

# Testing superscaling on nuclei at low $q$ with different nuclear models and predictions on electroweak cross section

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# Topics

- \* Scaling Procedure

- \* Nuclear Models

- RFG
- SM
- RPA
- MEC
- FSI

- \* Electron and Neutrino Cross Sections

- Scaling approach  $\iff$  full calculations
- Comparison with experiments (electron)
- Prediction on cross sections (neutrino)

# Scaling

## General procedure

\*Divide the experimental cross section (response) by an appropriate single-nucleon cross section (response), having contributions from  $Z$  protons and  $N$  neutrons with their corresponding form factors, to obtain a reduced cross section (response).

\*Plot the reduced cross section (response) versus one or more appropriately chosen variables: if the results do not depend on some of these variables and a universal behavior is found one says that **the results scale**.

Scaling of the **first kind**  $\iff$  Independence of the momentum transfer  $q$

Scaling of the **second kind**  $\iff$  Independence of  $k_F$  or  $A$

**Superscaling**  $\iff$  Both 1st- and 2nd- kind scaling

Scaling of **zeroth- kind**  $\iff$  Independence of type of **external operator**

# Relativistic Fermi Gas Model

**Universal Scaling Function**  $f^{RFG}(\Psi) = \frac{3}{4} (1 - \Psi^2) \theta(1 - \Psi^2)$

**Scaling variable**  $\Psi = \Psi_0 \left( 1 + \Psi_0 \frac{k_F}{2q} \sqrt{\frac{q^2}{m^2} + 1} \right)$

$$\Psi_0 = \frac{2m}{k_F} \left[ \sqrt{\left( \frac{\omega - E_{shift}}{2m} \right) \left( 1 + \frac{\omega - E_{shift}}{2m} \right)} - \frac{q}{2m} \right]$$

Some effects not considered in the RFG model:

- \*Finite size of the system
- \*Collective excitations
- \*Meson Exchange Currents (MEC)
- \*Final State Interactions (FSI)

**Scaling ?**

# Electron scattering

$$\frac{d^2\sigma}{d\theta d\omega} = \sigma_M \left\{ \frac{(\omega^2 - q^2)^2}{q^4} R_L(\omega, q) + \left[ \tan^2\left(\frac{\theta}{2}\right) - \frac{\omega^2 - q^2}{2q^2} \right] R_T(\omega, q) \right\}$$

**Longitudinal Response**  $R_L(\omega, q) = 4\pi \sum_{j=0} |\langle J_f || \mathcal{C}_J || J_i \rangle|^2$

**Transverse Response**  $R_T(\omega, q) = 4\pi \sum_{j=1} (|\langle J_f || \mathcal{E}_J || J_i \rangle|^2 + |\langle J_f || \mathcal{M}_J || J_i \rangle|^2)$

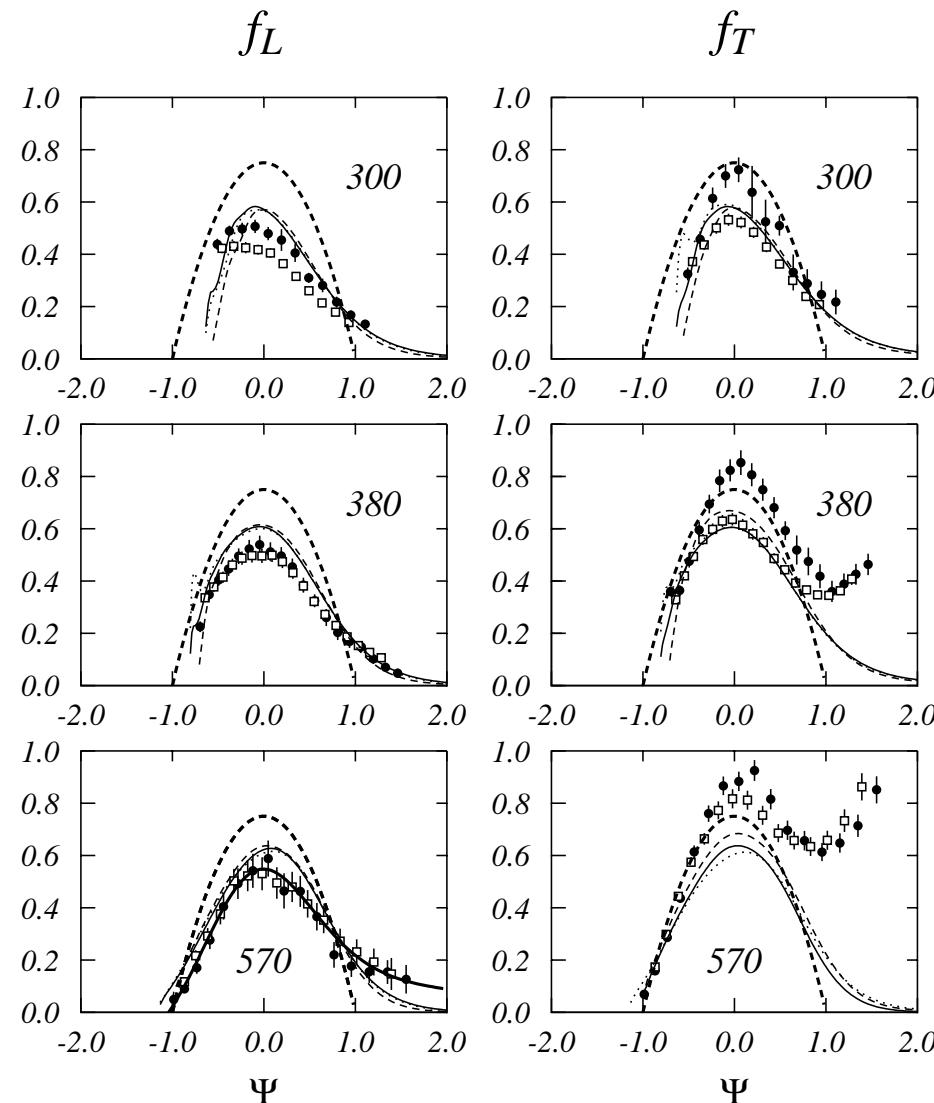
## Scaling functions

$$f_L(\Psi) = k_F \frac{q^2 - \omega^2}{q m} \frac{R_L(\omega, q)}{Z(G_E^p)^2 + N(G_E^n)^2}$$

$$f_T(\Psi) = 2 k_F \frac{q m}{q^2 - \omega^2} \frac{R_T(\omega, q)}{Z(G_M^p)^2 + N(G_M^n)^2}$$

## Semi-relativistic corrections

$$\epsilon_p \rightarrow \epsilon_p \left( 1 + \frac{\epsilon_p}{2m} \right); R_L(q, \omega) \rightarrow \frac{q^2}{q^2 - \omega^2} R_L(q, \omega); R_T(q, \omega) \rightarrow \frac{q^2 - \omega^2}{q^2} R_T(q, \omega)$$



Full circles:  $^{12}\text{C}$ ; Empty squares:  $^{40}\text{Ca}$ ; — - RFG; — Empirical Scaling Function;  
 —  $^{12}\text{C}$ , ...  $^{16}\text{O}$  and - -  $^{40}\text{Ca}$  with all the effects beyond RFG

# Superscaling beyond RFG model

Numerical comparison of the various scaling functions

$M$  scaling functions  $f_\alpha$ ;  $\alpha = 1, M$

each of them known for

$N$  values of the scaling variable  $\Psi_i$ ;  $i = 1, N$

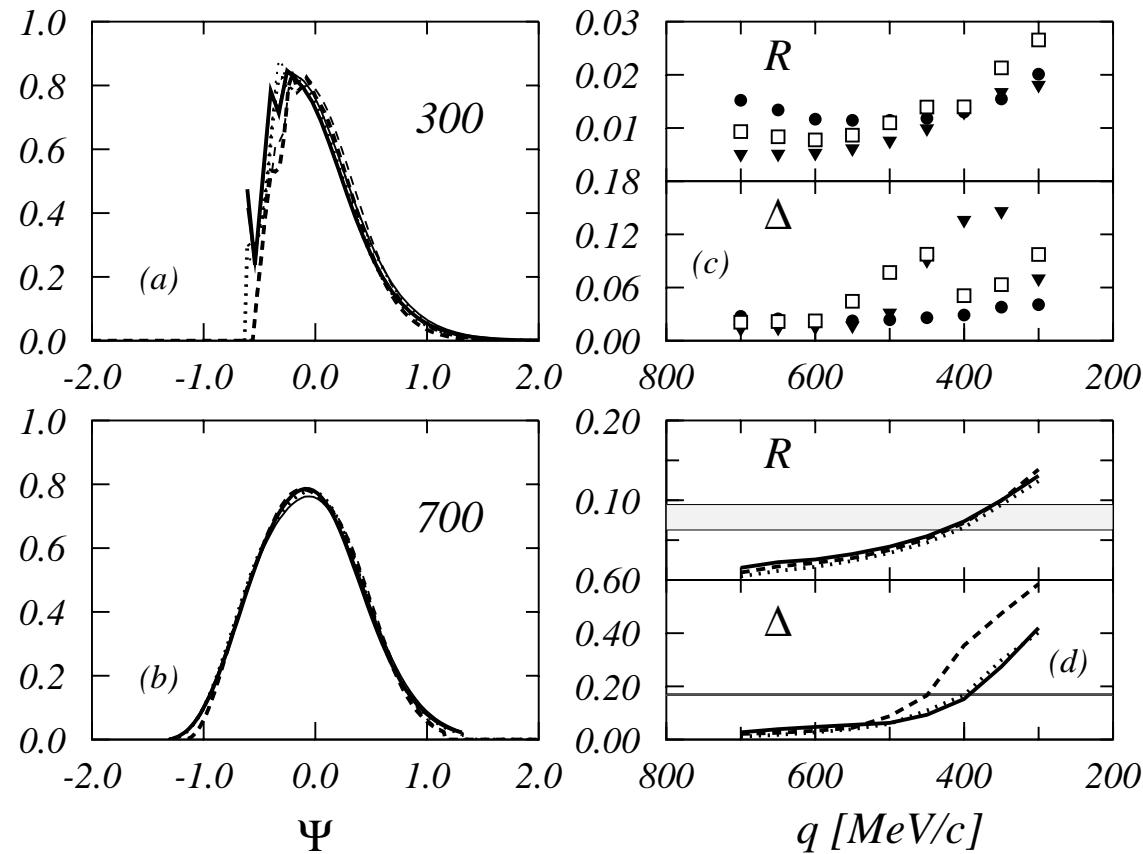
$$\Delta = \max_{i=1,N} [f^{max}(\Psi_i) - f^{min}(\Psi_i)] \text{ maximum distance}$$

$$\mathcal{R} = \frac{1}{Nf^M} \sum_{i=1,N} [f^{max}(\Psi_i) - f^{min}(\Psi_i)] \text{ global properties}$$

