorable scaling with dimensionless laser pulse fluence \mathcal{F}

“Perfect” monoenergeticity for “rigid”, coherent sail motion

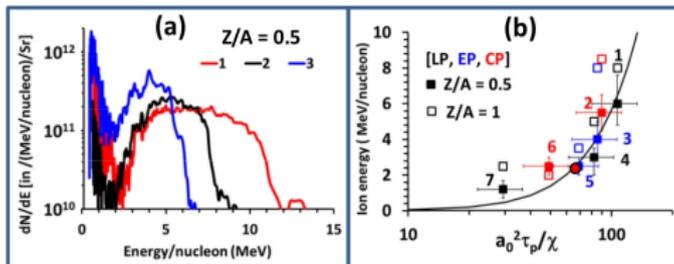
Need of ultrathin (nm) foils and ultrahigh contrast pulses

Issues: slow energy gain, heating, transparency, deformation ...



\mathcal{F}^2 scaling experimentally observed

VULCAN laser, RAL/GLF:
Laser pulse: $t_p \simeq 800$ fs
 3×10^{20} W cm $^{-2}$
 $\sim 10^9$ contrast
Target: ~ 0.1 μ m metal foil



Multispecies ($Z/A = 1, 1/2$) peaks observed with $\Delta\mathcal{E}/\mathcal{E} \simeq 20\%$
Up to $\simeq 10$ MeV/amu observed at high flux
Simulations suggest > 100 MeV/nucleon are within reach

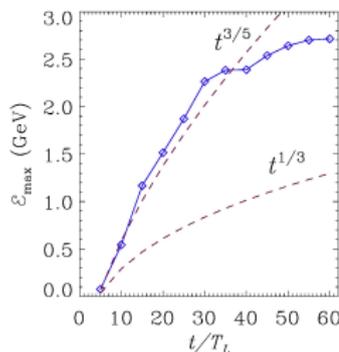
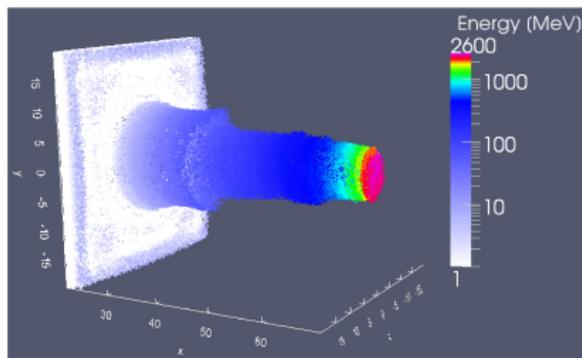
Kar, Kakolee, Qiao, Macchi, Borghesi et al PRL **109** (2012) 185006

Other recent expts: Steinke et al, PRST-AB **16** (2013) 11303;
Aurand et al, NJP **15** (2013) 33031

High energy gain in 3D simulation of ELI regime

Laser: 24 fs, 4.8 μm spot, $I = 0.85 \times 10^{23} \text{ W cm}^{-2} \implies 1.5 \text{ kJ}$

Target: 1 μm foil, $n_e = 1.1 \times 10^{23} \text{ cm}^{-3}$, $\zeta \simeq a_0 \simeq 200$

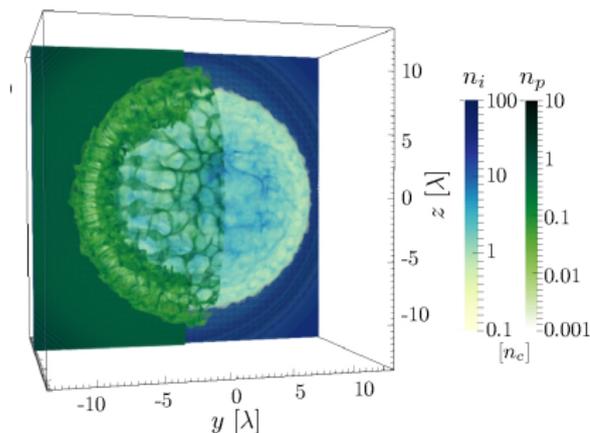


$E_{\text{max}} \simeq 2.6 \text{ GeV} > 4$ times 1D model prediction

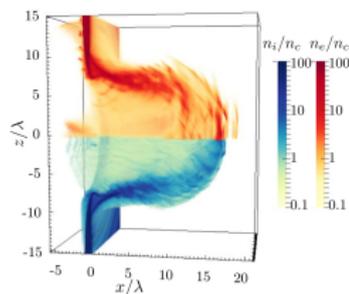
Macchi et al, Plasma Phys. Contr. Fus. **55** (2013) 124020

Sgattoni et al, Appl. Phys. Lett. **105** (2014) 084105

Transverse Structures in 3D Simulation



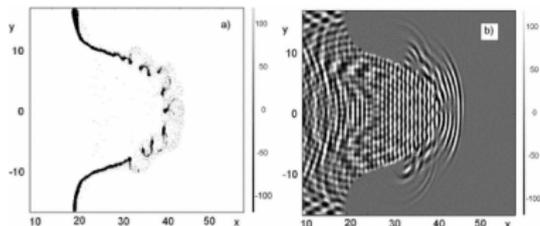
- Code: ALaDyn (originally developed by Benedetti, Londrillo, Sgattoni, Turchetti – Bologna University)
- Machine: FERMI BlueGene/Q at CINECA sponsored by PRACE
- Set-up: 4096×1792^2 grid points, 2×10^{10} particles, 16384 cores used
- $10^{23} \text{ W cm}^{-2}$ laser pulse on solid target (sub- μm thickness)



Formation of **net-like structures** with size $\sim \lambda$ (laser wavelength)
What is their origin?
What sets the dominant scale?

