

Cortona, 13 Ottobre 2006.

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Electromagnetic Form Factors of the nucleon in spacelike and timelike regions

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Outline

- Motivations
- A covariant expression for the nucleon electromagnetic current: the **MANDELSTAM FORMULA**
- Nucleon Bethe-Salpeter amplitude
- Quark-photon vertex function
- Projection on the Light-Front
- Nucleon electromagnetic form factors in the spacelike and in the timelike regions: **preliminary results**
- Conclusions & perspectives

Motivations

- Experimental puzzle: the ratio $\mu\text{pGEp}/\text{GMp}$ very different from 1 for $Q^2(\text{SL}) > 1$

The investigation of nucleon electromagnetic form factor in the spacelike and in the timelike regions, within the light-front dynamics:

- open a unique possibility to study the hadronic state, both in the valence and in the non-valence sector

$$\begin{aligned} |meson\rangle &= |q\bar{q}\rangle + |q\bar{q}q\bar{q}\rangle + |q\bar{q}g\rangle + \dots \\ |baryon\rangle &= \underbrace{|qqq\rangle}_{\text{valence}} + \underbrace{|qqq q\bar{q}\rangle + |qqq g\rangle}_{\text{nonvalence}} + \dots \end{aligned}$$

the Fock expansion
is meaningful
within LF framework!!

- yields the possibility to address the vast phenomenology of hadronic resonances (**Vector Meson propagation...**) in the timelike region, and then impose strong constraints on dynamical models pointing to a microscopical description of hadrons
- allows one to obtain insights into the two-body currents associated to the quark-antiquark pair production (very important in reference frame where $q^+ > 0$)

Mandelstam Formula for the EM nucleon current:SL

The matrix element in the spacelike region for the EM current is given by

$$\langle N; \sigma', p' | j^\mu | p, \sigma; N \rangle = \bar{U}_N(p', \sigma') \left[-F_2(Q^2) \frac{p'^\mu + p^\mu}{2M_N} + (F_1(Q^2) + F_2(Q^2)) \gamma^\mu \right] U_N(p, \sigma)$$

final and initial nucleon momenta

nucleon spinors

$F_1(Q^2)$ and $F_2(Q^2)$ are the Dirac and Pauli form factors. This matrix element can be approximated *microscopically* by the **Mandelstam formula**:

$$\langle N; \sigma', p' | j^\mu | p, \sigma; N \rangle = \int \frac{d^4 k_1}{(2\pi)^4} \int \frac{d^4 k_2}{(2\pi)^4} \text{Tr} \left\{ \bar{\Phi}_N^{\sigma'}(k_1, k_2, k_3, p') \times S^{-1}(k_1) S^{-1}(k_2) I^\mu(k_1, k_2, k_3, q) \Phi_N^\sigma(k_1, k_2, k_3, p) \right\}$$

$S(k_i)$ are the Feynman propagators, I^μ is the quark-photon vertex function, and Φ_N is the nucleon **Bethe-Salpeter amplitude**.

Nucleon Bethe-Salpeter amplitude

The nucleon Bethe-Salpeter amplitude is given by:

$$\Phi_N^\sigma(k_1, k_2, k_3, p) = \Lambda(k_1, k_2, k_3) \mathcal{P}(k_1, k_2, k_3) \chi_{\tau_N} U_N(p, \sigma)$$

where Λ is a momentum dependent function and $\mathcal{P}$ accounts for the Dirac structure of the qqg-N vertex.

$\mathcal{P}$ is based on the following **effective Lagrangian**:

$$\begin{aligned} \mathcal{L}_{eff}(x, \tau_1, \tau_2, \tau_3, \tau_N) = & \epsilon_{abc} \int d^4x_1 d^4x_2 d^4x_3 \mathcal{F}(x_1, x_2, x_3, x) \times \\ & \left[m_N \alpha \frac{1}{\sqrt{2}} \sum_{\tau_1, \tau_2, \tau_3} \bar{q}^a(x_1) T_{\tau_1}^\dagger \epsilon \tau_y T_{\tau_2}^* \gamma^5 q_C^b(x_2) \bar{q}^c(x_3) T_{\tau_3}^\dagger + \right. \\ & \left. - (1 - \alpha) \frac{1}{\sqrt{6}} \sum_{\tau_1, \tau_2, \tau_3} \bar{q}^a(x_1) T_{\tau_1}^\dagger \epsilon \tau_y T_{\tau_2}^* \gamma^5 \gamma_\mu q_C^b(x_2) \cdot \bar{q}^c(x_3) T_{\tau_3}^\dagger (-\epsilon \partial^\mu) \right] \psi_N(x) T_{\tau_N} + \\ & + \text{cyclic} + \text{h.c.} \end{aligned}$$

M. Beyer et al,
nucl-th/9804021

We take $\alpha=1$

Quark-photon vertex function

We have an **isoscalar** and an **isovector** contribution:

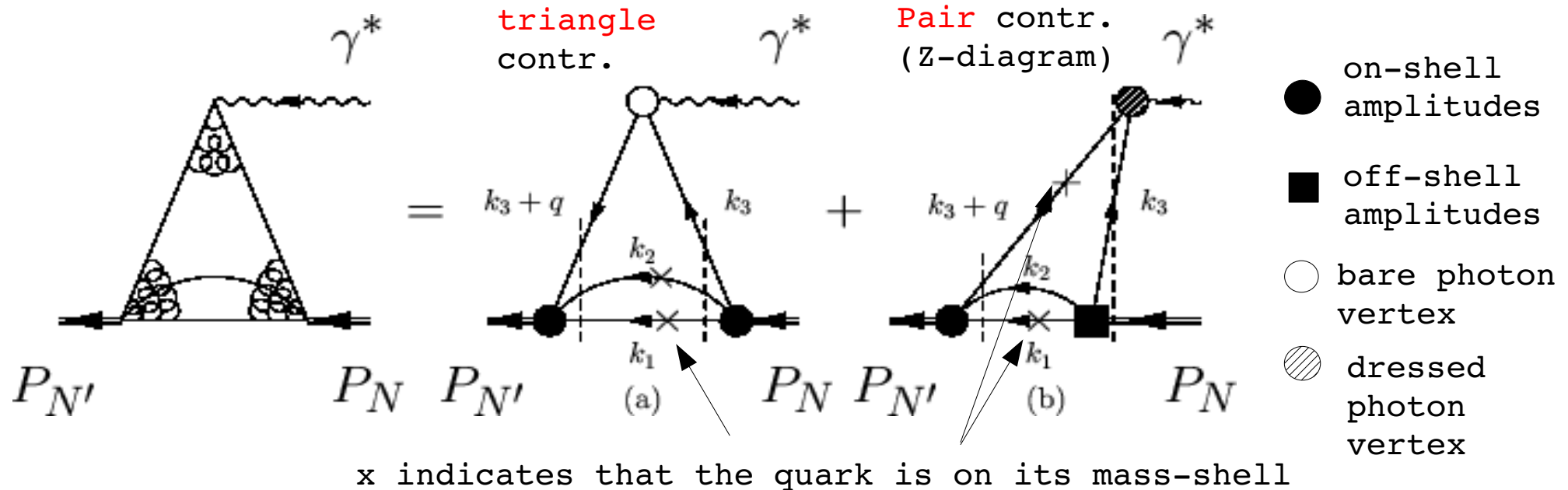
$$\begin{aligned} I_3^\mu(k_3, q) &= \left(\frac{1 + \tau_z}{2} \right) I_u^\mu(k_3, q) + \left(\frac{1 - \tau_z}{2} \right) I_d^\mu(k_3, q) = \\ &= \frac{I_u^\mu(k_3, q) + I_d^\mu(k_3, q)}{2} + \tau_z \frac{I_u^\mu(k_3, q) - I_d^\mu(k_3, q)}{2} = I_{IS}^\mu(k_3, q) + \tau_z I_{IV}^\mu(k_3, q) \end{aligned}$$

Each term contains a purely valence contribution (in the spacelike region only), i. e. an **elastic term** (**triangle**) and a **pair production term** (**Z-diagram**)