\* $\mathcal{R}=0, \Delta=0 \iff$  Perfect Scaling - Only in the RFG

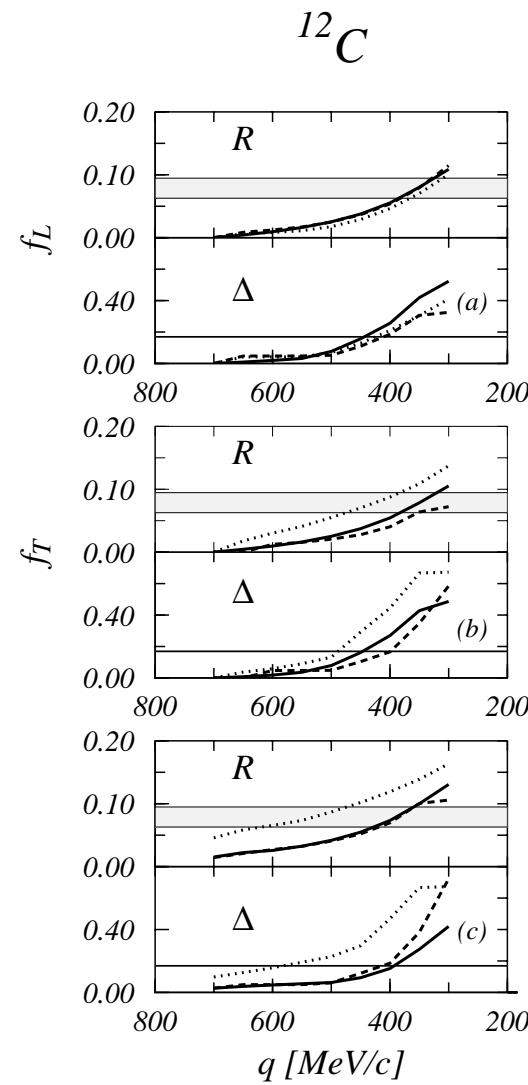
\* $\mathcal{R}=0.078 \pm 0.016, \Delta = 0.169 \pm 0.002 \iff$  2nd-kind scaling for  $f_L$  data at 570 MeV/c

# Continuum shell model result

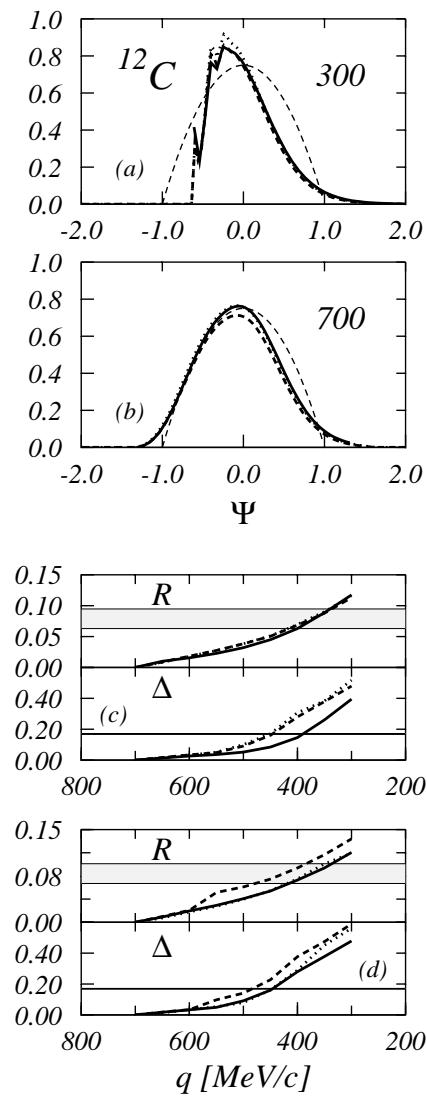


—  $^{12}\text{C}, \dots, ^{16}\text{O}$  and - -  $^{40}\text{Ca}$ ; thicker lines:  $f_L$ , thinner lines:  $f_T$   
 circles  $^{12}\text{C}$ , squares  $^{16}\text{O}$  and triangles  $^{40}\text{Ca}$ ; horizontal bands: empirical  $\mathcal{R}$ ,  $\Delta$

# RPA



# MEC



# Folding

Phenomenological model used to describe hadronic processes and electron scattering cross section in order to take into account effects beyond the RPA

G. Co', K.Q. Quader, R. Smith, J. Wambach, NPA 485 (1988) 61

J.E. Amaro, G. Co', A.M. Lallena, NPA 578 (1994) 365

**known nuclear response**  $S^0(q, E)$

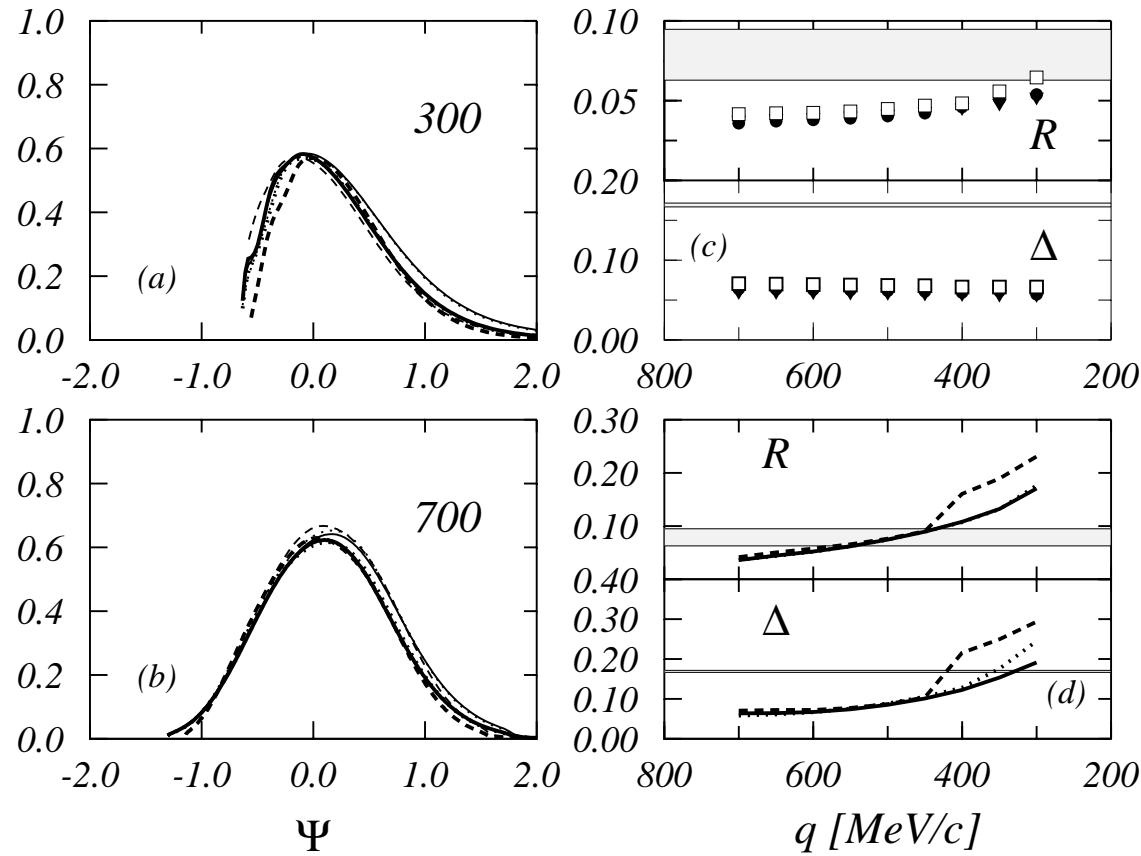
**folding integral**  $S^{\text{FSI}}(q, \omega) = \int_0^\infty dE S^0(q, E)[h(E, \omega) + h(E, -\omega)]$

**folding function**  $h(E, \omega) = \frac{1}{2\pi} \frac{\Gamma(\omega)}{[E - \omega - \Delta(\omega)]^2 + [\Gamma(\omega)/2]^2}$

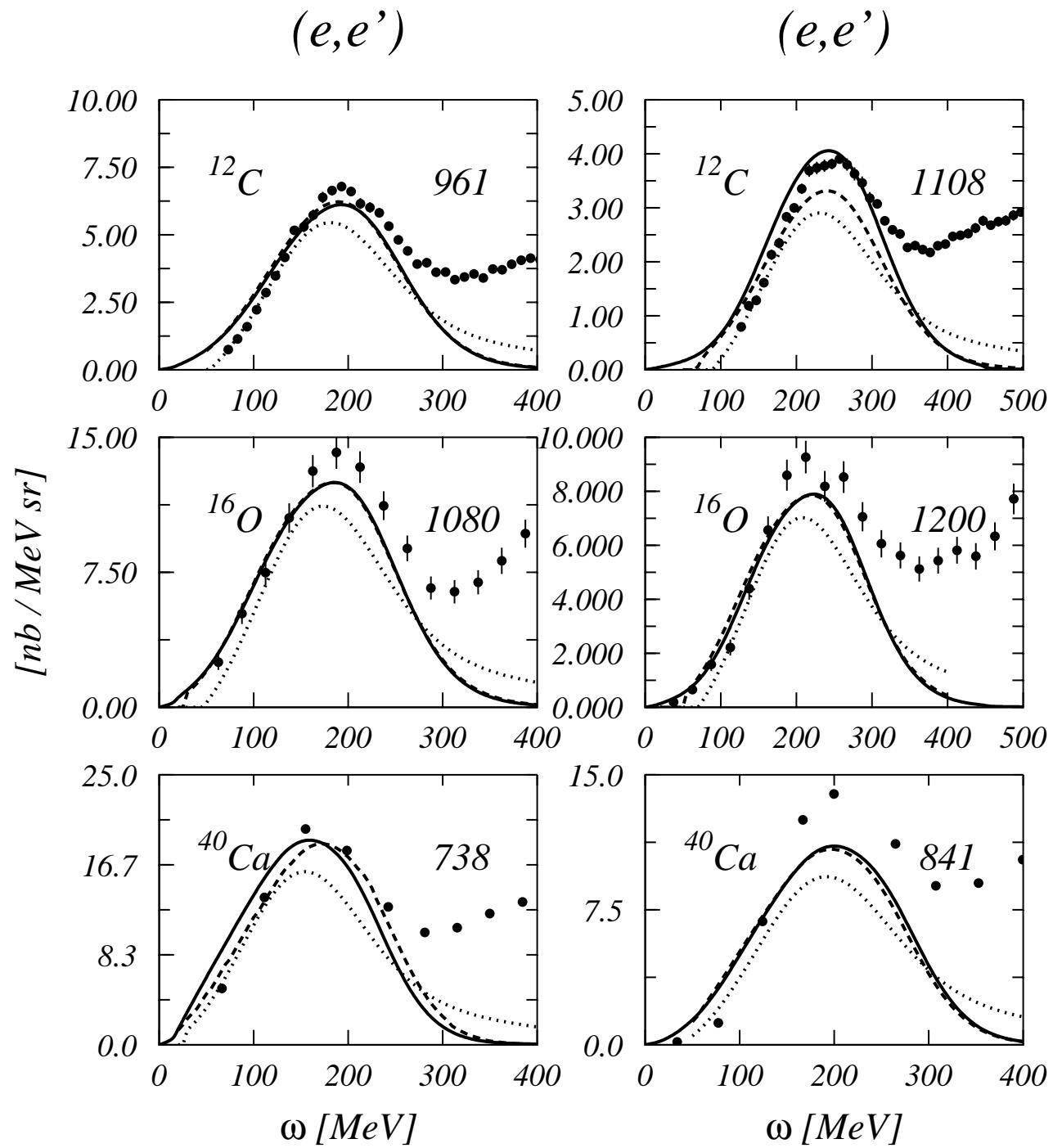
$$\Gamma(\omega) = \frac{1}{\omega} \int_0^\omega d\epsilon [\gamma(\epsilon + \omega) + \gamma(\epsilon - \omega)]$$

**single particle width**  $\gamma(\epsilon) = A \left( \frac{\epsilon^2}{\epsilon^2 + B^2} \right) \left( \frac{C^2}{\epsilon^2 + C^2} \right)$

# Scaling functions with the FSI



—  $^{12}\text{C}, \dots, ^{16}\text{O}$  and - -  $^{40}\text{Ca}$ ; thicker lines:  $f_L$ , thinner lines:  $f_T$   
circles  $^{12}\text{C}$ , squares  $^{16}\text{O}$  and triangles  $^{40}\text{Ca}$ ; horizontal bands: empirical  $\mathcal{R}, \Delta$



## Inclusive electron scattering cross section - Figure Caption

The  $^{12}\text{C}$  data, measured at  $37.5^\circ$ , are from R.M. Sealock et al. Phys. Rev. Lett. 62, 1350 (1989), those of  $^{16}\text{O}$ , measured at  $32.0^\circ$ , are from M. Anguiano, A.M. Lallena and G. Co', Phys. Rev. C 53, 3155 (1996) and those of  $^{40}\text{Ca}$ , measured at  $45.5^\circ$ , from C. Williamson et al., Phys. Rev. C 56, 3152 (1997). The full lines show the results of the complete calculations. The cross sections obtained with our  $f_L$  are shown by the dashed lines, and those obtained with the empirical scaling function are given by the dotted lines.