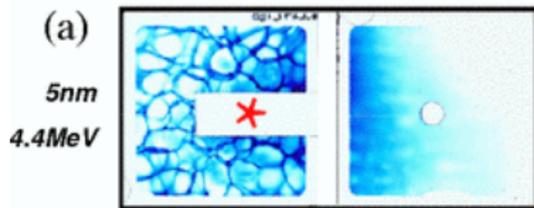
Rayleigh-Taylor Instability in RPA?

Thin foil target of areal density σ accelerated by a laser of intensity I is unstable with growth rate $\gamma = (P_0 q / \sigma)^{1/2}$ with $P_0 = 2I/c$ and q the wavevector [Ott, PRL **29** (10972) 1429]



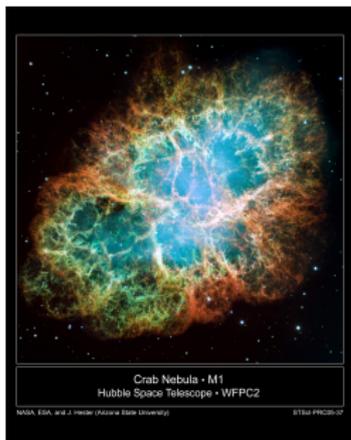
2D simulation

[F.Pegoraro & S.V.Bulanov,
PRL **99** (2007)065002]



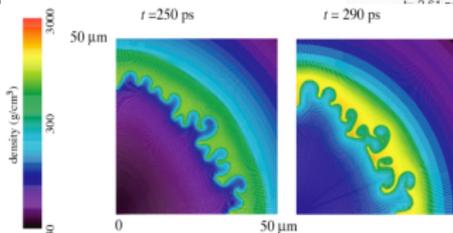
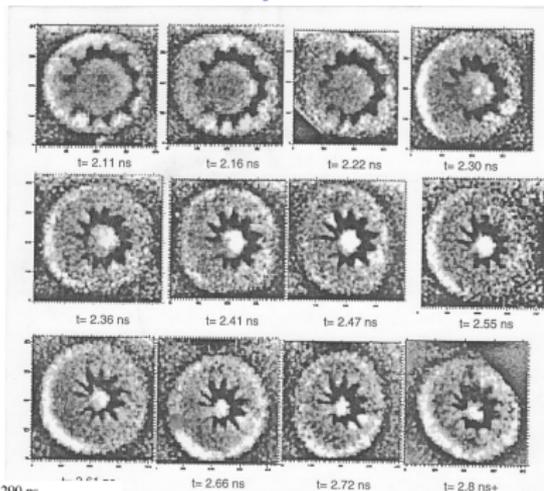
Experimental indication from
accelerated ion beam profile
structures [Palmer et al, PRL **108**
(2012) 225002]

Rayleigh-Taylor Instability in space and lab plasmas



Crab Nebula,
Hubble Space
Telescope

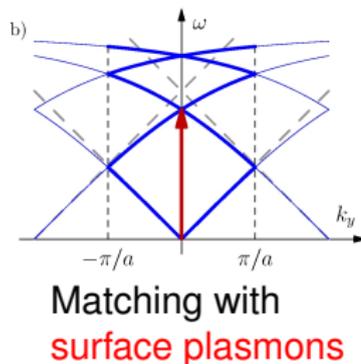
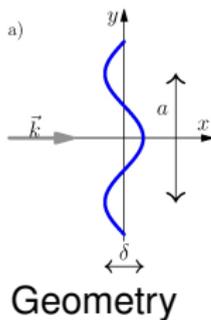
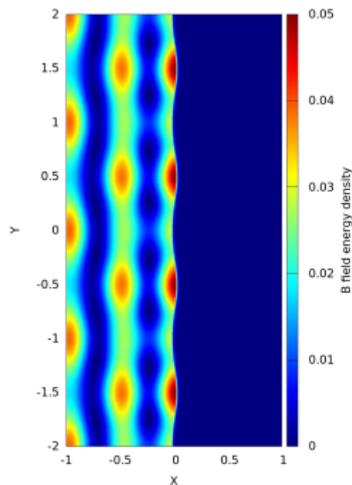
Laser-driven
implosion for
Inertial
Confinement
Fusion studies,
1995
(Wikipedia)



RTI simulation
GAPS group,
Roma
(S.Atzeni et al)