Projecting out the Mandelstam formula on the Light-Front: SL

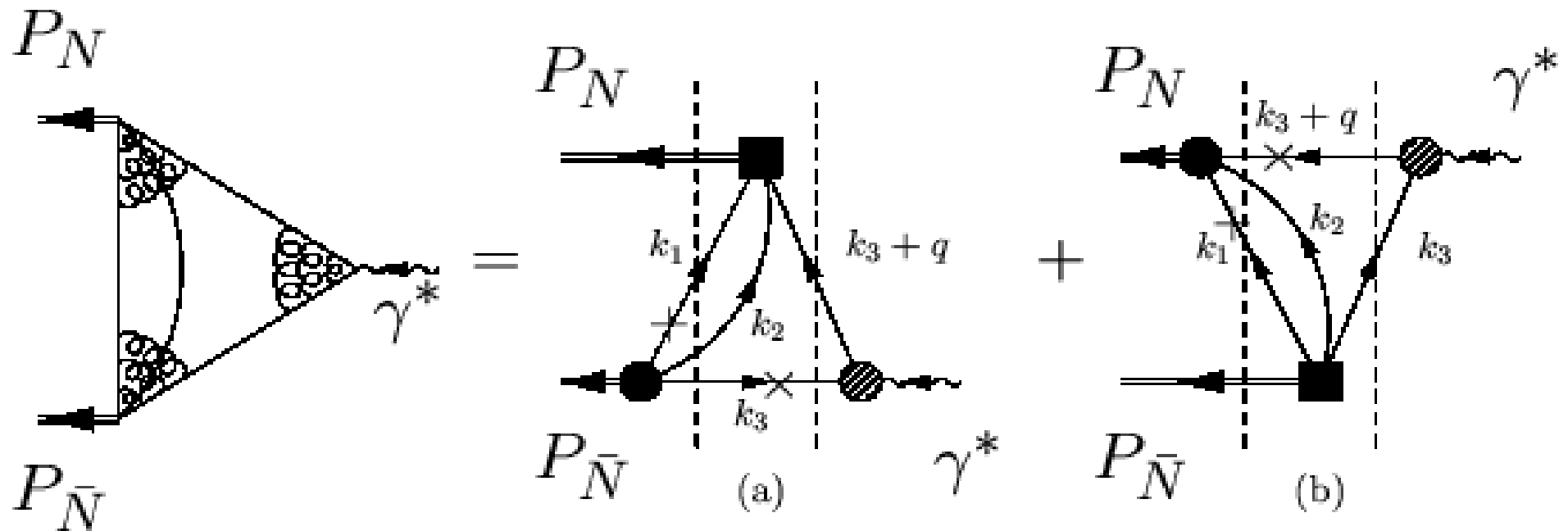
...through a k - integration, by adopting the so-called **PPA** (*Propagator Pole Approximation*), i.e. by disregarding the analytical structure form the Bethe-Salpeter amplitudes.
The reference frame we will adopt is $q^+ > 0$ and $q_{\text{perp}} = 0$.

Spacelike process:



Projecting out the Mandelstam formula on the Light-Front: TL

...and timelike process:



Phenomenological approximations: 1

We will introduce some phenomenological approximations: the momentum dependent parts of the nucleon Bethe-Salpeter amplitudes are approximated as follows:

Nucleon:

$$W_N \sim \frac{1}{[\beta^2 + M_0^2(1, 2, 3)]^3}$$

In the **valence sector** the spectator quarks are on their own k^- -shell, and **the momentum dependence**, reduced to a 3-momentum dependence through the LF projection, **is approximated through a Nucleon Wave Function a la Brodsky** (pQCD inspired)

In the **non-valence** vertex, needed to evaluate the Z-diagram contribution, **the momentum dependence is approximated by this expression:**

$$G_N \sim \frac{1}{[\beta^2 + M_0^2(1, 2)]^2} \left\{ \frac{1}{[\beta^2 + M_0^2(3', 2)]} + \frac{1}{[\beta^2 + M_0^2(3', 1)]} \right\}$$