# $(\nu_e, e^-)$ Reaction

## Differential Cross Section

$$\begin{aligned}\frac{d^2\sigma}{d\Omega d\omega} = & \frac{G^2 \cos^2 \theta_C}{(2\pi)^2} |\mathbf{k}'| \epsilon' F(Z', \epsilon') \left\{ \left( l_0 l_0^* + \frac{\omega^2}{q^2} l_3 l_3^* - \frac{\omega}{q} l_3 l_0^* \right) R_{CC}^V(\omega, q) \right. \\ & + l_0 l_0^* R_{CC}^A(\omega, q) + l_3 l_3^* R_{LL}^A(\omega, q) + 2 l_3 l_0^* R_{CL}^A(\omega, q) \\ & \left. + \left[ \frac{1}{2} (\vec{l} \cdot \vec{l}^* - l_3 l_3^*) \right] [R_T^V(\omega, q) + R_T^A(\omega, q)] + 2 \left[ -\frac{i}{2} (\vec{l} \times \vec{l}^*)_3 \right] R_{T'}(\omega, q) \right\}\end{aligned}$$

Vector ( $V$ ) and Axial-vector ( $A$ ) contributions  
in the Nuclear Response Functions

$$R_{CC}^V(\omega,q)=4\pi\sum_{j=0}|\langle J_f||\mathcal{C}_J^V||J_i\rangle|^2 \hspace*{0.2cm} ; \hspace*{0.2cm} R_{CC}^A(\omega,q)=4\pi\sum_{j=0}|\langle J_f||\mathcal{C}_J^A||J_i\rangle|^2$$

$$R_{CL}^A(\omega,q)=2\pi\sum_{j=0}(\langle J_f||\mathcal{C}_J^A||J_i\rangle^\star\langle J_f||\mathcal{L}_J^A||J_i\rangle+\langle J_f||\mathcal{C}_J^A||J_i\rangle\langle J_f||\mathcal{L}_J^A||J_i\rangle^\star)$$

$$R_{LL}^A(\omega,q)=4\pi\sum_{j=0}|\langle J_f||\mathcal{L}_J^A||J_i\rangle|^2$$

$$R_T^V(\omega,q)=4\pi\sum_{j=1}(|\langle J_f||\mathcal{E}_J^V||J_i\rangle|^2+|\langle J_f||\mathcal{M}_J^V||J_i\rangle|^2)$$

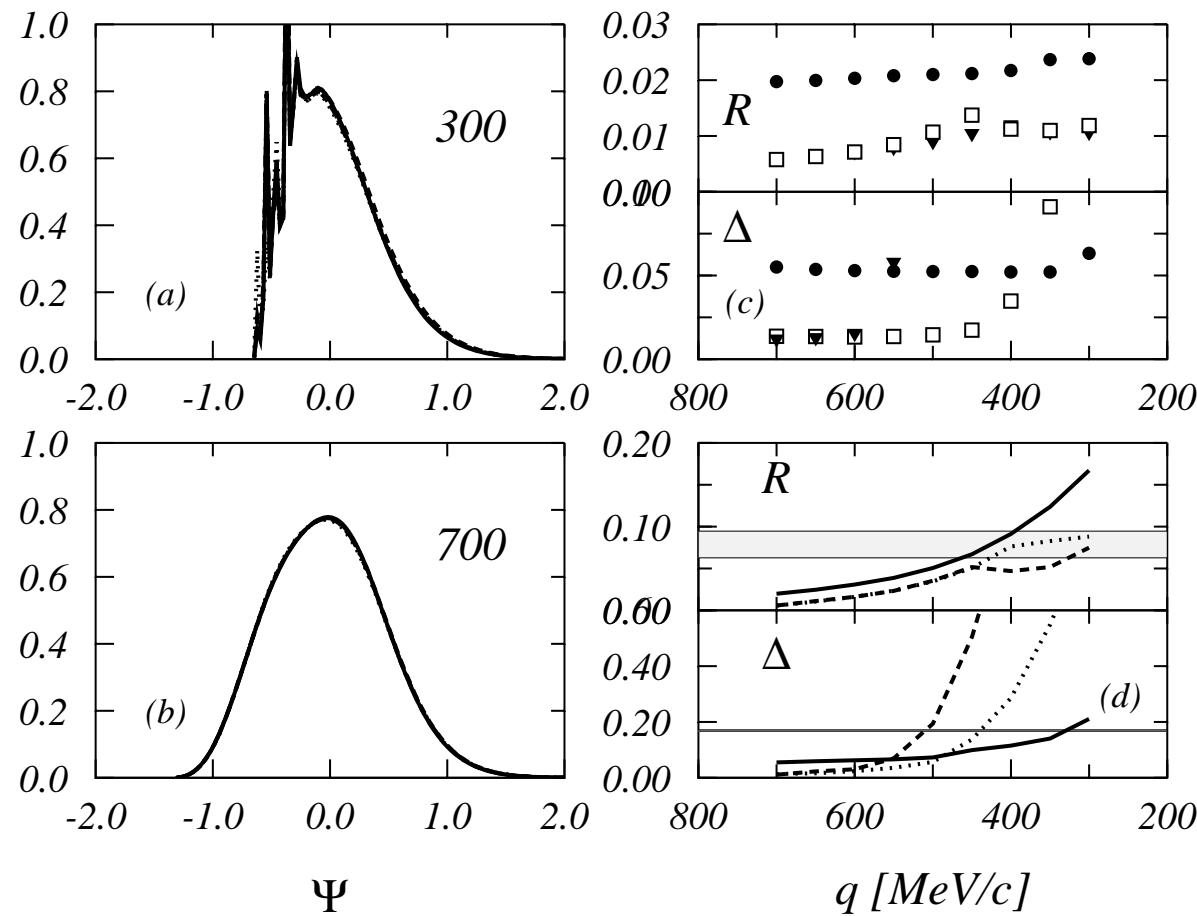
$$R_T^A(\omega,q)=4\pi\sum_{j=1}(|\langle J_f||\mathcal{E}_J^A||J_i\rangle|^2+|\langle J_f||\mathcal{M}_J^A||J_i\rangle|^2)$$

$$\begin{aligned} R_{T'}^{V,A}(\omega,q)=2\pi\sum_{j=1}&(\langle J_f||\mathcal{E}_J^V||J_i\rangle^\star\langle J_f||\mathcal{M}_J^A||J_i\rangle+\langle J_f||\mathcal{E}_J^V||J_i\rangle\langle J_f||\mathcal{M}_J^A||J_i\rangle^\star+\\&\langle J_f||\mathcal{E}_J^A||J_i\rangle^\star\langle J_f||\mathcal{M}_J^V||J_i\rangle+\langle J_f||\mathcal{E}_J^A||J_i\rangle\langle J_f||\mathcal{M}_J^V||J_i\rangle^\star) \end{aligned}$$

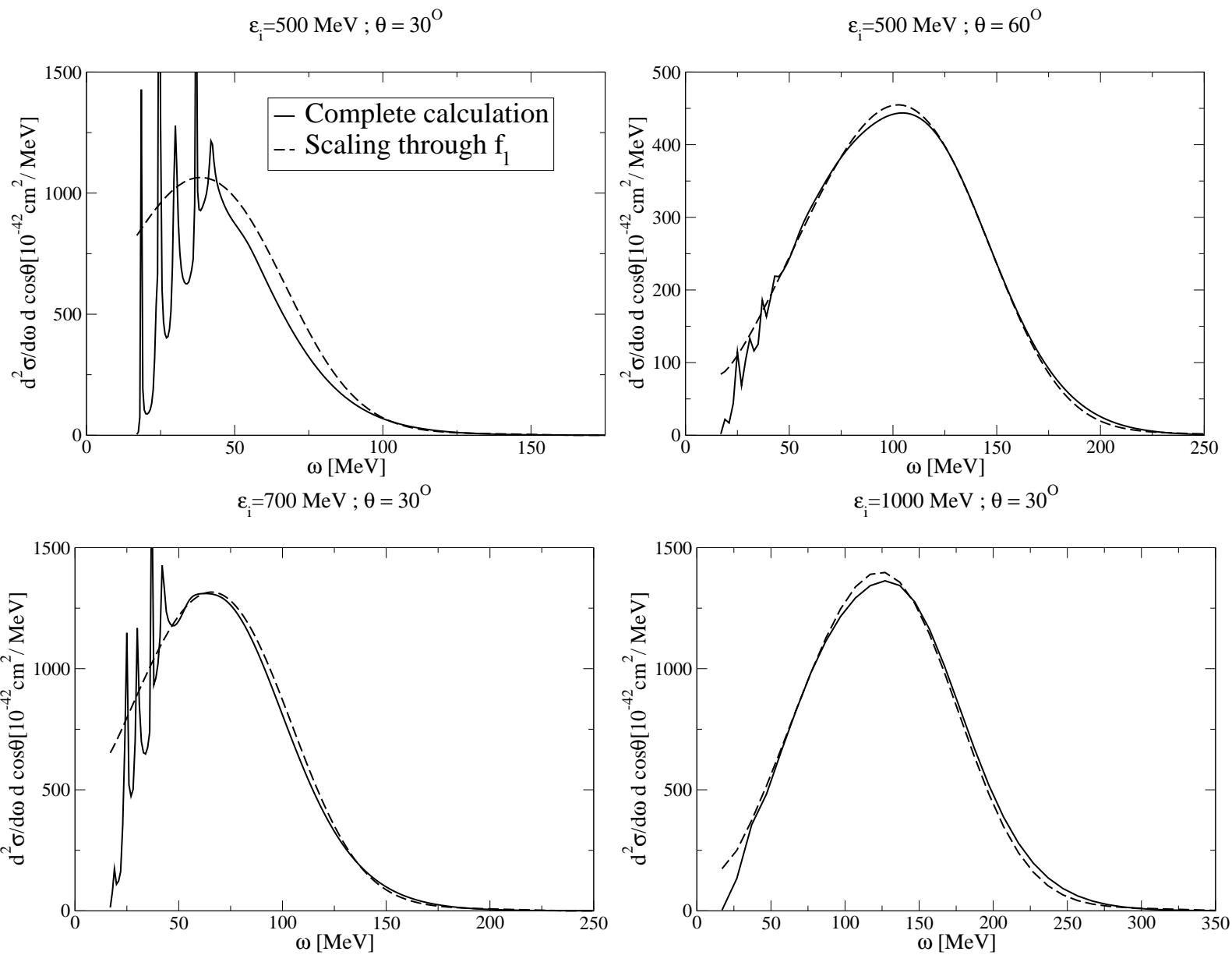
# Scaling Functions

$$\begin{aligned} f_{CC}^V(\Psi) &= k_F \frac{q^2 - \omega^2}{q m_N} \frac{R_{CC}^V(\omega, q)}{N(G_E^{(1)})^2} \\ f_{LL}^A(\Psi) &= 4 k_F \frac{q m_N}{4m_N^2 + q^2 - \omega^2} \frac{R_{LL}^A(\omega, q)}{N(G_A)^2} \\ f_T^V(\Psi) &= 2 k_F \frac{q m_N}{q^2 - \omega^2} \frac{R_T^V(\omega, q)}{N(G_M^{(1)})^2} \\ f_T^A(\Psi) &= 2 k_F \frac{q m_N}{4m_N^2 + q^2 - \omega^2} \frac{R_T^A(\omega, q)}{N(G_A)^2} \\ f_{T'}^{VA}(\Psi) &= 2 k_F \frac{q m_N}{\sqrt{q^2 - \omega^2} \sqrt{4m_N^2 + q^2 - \omega^2}} \frac{R_{T'}^{VA}(\omega, q)}{N G_M^{(1)} G_A} \end{aligned}$$

