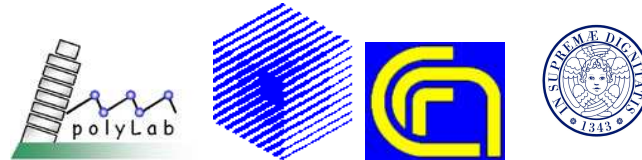


Introduction to Plasma Discharges for Material Processing

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References

- [1] J. Reece Roth, *Industrial Plasma Engineering*, Voll.1–2 (Institute of Physics Publishing, 2004)
- [2] M. A. Lieberman and A. J. Lichtenberg, *Principles of Plasma Discharges and Materials Processing*, Second Edition (John Wiley and Sons, 2005)
- [3] K. Ostrikov, “*Colloquium: Reactive plasmas as a versatile nanofabrication tool*”, *Rev. Mod. Phys* **77**, 489 (2005)
- [4] C.-M. Chan, T.-M- Ko, H. Hiraoka, “*Polymer surface modification by plasmas and photons*”, *Surf. Science Rep.* **24**, 1 (1996)
- [5] A. Macchi, *Appunti su scariche di plasma per applicazioni tecnologiche*, and references therein (2005–, in progress!)

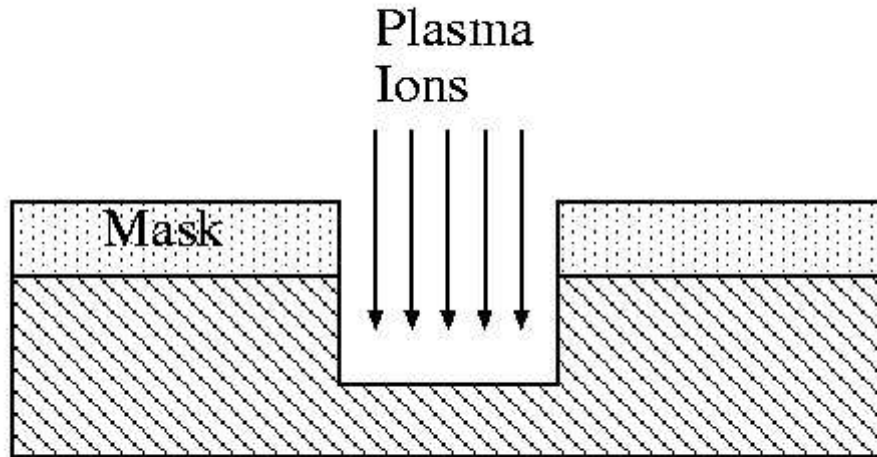
Example 0: nature



Principle of plasma chemistry: discharge creates chemically reactive species (either ions or neutrals)
Lightning creates ozone, fixed nitrogen and light ^a

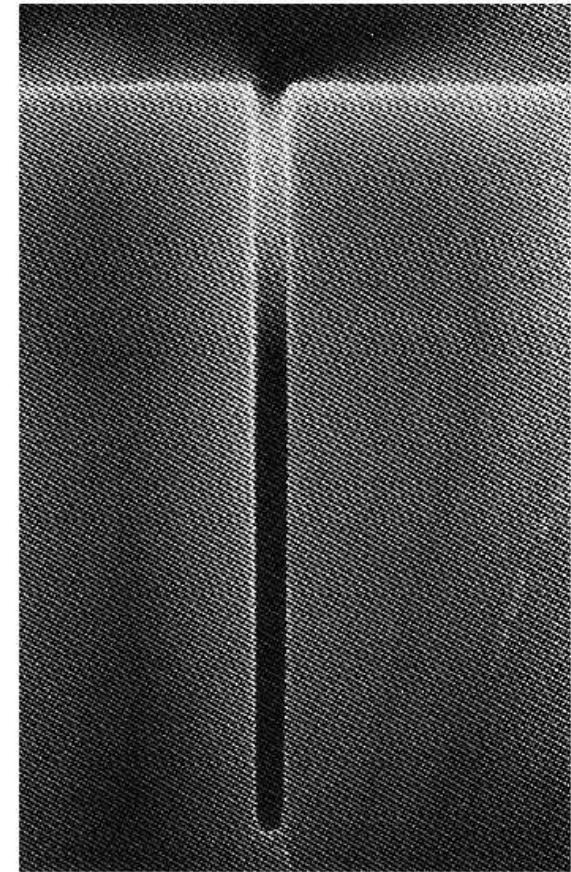
^asee e.g. M. A. Uman, *The lightning discharge* (Dover, 2001)

