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Electroweak Reactions

on Few-Nucleon Systems

- 1. Introduction: Why studying Few-Nucleon Physics?
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- 3. More complicated Few-Nucleon Systems
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1. Introduction:

Why studying Few-Nucleon Physics?

- Principal goal of nuclear physics: understanding of hadronic systems, in particular their internal structure
- Fundamental theory: **QCD**
- Nonperturbative regime: QCD not applicable
 → effective theories based on nucleons, mesons,
 isobars as effective degrees of freedom
- Ideal test object for effective theory: **few-nucleon systems**
 - cleanest theoretical treatment available
 - may serve as effective neutron targets
- Some selected questions:
 - How far in energy and momentum transfer can the effective picture be pushed?
 - Up to which mass number A are **microscopic** calculations feasible?
 - What is the role of relativity?

2. The Two-Nucleon System

Possible reactions in two-nucleon system up to the Δ -region:

 $\begin{array}{lll} \mathrm{NN}\mbox{-scattering:} & NN \rightarrow NN \\ \mathrm{Compton-Scattering:} & \gamma d \rightarrow \gamma d \\ \mathrm{Deuteron \ breakup:} & \gamma d \rightarrow NN \ , & ed \rightarrow e'NN \\ \mathrm{Elastic \ scattering:} & ed \rightarrow e'd \\ \mathrm{Bremsstrahlung:} & NN \rightarrow \gamma NN \\ \mathrm{Pionproduction:} & \pi d \rightarrow \pi d \ , \ \pi d \rightleftharpoons NN \ , \\ & NN \rightarrow \pi NN \ , \ \pi d \rightarrow \pi NN \ , \\ & \gamma d \rightarrow \pi d \ , \ \gamma d \rightarrow \pi NN \ , \ldots \end{array}$

Note: different reactions are linked via unitarity $(\rightarrow \text{Optical Theorem})$, i.e.

Im
$$T(\gamma d \to \gamma d; \theta = 0) \sim \sigma_{tot}(\gamma d \to NN, \pi d, \pi NN, ...)$$

Consequence: Different possible reactions should be treated within **one** framework!

Recently finished (M. S.: habilitation thesis, Universität Mainz):

Approach which is able to describe consistently **all** possible reactions in the two-nucleon system beyond π -threshold up to the Δ -region.

General structure:

• Introduction of XNN-vertices $(X \in \{\pi, \rho, \omega, \sigma, ...\})$ known from potential theory (i.e. Bonn potential) as basic ingredients of the hadronic interaction V



- Relevant equations:
 - Schrödinger equation for bound state:

$$(H_0 + V)|B\rangle = E|B\rangle$$

 Lippmann-Schwinger equation for scattering states:

$$|\Psi\rangle^{\pm} = (1 + G_0(z)T(z))|\phi\rangle^{PW}$$

with

$$T(z) = V + VG_0(z)T(z) , \quad z = E \pm i\epsilon ,$$

with $G_0(z) = \frac{1}{z - H_0}$:
$$\underbrace{\mathbf{T}}_{=} \underbrace{\mathbf{v}}_{=} + \underbrace{\mathbf{v}}_{=} \underbrace{\mathbf{v}}_{=} \underbrace{\mathbf{v}}_{=} + \underbrace{\mathbf{v}}_{=} \underbrace{\mathbf{v$$

 Iteration of V via Lippmann-Schwinger equation generates energy-dependent retarded NN-potential as well as self-energy loop diagrams (→ distinction between bare and physical nucleons necessary):



• Electromagnetic currents are implemented via minimal substitution

 \rightarrow gauge invariance fulfilled



- Note: NN-interaction, internal nucleon structure described by pion cloud, meson exchange currents and e.m. loop diagrams are based on same vertices → unitarity fulfilled!
- Δ -degree taken into account **nonperturbatively**, free parameters are extracted from elastic πN scattering and elementary photopionproduction



• Mutual interactions within the XNN-system are considered using standard three-body techniques



- Applicable for all hadronic and electromagnetic reactions till about 400-500 MeV excitation energy
- πNN -retardation taken into account



Exact, retarded propagator for πNN -state:

$$G_0(z) = (z - E_N(1') - E_N(2) - E_\pi)^{-1}, \quad z = E + i\epsilon$$

Characteristic features:

- Nonlocal
- Nonhermitean
- Existence of singularities above pion-threshold

 \rightarrow low energy approximation: **static limit**, i. e. nucleons are infinitely heavy during meson exchange:

$$G_0(z) \to G_0^{stat} := -E_\pi^{-1}$$

- Advantage: local, energy independent
- Disadvantages: violation of unitarity, pion is no longer dynamic degree of freedom