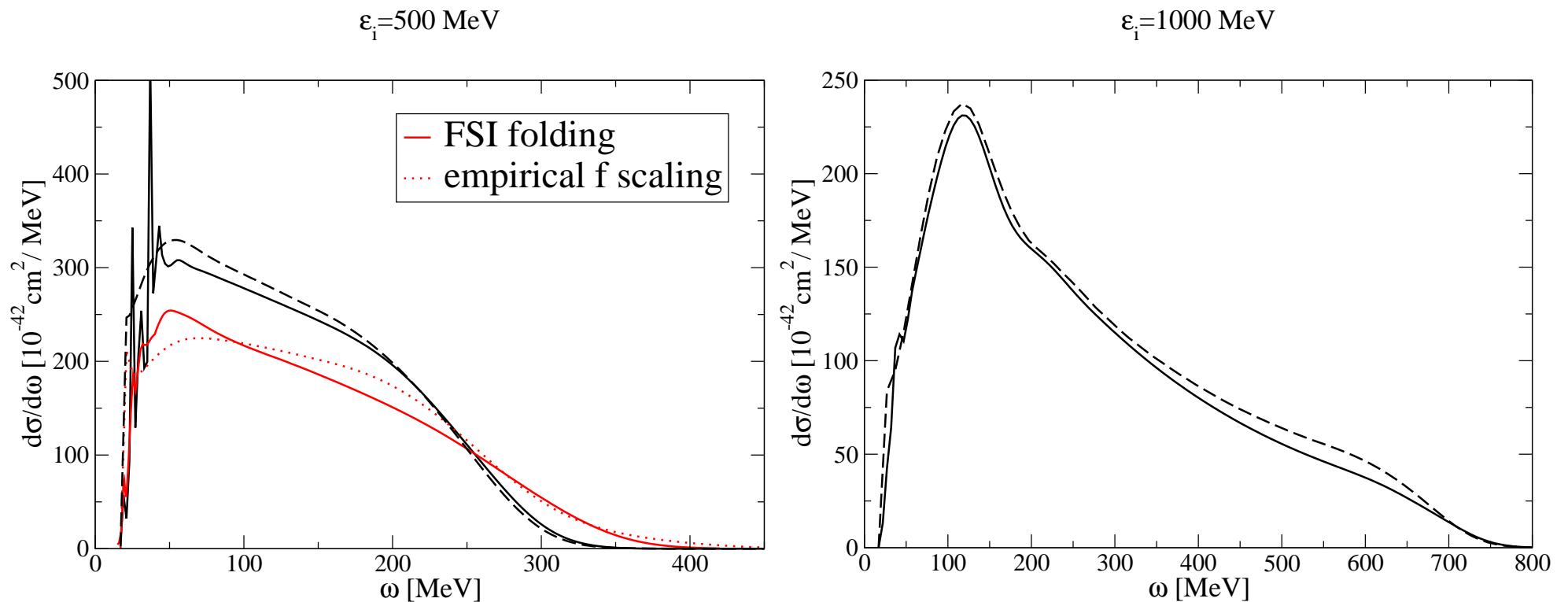
# Scaling analysis



(a), (b):  $f_{CC}^V f_{LL}^A f_T^V f_T^A f_{T'}^{VA}$  for  $^{16}\text{O}$ ; (c): circles  $^{12}\text{C}$ , squares  $^{16}\text{O}$ , triangles  $^{40}\text{Ca}$   
 (d): —  $^{12}\text{C}$ , ...  $^{16}\text{O}$  and - -  $^{40}\text{Ca}$ ; horizontal bands: empirical  $\mathcal{R}$ ,  $\Delta$