Example 1: *etching* on Silicon



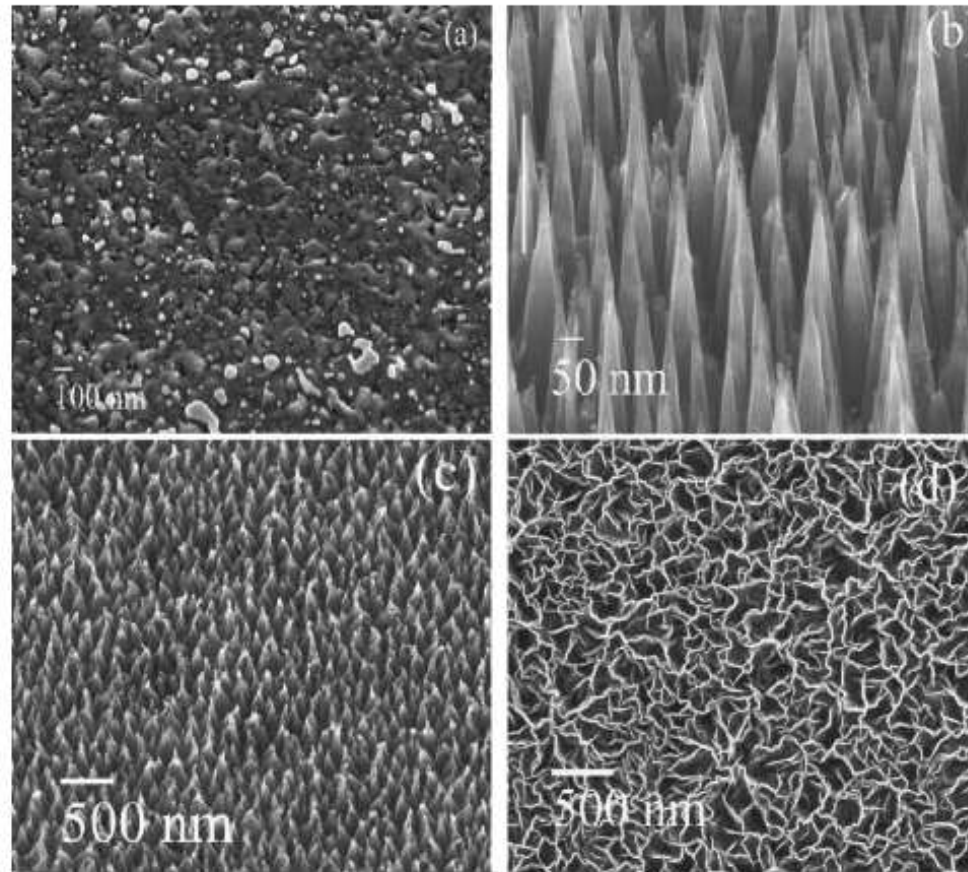
Scheme of directional plasma etching on the surface of Si, for the production of microchips

(see [1], p.2)



Result example in monocrystalline Si (dimensions: $0.2 \times 4 \mu\text{m}$)

Example 2: Carbon nanostructures



Scanning electron micrograph [3] of
a) Silicon surface *before* hydrocarbon plasma processing
b) nanotips c) nanopyramids d) nanowalls

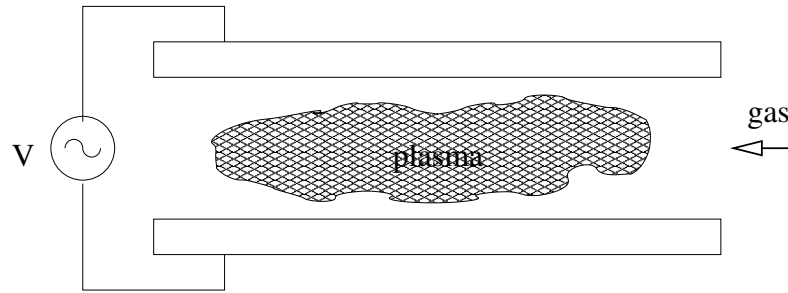
A list of possible applications

- Microelectronics etching, deposition, oxidation, implantation, passivation
- Liquid crystal display and solar cell depositions
- Aerospace and automotive ceramic and metal coatings, films, paints
- Metallurgical melting, refining, welding, cutting, hardening
- Ceramics synthesis, ultrapure powders, nanopowders
- Food packaging permeability barriers
- Textile adhesion treatments
- Medical materials bio-compatibility treatments, sterilization, cleaning
- Architectural and automotive glass coatings
- . . .

Why plasmas?

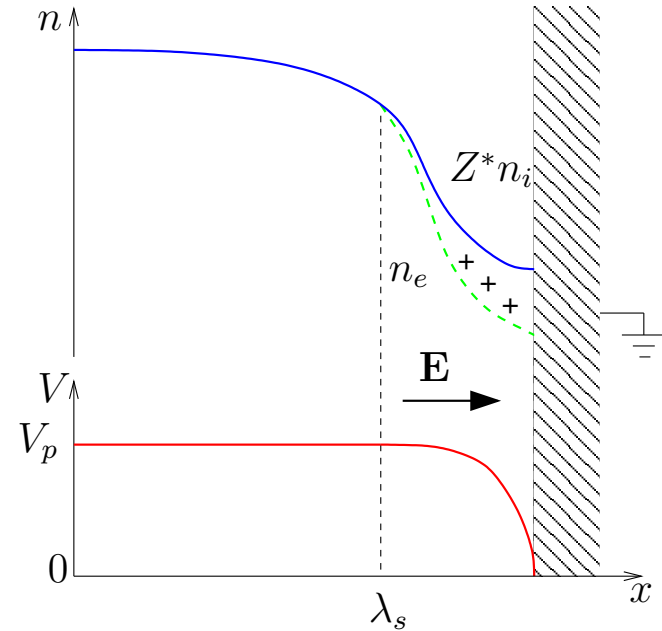
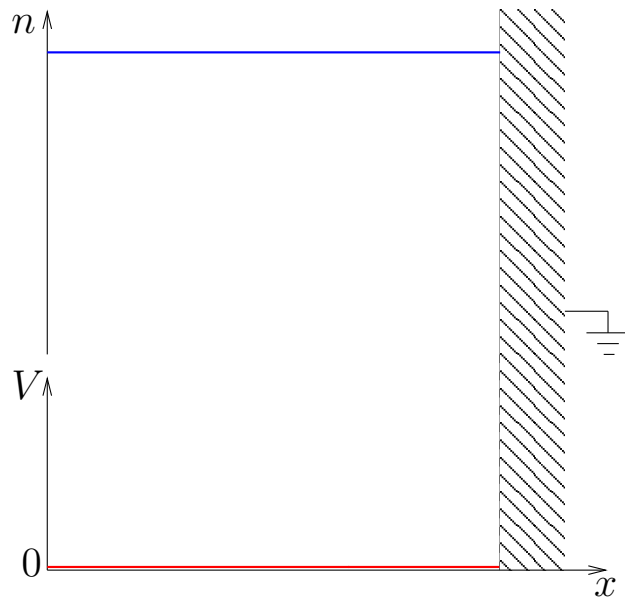
- In a discharge, energetic particles (electrons or ions) drive *chemical reactions with high activation energy* which may be problematic to obtain in other ways
- An important issue is often the *absence of thermal equilibrium*: active species particles may have high energy but the overall energy density is low (example: avoid heating surfaces to high temperatures)
- plasma chemistry is activated and controlled by electric fields; plasma processes are much more **clean** than wet (chemical) ones!