Extension to Electrodisintegration

unpolarized differential cross section in d(e, e'p)n:

 $\frac{d^{3}\sigma}{dk_{2}^{lab}d\Omega_{e}^{Lab}d\Omega_{np}} = c\left[\rho_{L}f_{L}(\omega,q,\theta) + \rho_{T}f_{T}(\omega,q,\theta) + \right]$

 $\rho_{LT} f_{LT}(\omega, q, \theta) \cos(\phi) + \rho_{TT} f_{TT}(\omega, q, \theta) \cos(2\phi)]$

NIKHEF-Experiment: $E_{np} = 280 \text{ MeV}, q^2 = 2.47 \text{ fm}^{-2}$



Comparison to photodisintegration at $E_{\gamma} = 300 \text{ MeV}$



full: retarded approach (2006) data: • Daphne (1996/1999), \circ LEGS (1995)







3. More complicated Few-Nucleon Systems

Interesting questions:

- Role of ³He as an effective neutron target
- What is in general the role of three (four,..)-body forces?
- Miscroscopic understanding of **nuclear structure**, i.e.
 - Up to which mass number A are **microscopic** calculations feasible?
 - In the long-term run: Test of typical many-body approximations by building a bridge between classical few-body systems (A = 2, 3) and complex nuclei $(A \ge 12)$
- astrophysical applications

• ...

• Most important observable: inclusive Response function for $\gamma A \to X$

$$R(W) = \int \Psi_f |\langle \Psi_f | \epsilon_\mu J^\mu | \Psi \rangle|^2 \delta(E_f - E_0 - W)$$

with

$$R(W) \sim \sigma_{tot}(\gamma A \to X, W)$$

• Optical Theorem:

 $\operatorname{Im} \lim_{\epsilon \to 0} T(\gamma A \to \gamma A; W + i\epsilon, \theta = 0) \sim R(W)$



• Problem: Pole structure of amplitude becomes more and more involved with increasing A! Elegant solution: Lorentz Integral
Transformation (Efros, Leidemann & Orlandini)

$$L(\sigma_r, \sigma_i) = \int_{E_{thres}}^{\infty} dW \frac{R(W)}{(W - \sigma_r)^2 + \sigma_i^2}$$

~ Im $T(\gamma A \to \gamma A; \sigma_r + i\sigma_i, \theta = 0)$

compare: $\sigma_{tot}(W) \sim \operatorname{Im} \lim_{\epsilon \to 0} T(\gamma A \to \gamma A; W + i\epsilon, \theta = 0)$

• Advantages:

- Due to $\sigma_i \neq 0$ fixed (typically 10-20 MeV), no pole structure in $T(\gamma A \rightarrow \gamma A)$

\longrightarrow bound state problem!

- FSI automatically included
- microscopic calculation of electromagnetic reactions on "complex" few-nucleon systems feasible (presently up to A = 7), see also contribution by Sonia Bacca
- Applicable also for exclusive cross reactions
- Price to be paid: **inversion** $L \to R$ necessary

Characteristic features of presently available LITcalculations:

- Realistic hadronic interaction (i.e. AV18, Urbana IX) used, at least for A = 3, 4
- In general, calculations are performed in configuration space, taking only **nucleonic** degrees of freedom into account

Planned Extension:

Application of LIT-method in **momentum** space, taking into account Δ and mesonic degrees of freedom

Suitable starting point: ${}^{3}\text{He}$

- standard Faddeev-calculations only applicable till $\pi\text{-threshold}$
- \rightarrow extension planned to higher energy/momentum transfer:
 - Inclusive cross section in the Δ -region
 - Extraction of G_{E_n} at large Q^2
 - GDH sum rule on the neutron using 3 He as effective neutron target:

$$I^{GDH} = \int_0^\infty \frac{dk}{k} \left(\sigma^P(k) - \sigma^A(k) \right) = 4\pi^2 \kappa^2 \frac{e^2}{m^2} S$$

4. Summary and Outlook

- Few-nucleon systems important for understanding of strong interaction, nuclear structure and extraction of neutron properties → very active field
- Self-consistent approach up to resonance region available for A = 2, to be extended for A = 3 (presently under construction)
- Some recent theoretical progress:
 - Role of retardation clarified
 - Lorentz integral method
 - ...
- Experimental challenge: to provide data in order to test theory; various activities at TJNAF, MAMI, TUNL,...