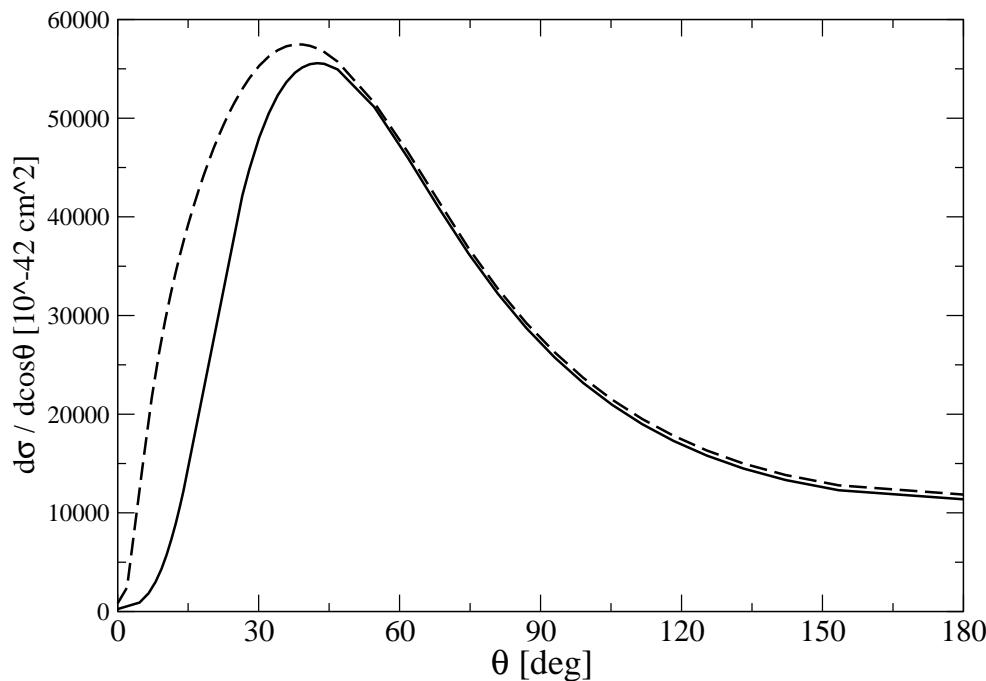


# $\frac{d\sigma}{d\omega}$ for $^{16}\text{O}$

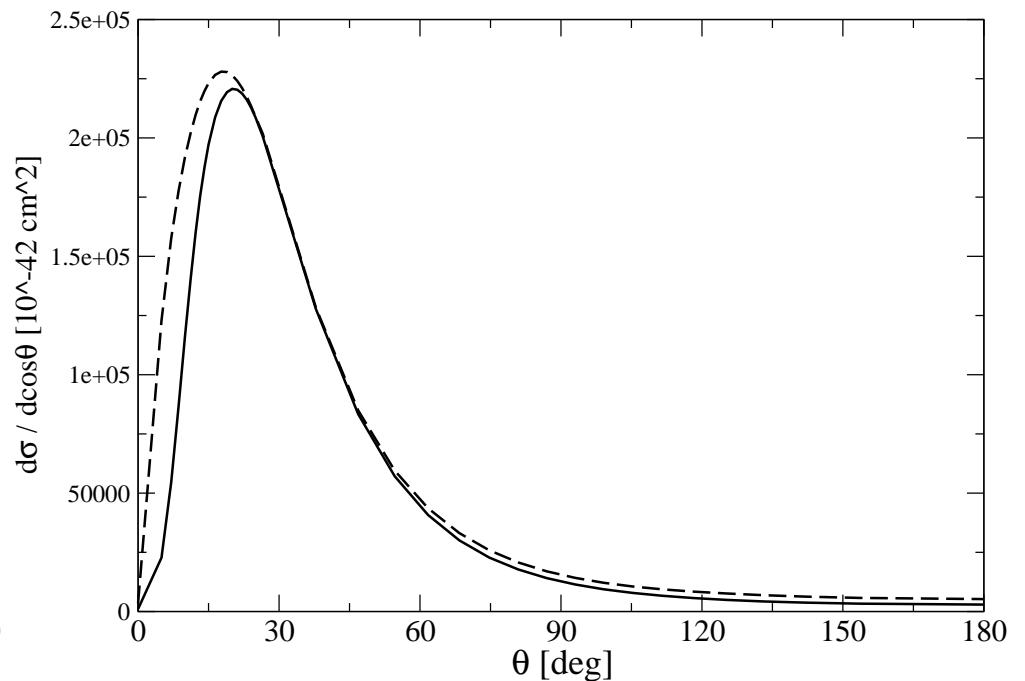


$\frac{d\sigma}{d \cos \theta}$  for  $^{16}\text{O}$

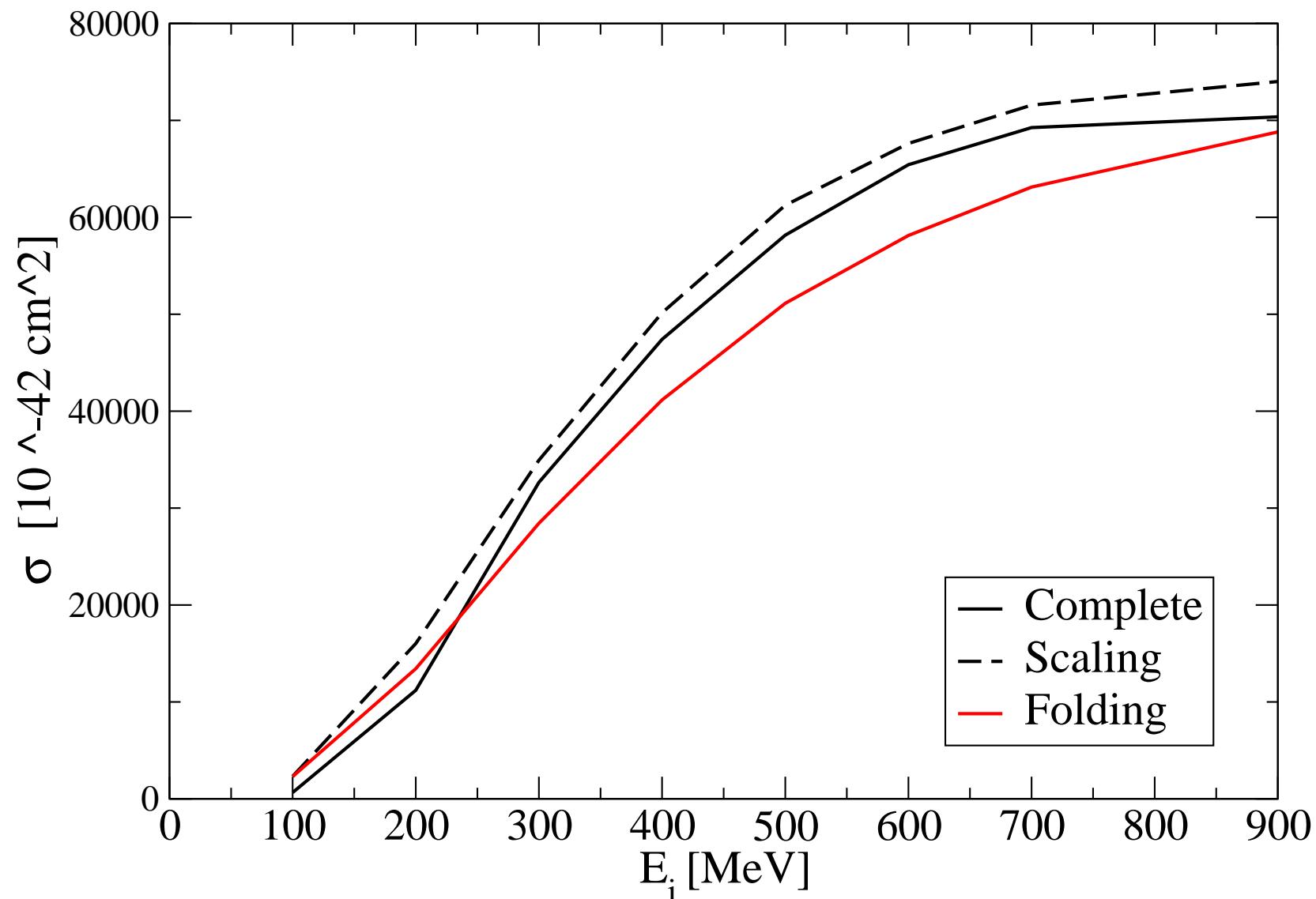
$\varepsilon_i = 500 \text{ MeV}$



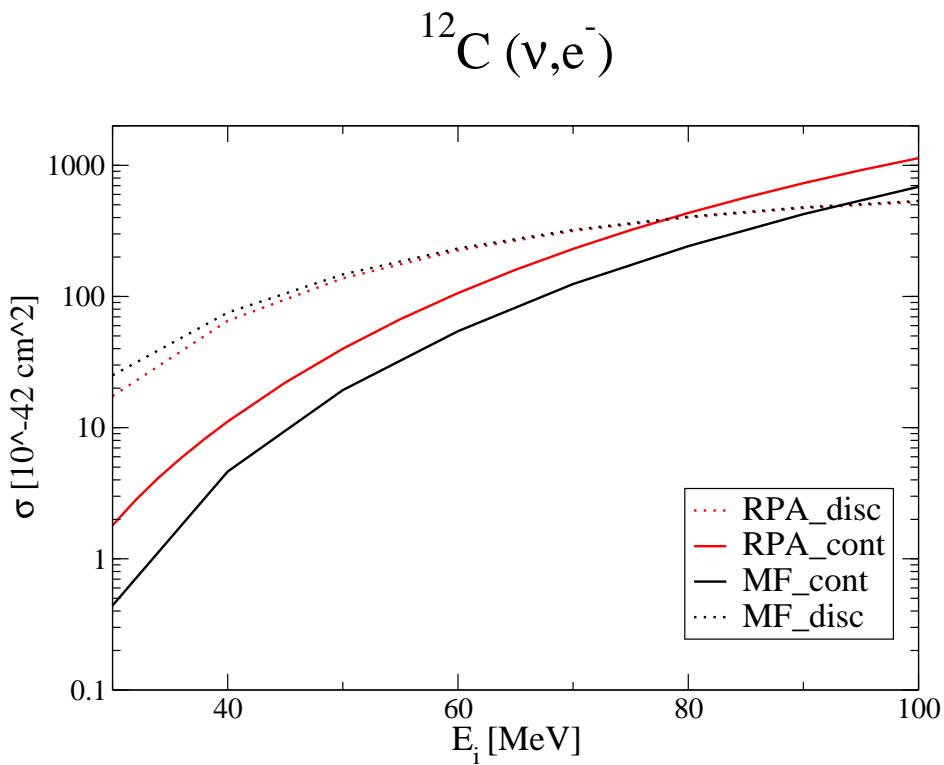
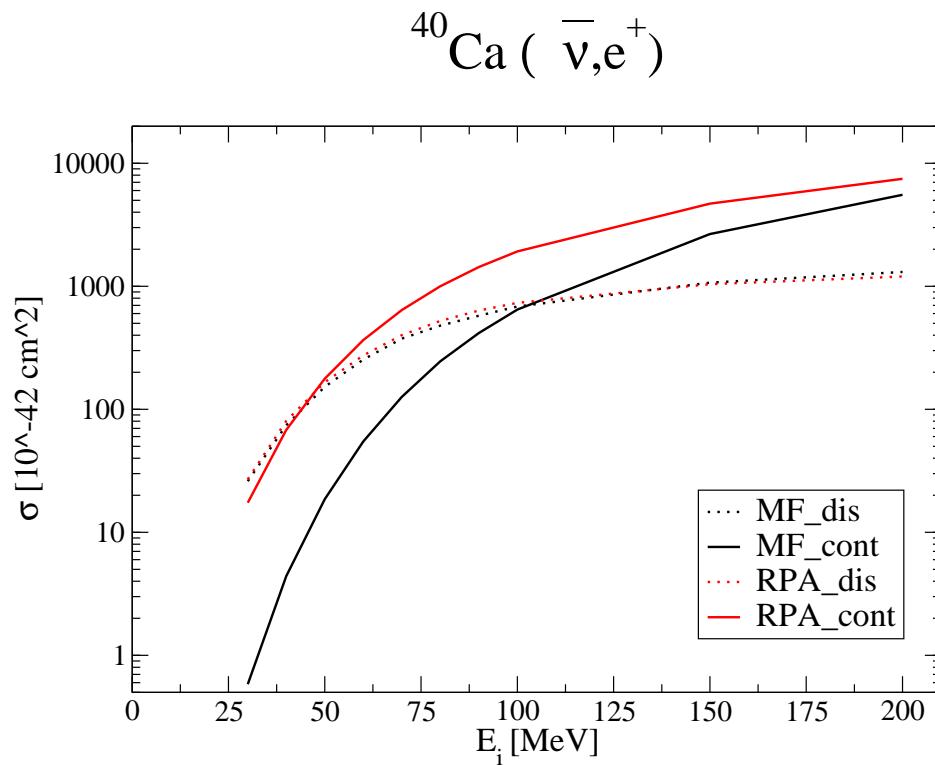
$\varepsilon_i = 1000 \text{ MeV}$



# Total Cross Section



# Warning: low energy



Discrete vs Continuum

MF vs RPA

# Summary

## \*Scaling

**SM:** 1st-kind:  $q \gtrsim 400$  MeV/c;

2nd-kind:  $q=700$  MeV/c yes,  $q=300$  MeV/c no; 0-kind: good

**RPA:** for  $q \gtrsim 500$  MeV/c effects negligible; 1st-kind:  $q \gtrsim 400$  MeV/c; worsening with zero range interaction; 0-kind: slightly ruined

**MEC:** 1st-kind:  $q \gtrsim 400$  MeV/c;

inclusion of the  $\Delta$  currents slightly decreases  $f_T$

**FSI:** responsible for the largest modifications of the SM results; do not strongly change scaling properties

## \*Electron and Neutrino Cross Sections

**theoretical scaling functions**

**empirical scaling function**

**full calculations**

# Appendix

## Semi-relativistic corrections in neutrino case

$$R_{CC}^V(q, \omega) \rightarrow \frac{q^2}{q^2 - \omega^2} R_{CC}^V(q, \omega)$$

$$R_{LL}^A(q, \omega) \rightarrow \left(1 + \frac{q^2 - \omega^2}{4m_N^2}\right) R_{LL}^A(q, \omega)$$

$$R_T^V(q, \omega) \rightarrow \frac{q^2 - \omega^2}{q^2} R_T^V(q, \omega)$$

$$R_T^A(q, \omega) \rightarrow \left(1 + \frac{q^2 - \omega^2}{4m_N^2}\right) R_T^A(q, \omega)$$

$$R_{T'}^{VA}(q, \omega) \rightarrow \sqrt{\frac{q^2 - \omega^2}{q^2}} \sqrt{1 + \frac{q^2 - \omega^2}{4m_N^2}} R_{T'}^{VA}(q, \omega)$$

## Equivalent definition of scaling functions

$$f_L(\Psi) = k_F \frac{q}{m} \frac{R_L(\omega, q)}{Z(G_E^p)^2 + N(G_E^n)^2}$$

$$f_T(\Psi) = 2 k_F \frac{m}{q} \frac{R_T(\omega, q)}{Z(G_M^p)^2 + N(G_M^n)^2}$$

$$f_{CC}^V(\Psi) = k_F \frac{q}{m} \frac{R_{CC}^V(\omega, q)}{N(G_E^{(1)})^2}$$

$$f_{LL}^A(\Psi) = k_F \frac{q}{m} \frac{R_{LL}^A(\omega, q)}{N(G_A)^2}$$

$$f_T^V(\Psi) = 2 k_F \frac{m}{q} \frac{R_T^V(\omega, q)}{N(G_M^{(1)})^2}$$

$$f_T^A(\Psi) = 1/2 k_F \frac{q}{m} \frac{R_T^A(\omega, q)}{N(G_A)^2}$$

$$f_{T'A}^{VA}(\Psi) = k_F \frac{R_{T'}^{VA}(\omega, q)}{N G_M^{(1)} G_A}$$

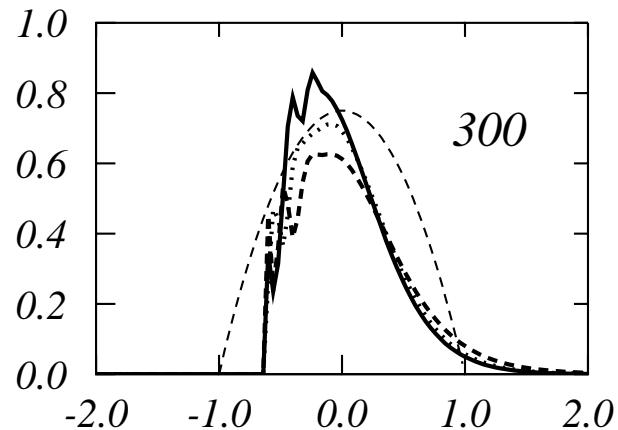
## Empirical scaling function

$$f^{\text{emp}}(\psi) = \frac{A \exp(-\psi^2) + B\psi^2 + C\psi + D}{(\psi + E)^2 + F^2}$$

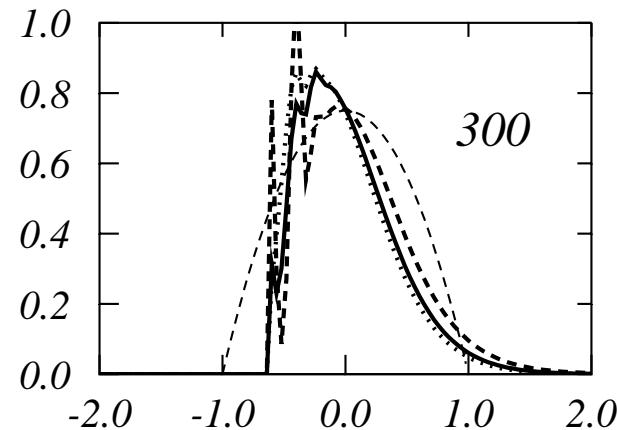
$$A=0.971, B=-0.067, C=0.385, D=0.145, E=0.366, F=1.900$$