Characterization of discharge plasmas



- Background gas pressure $P \approx 10^{-4} \div 1$ atm
- Free electron density $n_e \approx 10^8 \div 10^{14} \text{ cm}^{-3}$ ($Z^* \ll 1$)
- Electron Temperature $T_e \approx 1 \div 10$ eV
- Device size $L \approx 10 \div 10^2$ cm
- AC frequency: $\omega \approx 10$ Mhz \div 10 Ghz (rf to μ Ws)
- Surface interaction with walls (electrodes, processing material)

Elementary discharge model



A neutral plasma with $n_e = Z^*n_i$ is created between grounded walls at $t = 0$

Electrons drift to the walls, creating a negative *sheath* where ions are accelerated

Ions bombard the walls with energy $\mathcal{E}_i = V_p$ (plasma potential)

Sheath modeling

- Sheath modeling (even in DC and thermal equilibrium conditions) is an old, classical, and still discussed problem of plasma physics ^a
- For plasma *etching* and *nanofabrication*, the desired regime is that of *non-collisional, high voltage* ($eV_s \gg T_e$), DC sheaths, to have well-oriented ion fluxes on the surface
- Depending on the application, modeling of AC sheaths, collisional sheaths, . . . , is desirable

^a[see e.g. Riemann et al, Plasma Phys. Control. Fusion **47**, 1949 (2005)]

Child-Bohm's sheath modeling

Assumptions: $eV_s \gg T_e$, $n_e \ll n_i$

λ_s : sheath thickness $V(x)$: electric potential

J_0 : ion current V_0 : bias potential of the wall

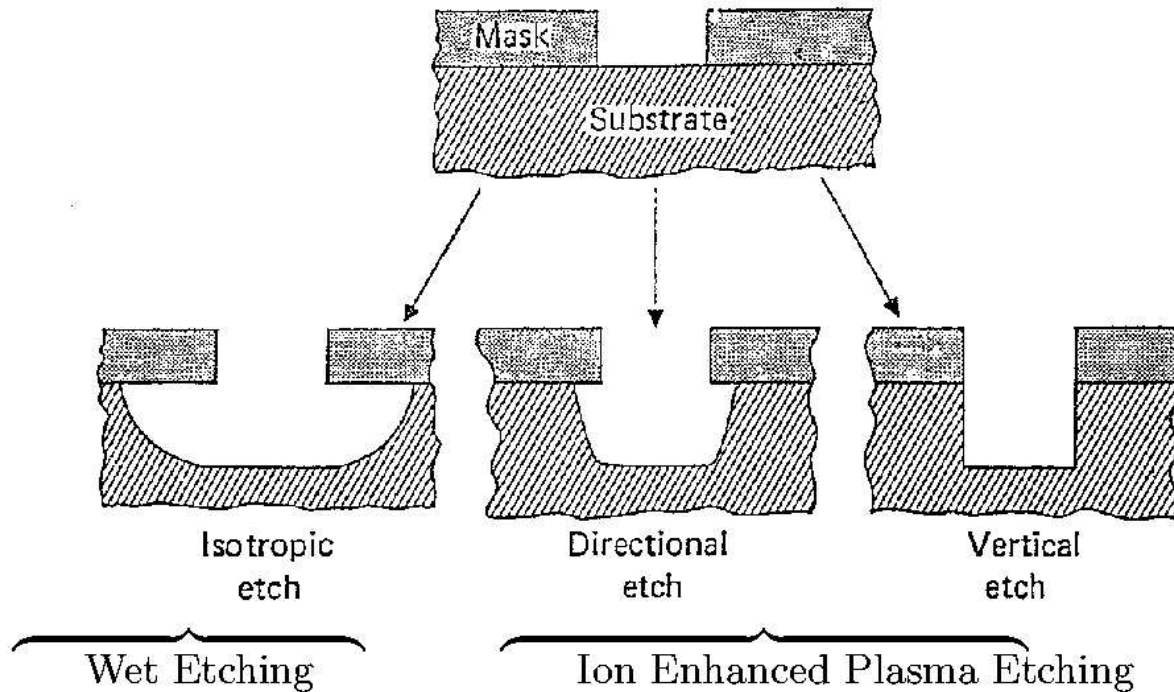
$$\lambda_s = \left(\frac{4\sqrt{2}\epsilon_0}{9en} \right)^{1/2} T_e^{-1/4} V_0^{3/4} \quad V(x) = -V_0 \left(\frac{x}{\lambda_s} \right)^{4/3}$$

$$J_0 = Zen_0 v_b = Zen_0 \sqrt{\frac{ZT_e}{M_i}}$$

See e.g. [1], vol.I, par.9.4.5

Plasma anisotropy is the key to etching

Discharge creates ions → sheath accelerates ions → ions catalyze surface reaction creating volatile specie