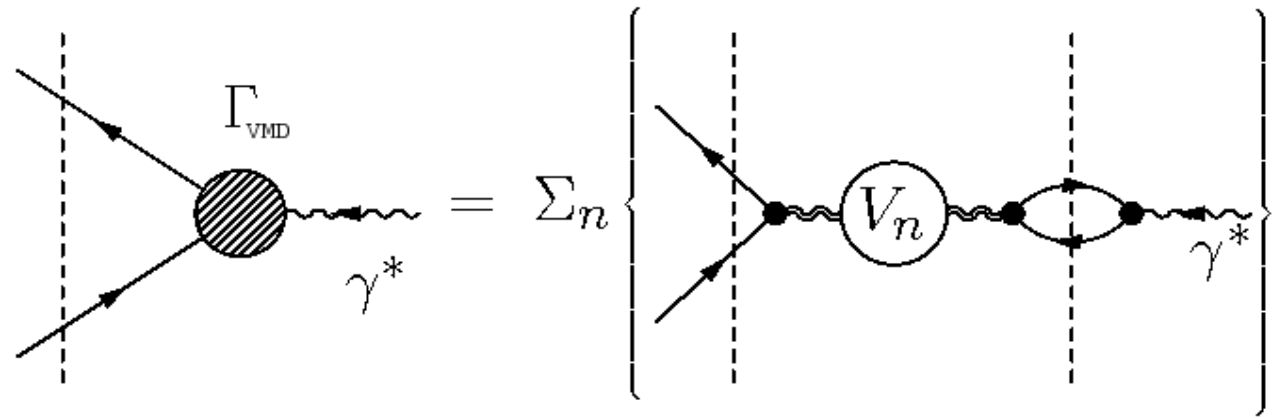
Quark-photon vertex function

It is composed by a **bare** term and by a **dressed** term:

$$\mathcal{I}_{IS(IV)}^\mu(k, q) = \frac{1}{2} \theta(p^+ - k^+) \theta(k^+) \gamma^\mu + \theta(q^+ + k^+) \theta(-k^+) \left[\frac{Z_B}{2} \gamma^\mu + Z_{VMD} \Gamma_{VMD}^\mu(k, q, IS(IV)) \right]$$

bare terms
proportional to γ^μ
dressed term

The dressed term
is described
through a
microscopical
VMD approximation:



T. Frederico et al, PRD 66 (2002) 116011

where the **vertex function** is taken as

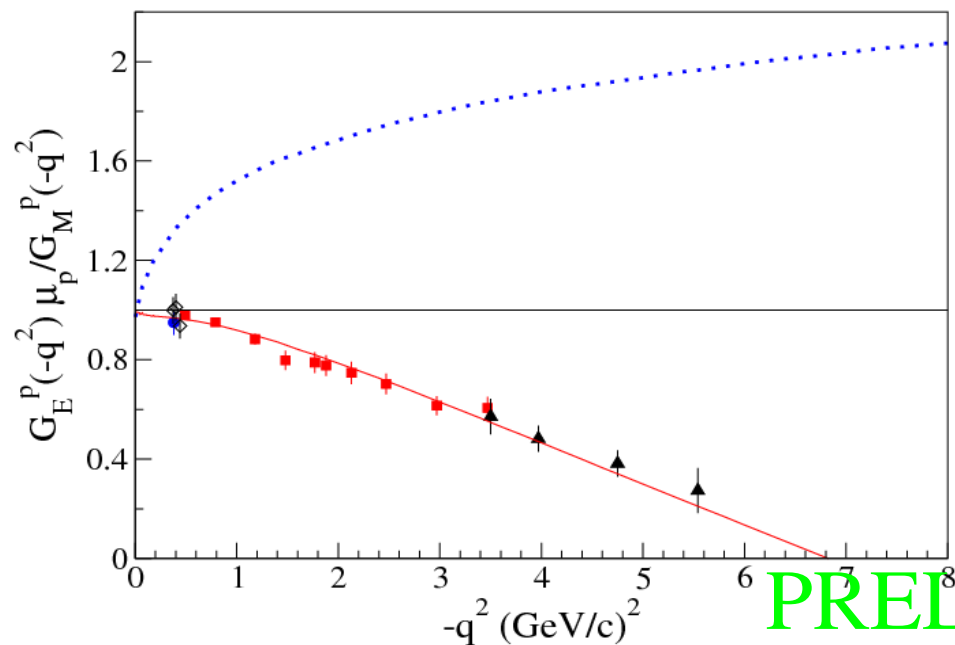
$$\hat{V}_n^\mu(k, k - P_n) = \gamma^\mu - \frac{k_{on}^\mu - (q - k)_{on}^\mu}{M_0(k^+, \mathbf{k}_\perp; q^+, \mathbf{q}_\perp) + 2m}$$