# RPA

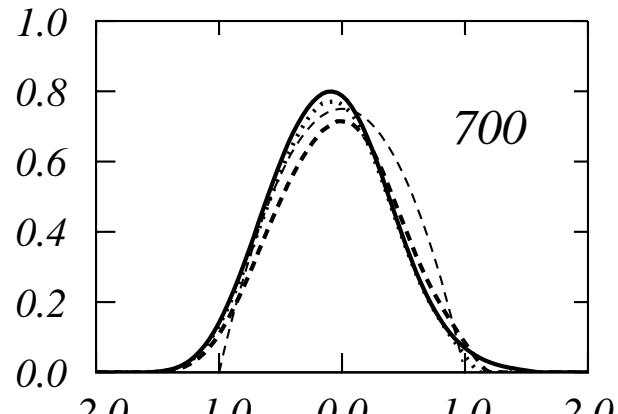
$f_L$



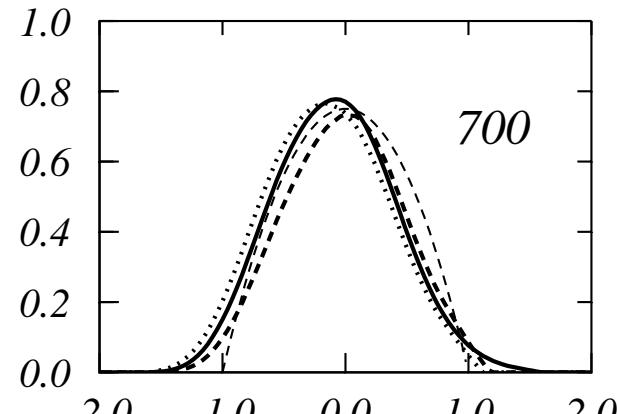
$f_T$



700

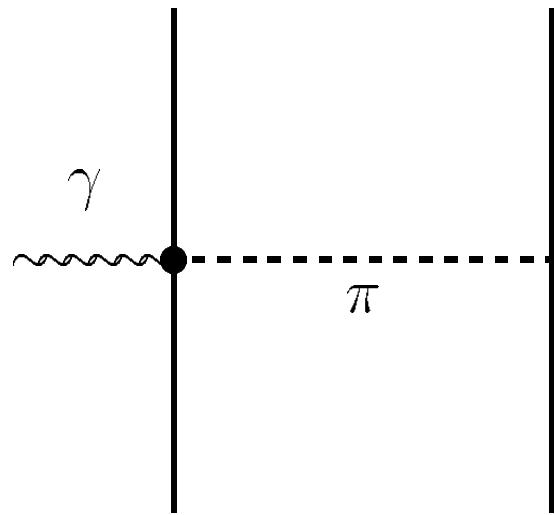


700

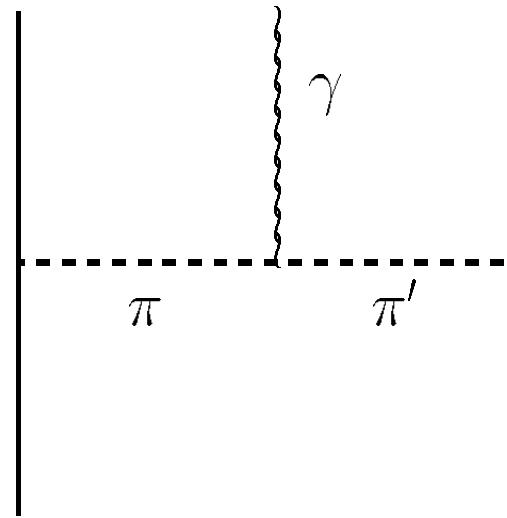


$\Psi$

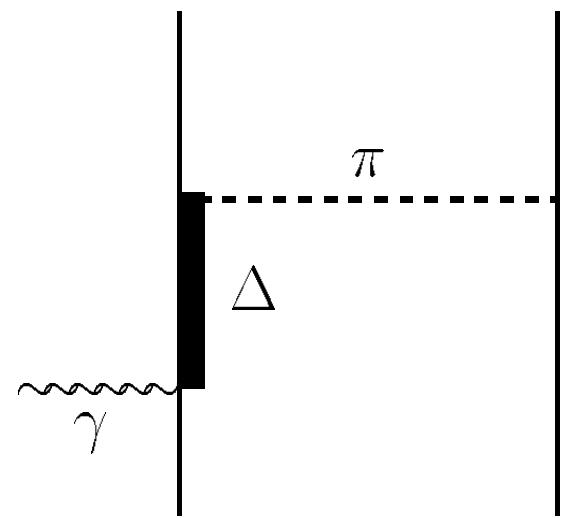
# MEC



(a)

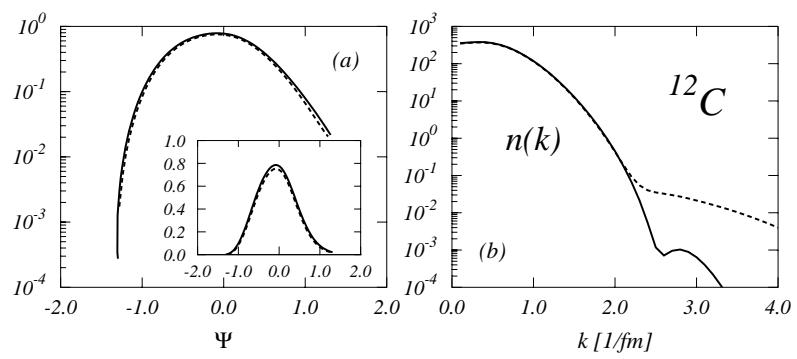


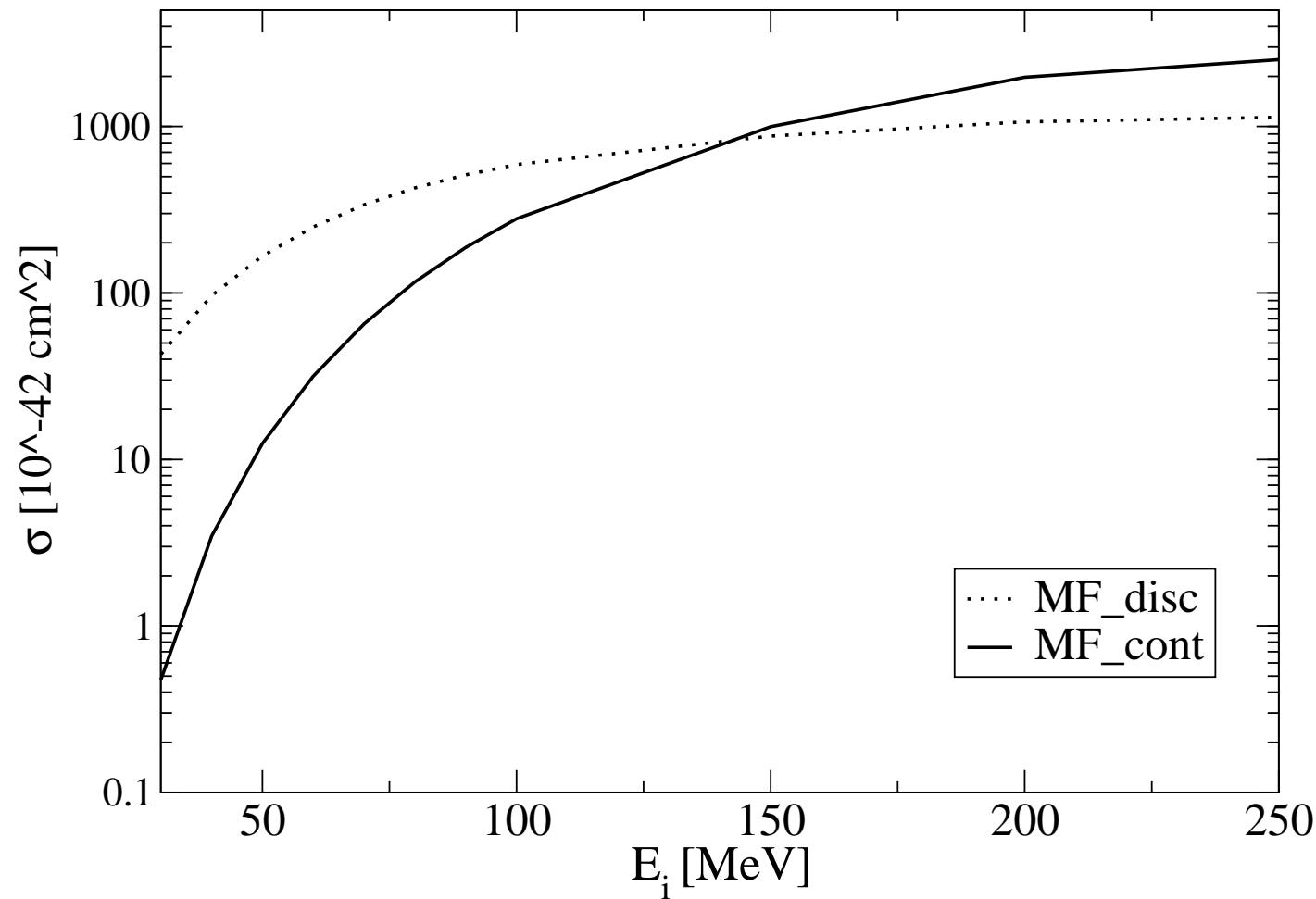
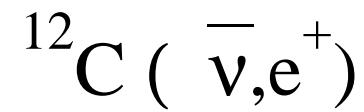
(b)



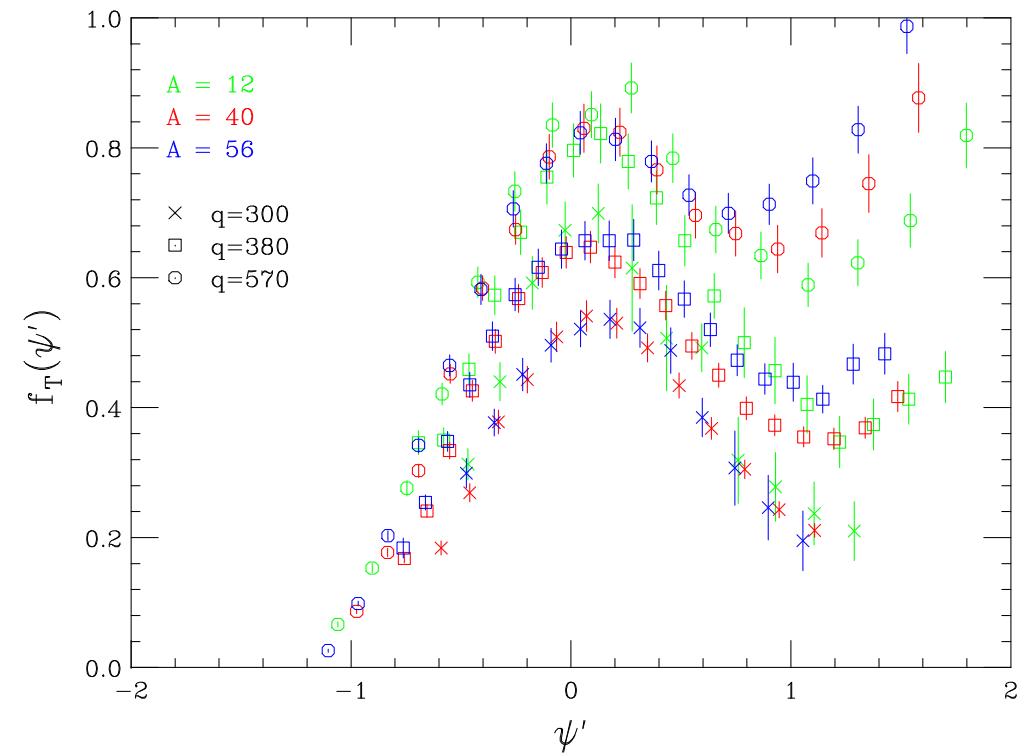
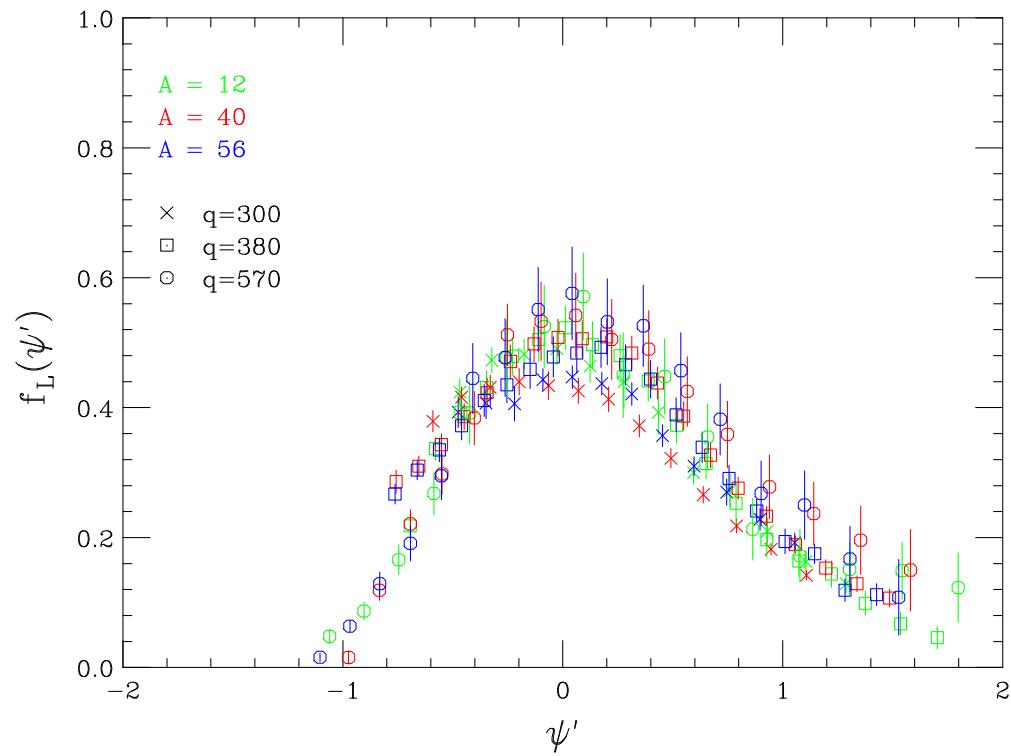
(c)