Plasma treatment of surfaces

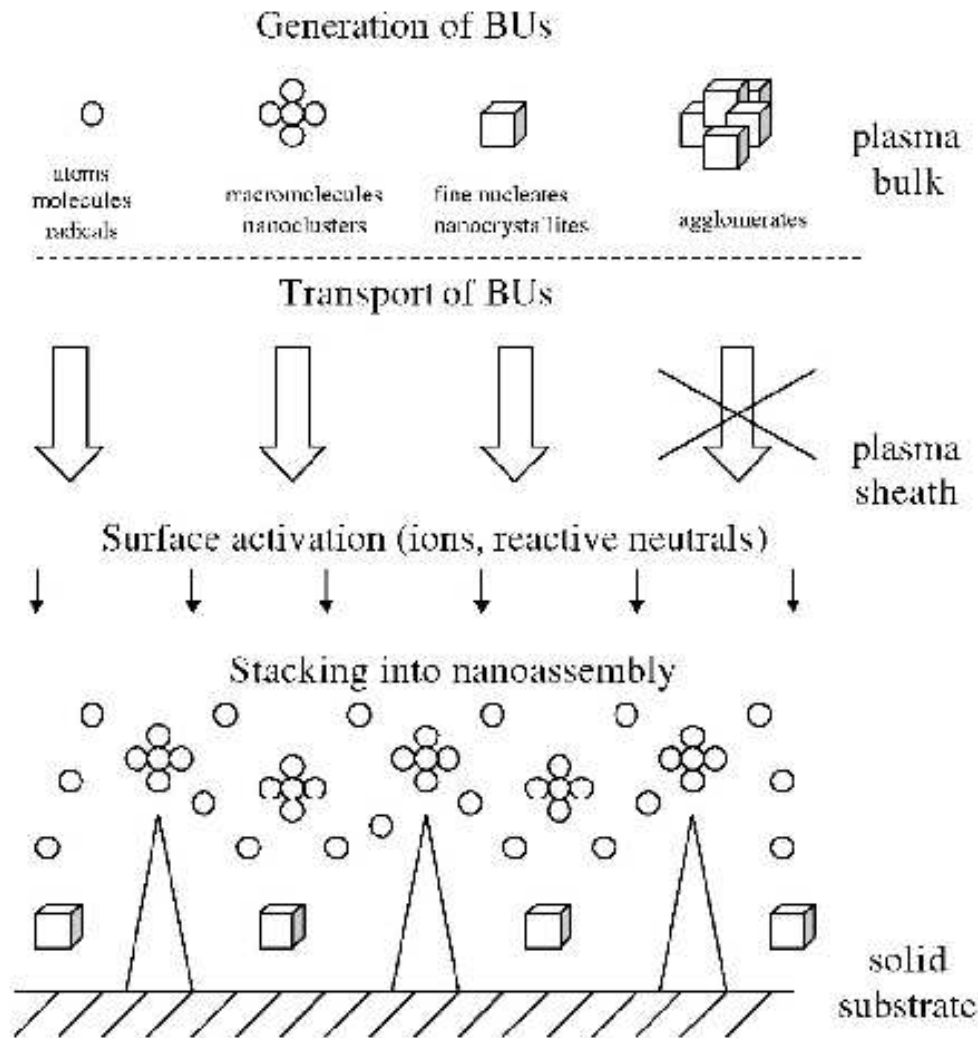
The surface energy may be modified by plasma processing in various ways (depending on exposure time and other parameters)

- plasma *cleaning* (active/passive): removal of surface contaminants
- attachment of polar groups (C=O, OH, COOH) yielding hydrogen bonding
- rotation of bulk polar group by the electric fields at the surface
- surface etching

Plasma deposition

- Plasma Chemical Vapor Deposition (PCVD) is used to produce *thin film* with special properties.
- In a discharge, the feedstock gas is dissociated into molecular fragments and eventually monomers which recombine on the surface of the substrate (*plasma polymerization*).
- The physical details are now well understood, but several “recipes” (chemical composition, plasma parameters, reactor geometry . . .) for thin film fabrication have been found.

Plasma-based nanofabrication



- Building Units (BUs) are
- generated in plasma discharge bulk
 - accelerated (when charged) in plasma sheath
 - transported to the surface
 - sometimes able to drive surface activation (no external heating necessary)

The role of sheath fields

The electric field \mathbf{E} in a collisionless, high-voltage sheath controls

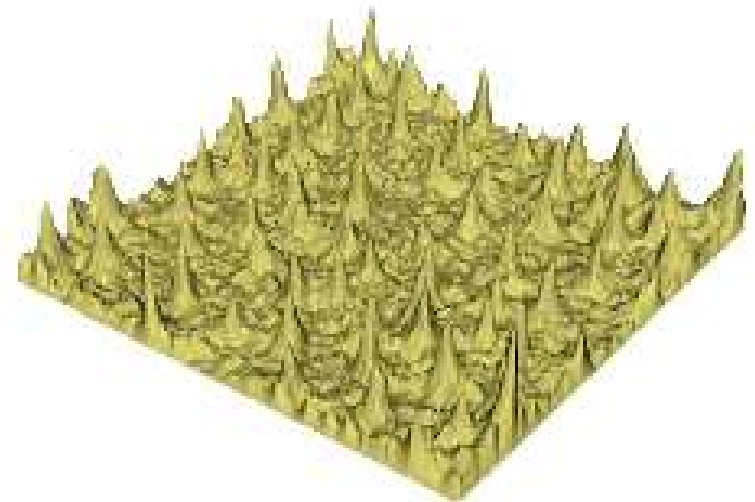
- the flux of ions to the surface
- the landing energy (just that required for activation)
- the orientation (normal to the surface)