Plasmonic modulation of radiation pressure

The EM field at a rippled surface (e.g. 2D reflecting, sinusoidal grating of period d) is modulated with **plasmonic enhancement** of the P -component when $d \sim \lambda$



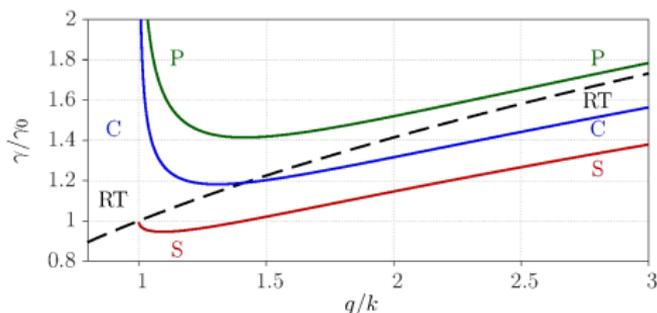
The resulting **modulation of laser light pressure** provides a **spatial seed** for RTI
[A.Sgattoni et al, [arXiv:physics/1404.1260](https://arxiv.org/abs/1404.1260)]

Thin foil RTI with self-consistent pressure modulation

Model: reflection from shallow 2D grating of depth δ (first order in δ/λ) + modified Ott's theory with modulated pressure:

$$P \simeq P_0(1 + K(q)\delta \cos qy), \quad K(q) = \begin{cases} -(q^2 - k^2)^{1/2} & (S) \\ k^2 q (q^2 - k^2)^{-1/2} & (P) \\ (k^2 - q^2/2)(q^2 - k^2)^{-1/2} & (C) \end{cases}$$

$$\gamma = (P_0/\sigma)^{1/2} \left[(q^2 + K^2(q)/4)^{1/2} + K(q)/2 \right]^{1/2}$$



S-polarization

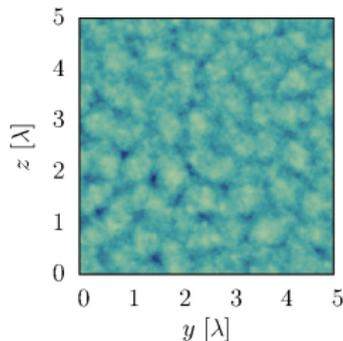
P-polarization

C-ircular polarization

RT: no modulation ($\delta = 0$)

Symmetry of RTI structures

Nonlinear hexagonal-like structures generated by RTI have “wallpaper” p6m symmetry



3D sim. (plane wave)



Persian glazed tile

An example of **spontaneous symmetry breaking** in a **classical** system (there's not only quantum physics and the Higgs ...)

EXTRA SLIDES

The paradigm of “coherent” acceleration

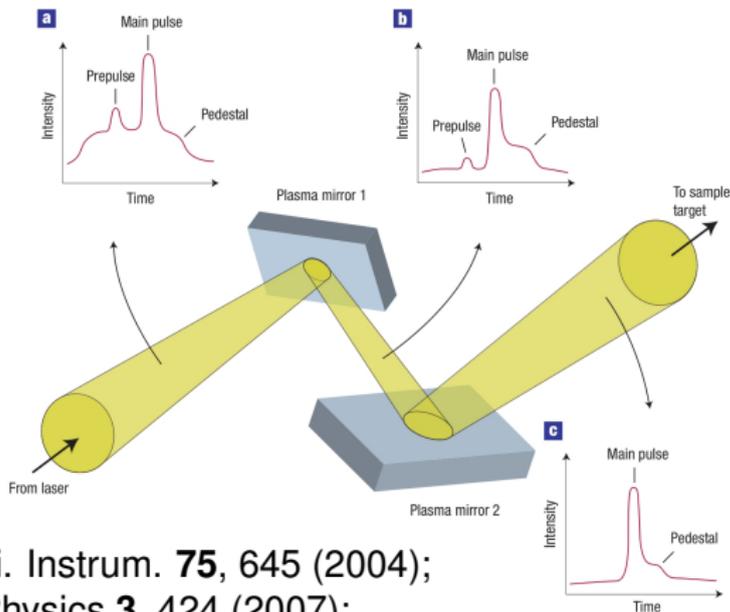
Early vision: V. Veksler, At. Energ. **2** (1957) 525

- ▶ accelerating field on each particle proportional to the number of accelerated particles
- ▶ automatic synchrony between the particles and the accelerating field
- ▶ field localization in the region where the particles are
- ▶ acceleration of quasi-neutral bunches with large numbers of particles
- challenge: control of the dynamics of the plasma (complex many-body system, leaning to instability)

Need for ultraclean pulses: plasma mirrors

Plasma mirrors

yielding $\sim 10^{12}$ pulse-to-prepulse contrast allow to preserve target structuring until the short pulse interaction



B. Dromey et al, Rev. Sci. Instrum. **75**, 645 (2004);

C. Thaury et al, Nature Physics **3**, 424 (2007);

figure from P. Gibbon, *ibid.*, 369.