SRC



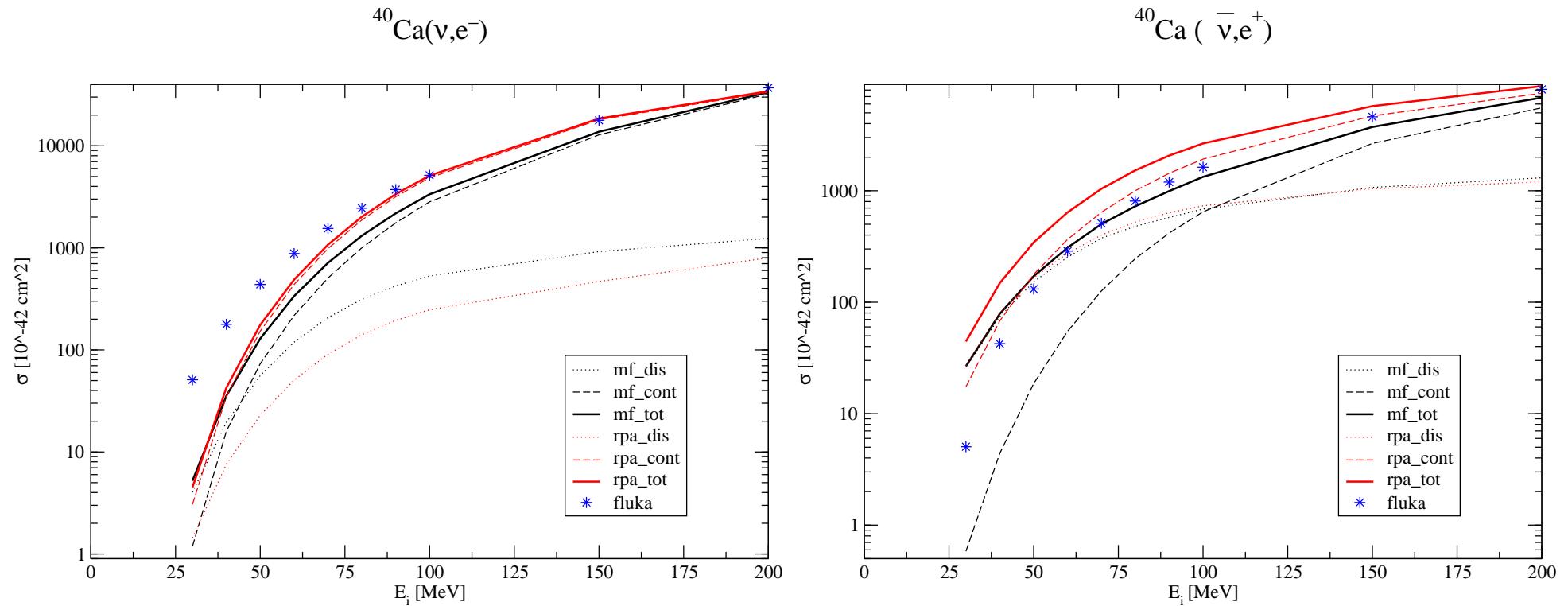


# Electron scattering



The longitudinal part superscales

# $^{40}\text{Ca}$ : MF-RPA vs FLUKA



$(\nu, e^-)$ : \* good agreement above 80 MeV between RPA and FLUKA

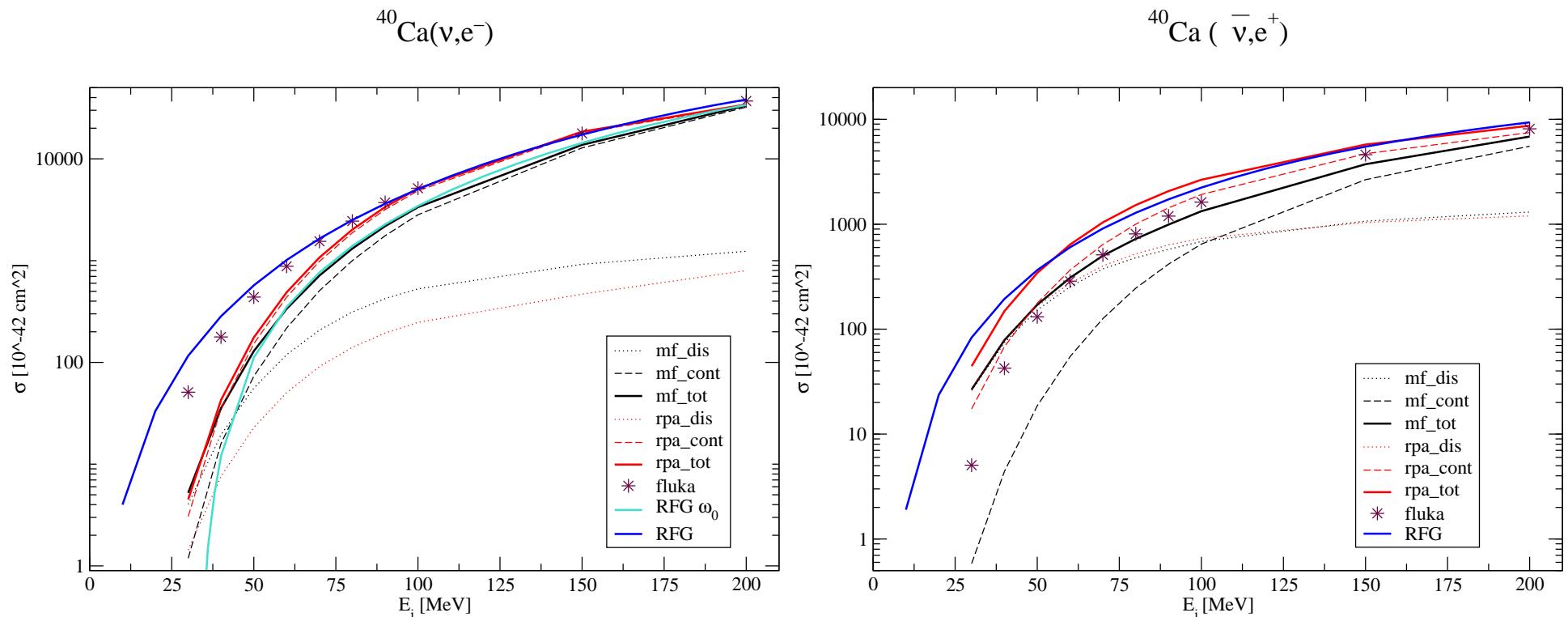
\* RPA=MF=FLUKA @  $E_i = 200$  MeV

$(\bar{\nu}, e^+)$ : \* FLUKA  $\simeq$  continuum RPA

\* relevant discrete contribution up to  $E_i = 100$  MeV

\* RPA  $\simeq$  MF in discrete case

# $^{40}\text{Ca}$ : MF-RPA vs RFG

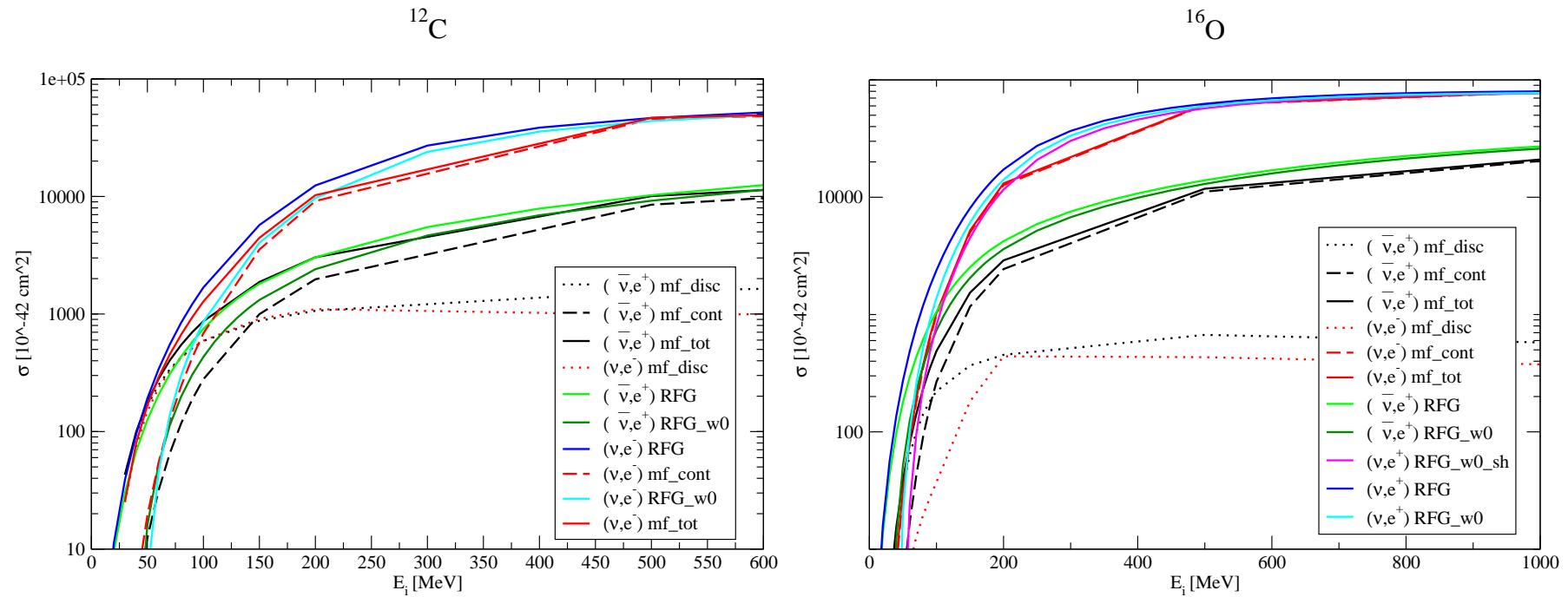


$(\nu, e^-)$ : \* good agreement above 60 MeV between RFG and FLUKA

\* MF  $\simeq$  RFG  $\omega_0$  above 50 MeV

PRC 71, 015501 (2005), Amaro et al.:  $\omega_0 \equiv \frac{1}{2M_i}(M_f^2 - M_i^2) \geq 0$ ;  $\omega = \omega_0 + \frac{|Q^2|}{2M_i}$