“Lucky” effect: lines of \mathbf{E} converge towards sharp ends

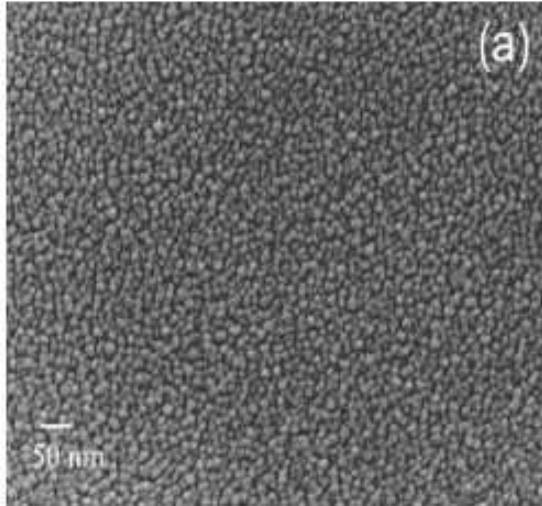
→ **nano-focusing** of BUs

to desired sites

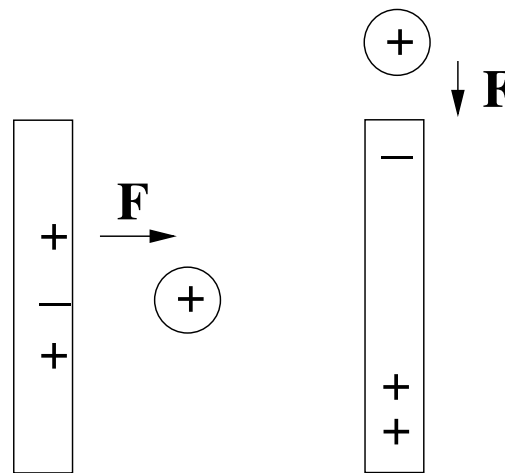
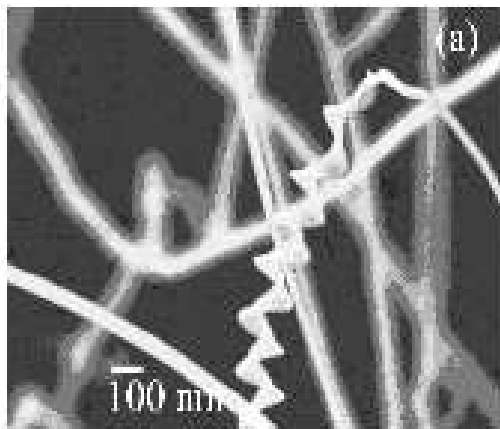
(example: current distribution from simulations of carbon nanotips fabrication) [3]



Other nanofabrication examples



Quantum dots produced by plasma RF sputtering [3]



Silicon nanowires grown by charged nanocluster BUs in a plasma: electrical induction favors one-dimensional growth [3]

Atmospheric plasma discharges

Plasma processing at 1 atm is desirable for industrial applications because

- vacuum operation not needed: **reduce costs**
- continuous (instead of batch) processing possible: **increase productivity**
- higher flux of particles possible: **reduce processing time**

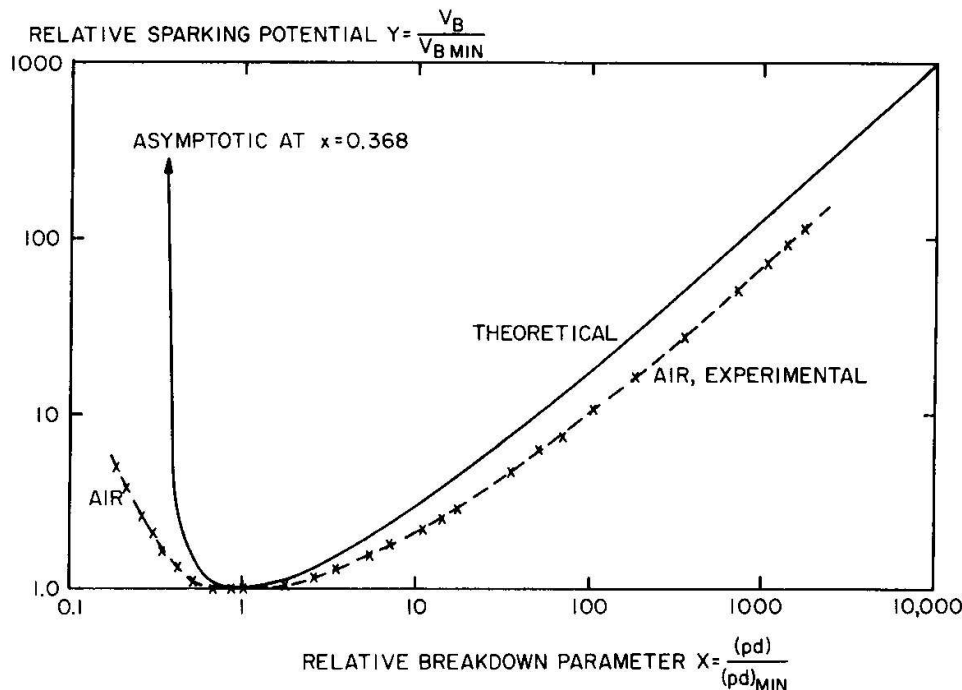
Drawbacks:

- plasma control is more complex
- higher breakdown voltage is required

Paschen law of breakdown

$$V_b = \frac{Cpd}{\ln[Apd / \ln(1 + 1/\gamma)]} = f(pd)$$

V_b : voltage, p : pressure, d : electrode spacing,
 γ : secondary emission coefficient from cathode



Optimal condition for ionizing collisions

$$eEl_c = I_z; \quad E = V_b/d,$$

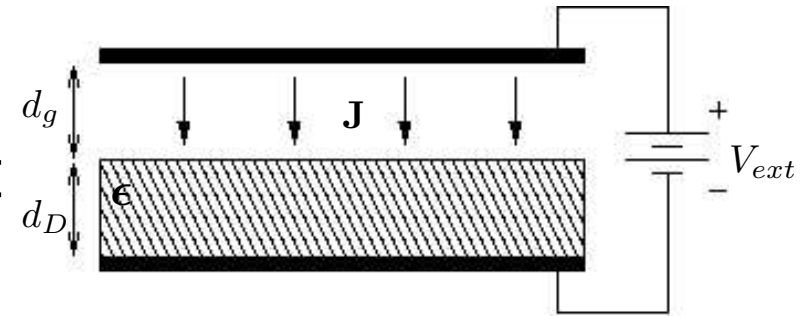
$$l_c \propto n^{-1} \propto p^{-1}$$

Secondary emission provides positive feedback

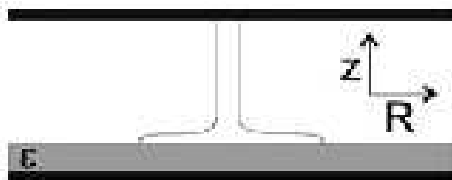
Dielectric barrier discharge (DBD)

A **dielectric** layer (*barrier*) is posed over (both) electrode(s) to:

- provide a non-dissipative current limiter (ballast)
- distribute charges over the surface



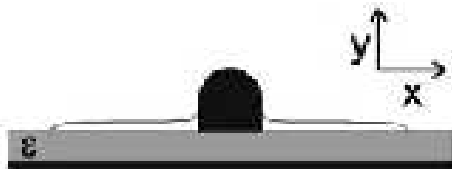
Several DBD schemes exist:



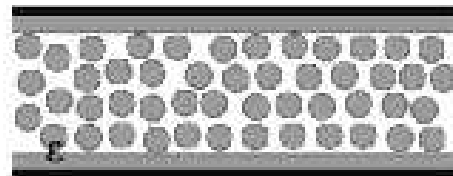
(a)



(b)



(c)



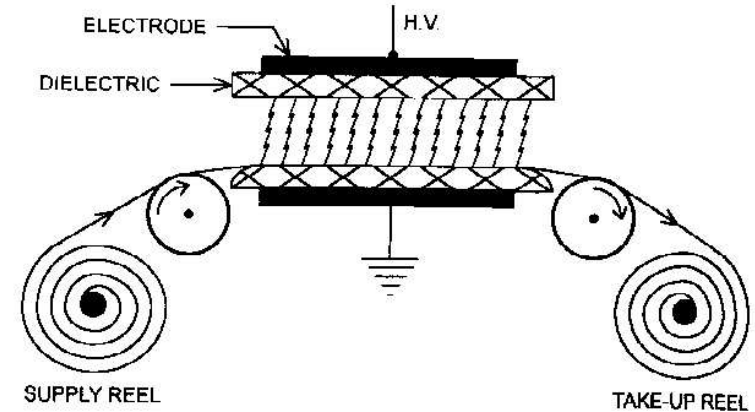
(d)

- a) simple planar cell
- b-c) coplanar electrodes cell
- d) “packed bed” reactor

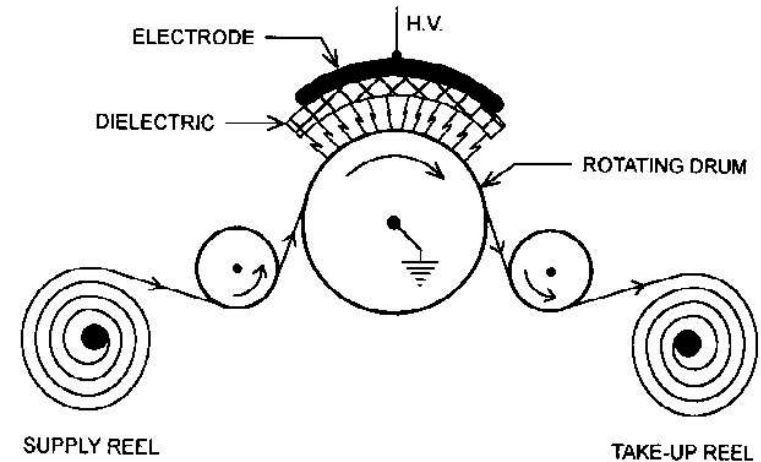
DBD reactors for bulk processing

Parallel plate reactor
for non-adherent webs
processing

Cylindrical annulus reactor
with rotating cylinder
for adherent web processing



a) PLANAR EXPOSURE CONFIGURATION



b) ROTATING CYLINDER DBD CONFIGURATION

DBD phenomenology

The discharge is **filamentary** and short-lived:

- microdischarge initiated by avalanche ionization
- electrons on avalanche head reach the barrier depositing surface charge
- electric field is quenched; discharge stops