# $^{12}\text{C}$ and $^{16}\text{O}$ : disc. vs cont., $(\bar{\nu}, e^+)$ - $(\nu, e^-)$ , MF vs RFG



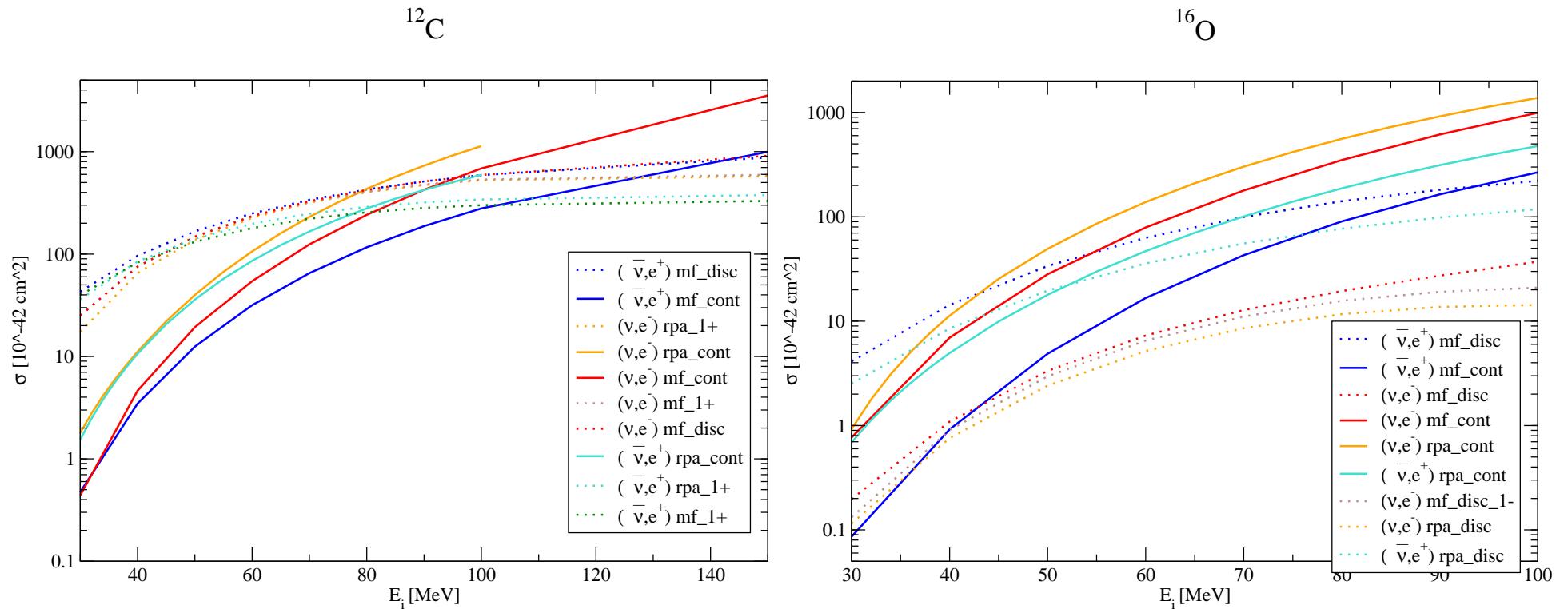
$^{12}\text{C}$ :

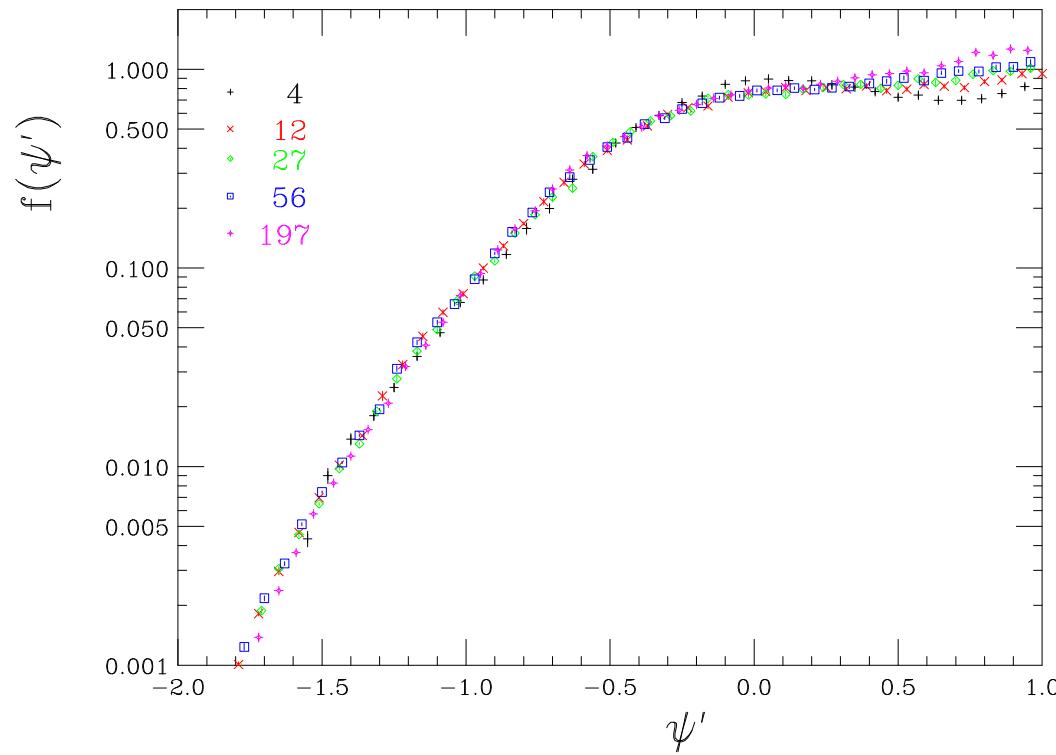
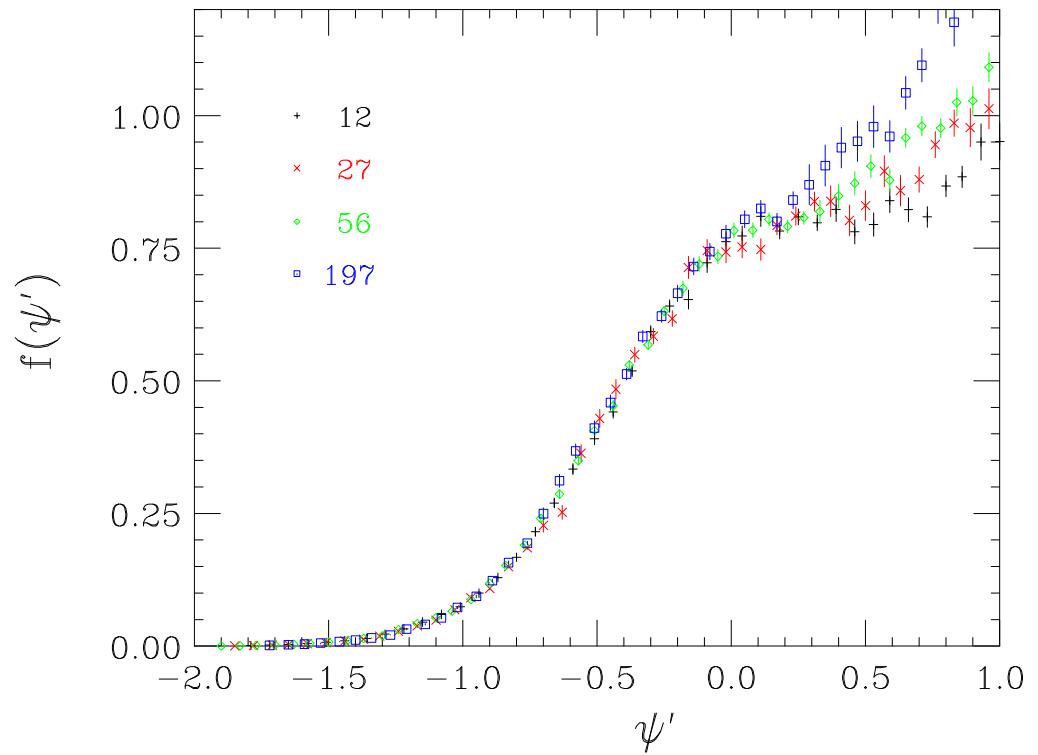
\* relevant discrete contribution up to  $E_i = 100$  or  $150$  MeV

\*  $\sigma(\nu, e^-)_{DISC} \simeq \sigma(\bar{\nu}, e^+)_{DISC}$

\*  $(\nu, e^-)_{MF_{CONT}} \simeq RFG_{\omega_0}$

# Nuclear Structure

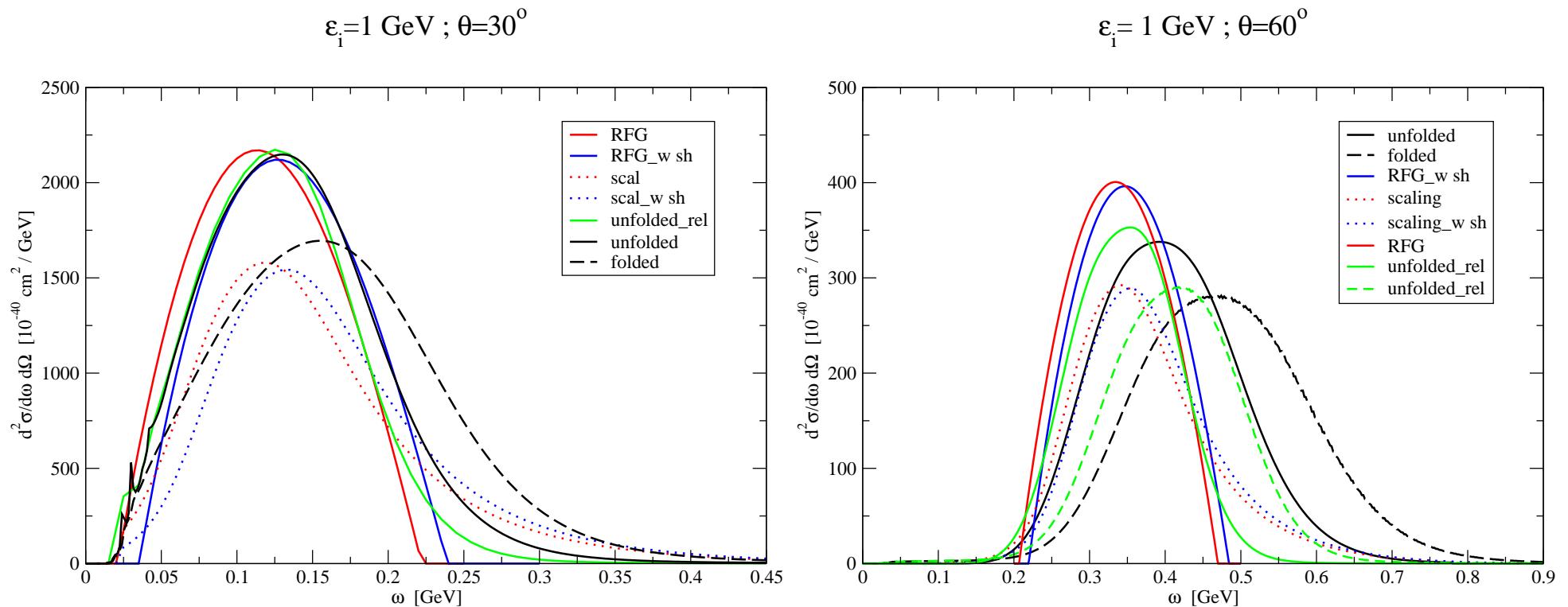




# Scaling of the second kind

# $^{16}\text{O}$ : MF vs RFG, Folding vs Scaling

## Double differential cross section



Relativistic kinematics:  
 $\epsilon_p \rightarrow \epsilon_p(1 + \epsilon_p/2m_N)$