Typical microdischarge parameters

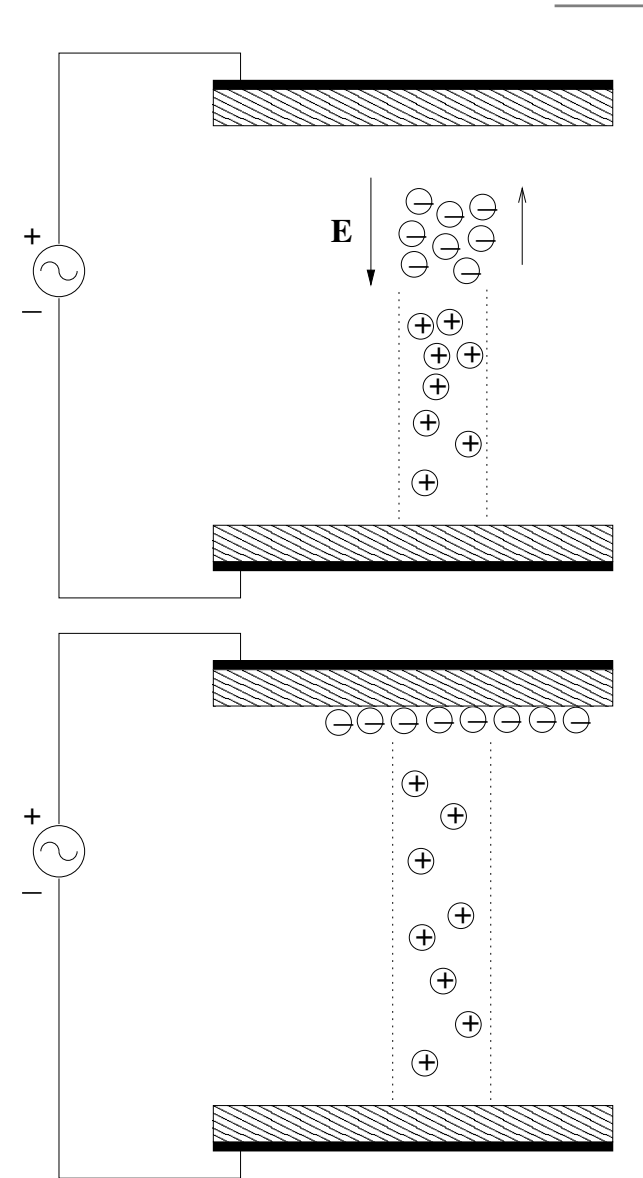
@ $p = 1$ atm:

$$\tau \approx 10 \div 100 \text{ ns} \quad n_e \approx 10^{14} \div 10^{15} \text{ cm}^{-3}$$

$$T_e \approx 1 \div 10 \text{ eV} \quad r_s \approx 100 \text{ } \mu\text{m}$$

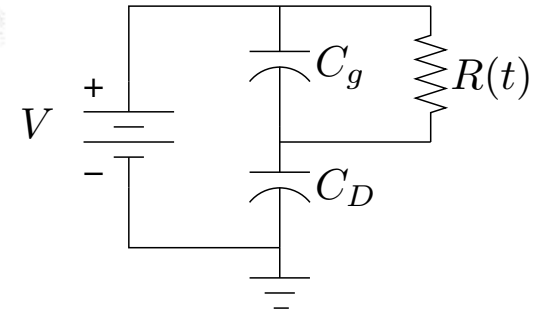
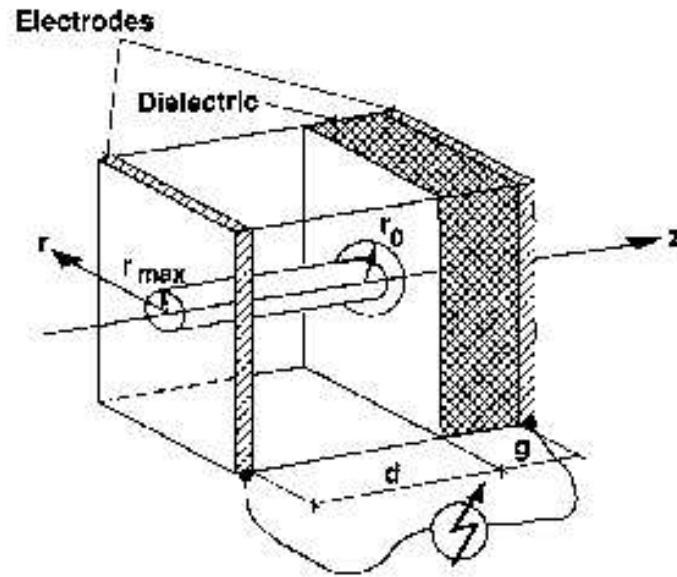
$$Q \approx 1 \text{ pC} \quad U \approx 1 \text{ } \mu\text{J}$$

$$E \approx 10^4 \text{ V cm}^{-1}$$



Simple model of DBD

Consider single filament and an equivalent circuit:



$$S_c = \pi r_s^2, \quad S_d = \pi r_0^2, \quad C_d = \epsilon S_d / d_d, \quad C_g = \epsilon S_d / d_g$$

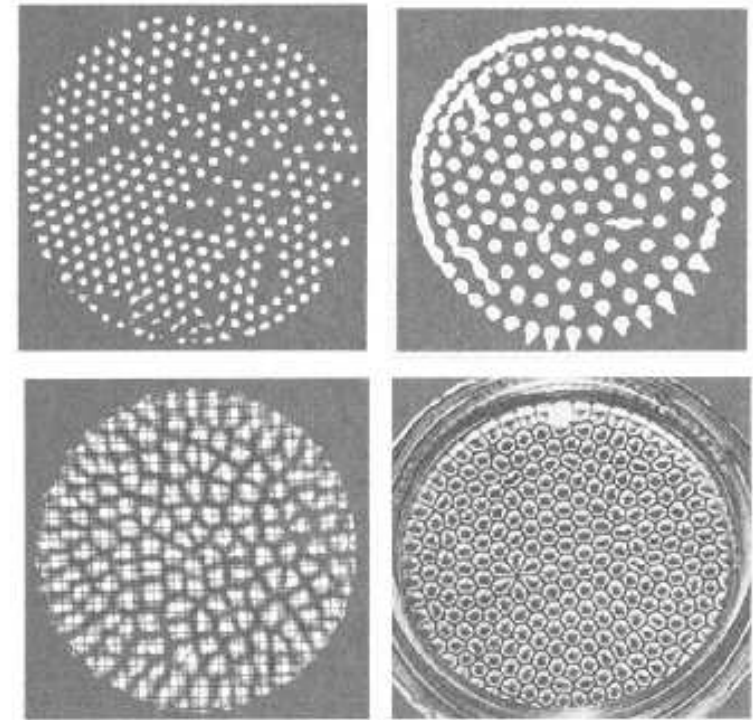
Deposited charge Q and absorbed energy U_{ass} do *not* depend on $R(t) = \rho(t)d_g/S_c$:

$$Q \simeq \epsilon E_b (d_g / d_d) S_{dep}, \quad U_{ass} = \frac{1}{2} \frac{C_d^2}{C_g + C_d} V_{min}^2, \quad V_{min} \simeq E_b d_g$$

$$\text{Duration } \tau \simeq \bar{\rho} \epsilon (d_g / d_d) (S_d / S_c) \simeq E_b (\epsilon / e n_e v_d) (S_d / S_c)$$

Collective effects in DBD

- Space-charge fields drive ionization waves: *streamers*
- Competition of microdischarges for available dielectric surface yields spatial (anti)correlation
- Interaction between microdischarges leads to 2D **pattern formation**: similarity with Bènard convection cells



Atmospheric diffuse (“glow”) discharges

- Issue: transition from filamentary to uniform discharge
- Simple criterion: avalanche electron “heads” must overlap before streamer phase

Spherical head model:

$$r_h \approx \sqrt{z_a \ell_d} \text{ (diffusion), } N_a \approx e^{\alpha z_a},$$

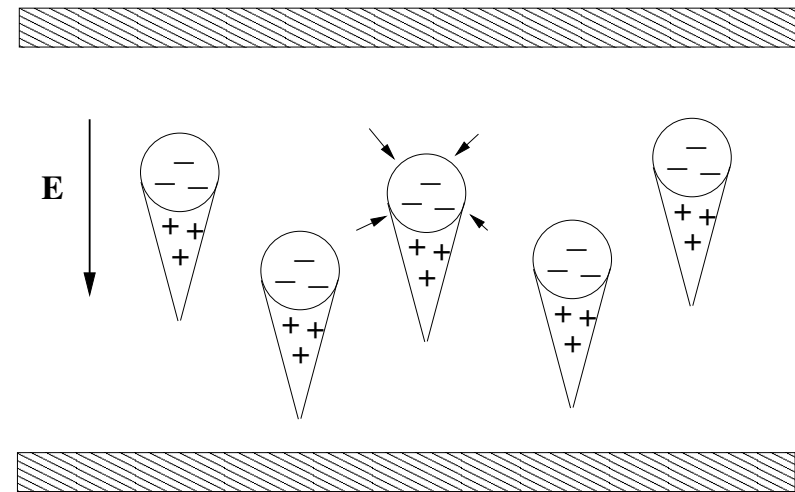
$$E_a \approx \frac{1}{4\pi\epsilon_0} \frac{N_a e}{r_h^2} \doteq E_b \text{ for streamers}$$

→ yields “critical distance” z_{cr}

Heads must overlap:

$$n_{e0} > r_{cr}^{-3} = (\ell_d z_{cr})^{-3/2}$$

n_{e0} : “preionization” density



An example: OAUGDP™

The “One Atmosphere Uniform Glow Discharge Plasma” scheme has been patented and registered by Roth [1].
Concept: there is a frequency range (RF) in which ions are *trapped* between electrodes (planar, symmetric, capacitive DBD configuration) and electrons are not.

$$\bar{x}_{osc,i} \leq d/2 \quad \bar{x}_{osc,e} \geq d/2$$

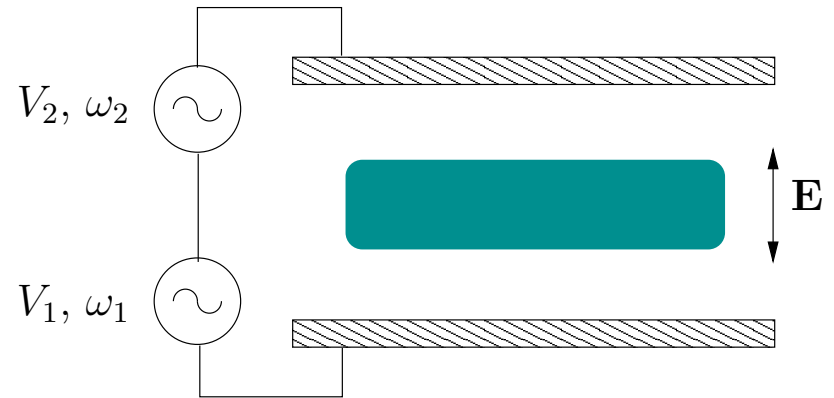
$$\frac{2e\bar{V}}{m_i \nu_{ci} d^2} \leq \omega \leq \frac{2e\bar{V}}{m_e \nu_{ce} d^2}$$

Outside the frequency range: filamentation instabilities
Underlying physics (and chemistry?) needs to be studied
Applications: etching, surface treatment, sterilization, . . .

Dual frequency discharges

Purpose: enhanced control of ion flux for etching (Lieberman)

Plasma density and ion energy are controlled independently



$$n_e \propto P_{abs} \propto \omega^2 V, \quad \mathcal{E}_i \propto V = V_h + V_l$$

$$\omega_h^2 V_h \gg \omega_l^2 V_l \rightarrow \boxed{V_h \text{ controls } n_e}$$

$$V_l \gg V_h \rightarrow \boxed{V_l \text{ controls } \mathcal{E}_i}$$

Detailed scheme involves the physics of *double AC sheath* structure, *stochastic* and *anomalous skin effect* heating, electromagnetic *surface wave* coupling of RF power to the plasma, ...

Conclusions

- Plasma discharges work very well :-) ...
although we do not fully understand why :-(
- Recipes for specific applications can be found
- Current research has lot of issues: empirical application and optimization, enhanced control of plasma parameters, numerical modeling, understanding of basic physics

Let's start the discussion!