Adjusted parameters

In this model, the following parameters appear:

- **constituent quark mass** $m_u = m_d = 220 \text{ MeV}$
- the **oscillator strength** $\omega = 1.556 \text{ GeV}^2$ appearing in the VMD approximation, chosen in order to reproduce the e.m. decay widths of the first four vector mesons
- **$\beta = 0.118 \text{ GeV}$** , fixed through the anomalous magnetic moments: $\mu_p = 2.878$ (Exp. 2.793), $\mu_n = -1.859$ (Exp. -1.913)
- the **widths for the vector meson** $\Gamma_n = 0.15 \text{ GeV}$
- the **renormalization constants** for the pair production term Z_B and Z_{VMD}

Results in the spacelike region: PROTON



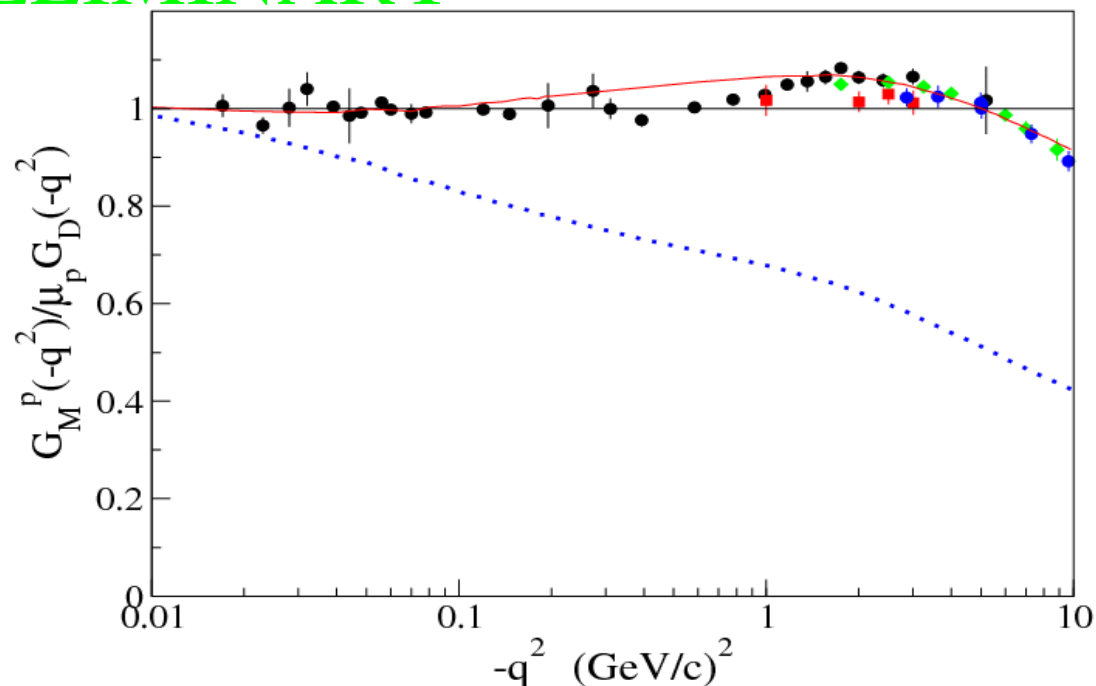
blue dotted line: triangle

red line: triangle + Z-diagram

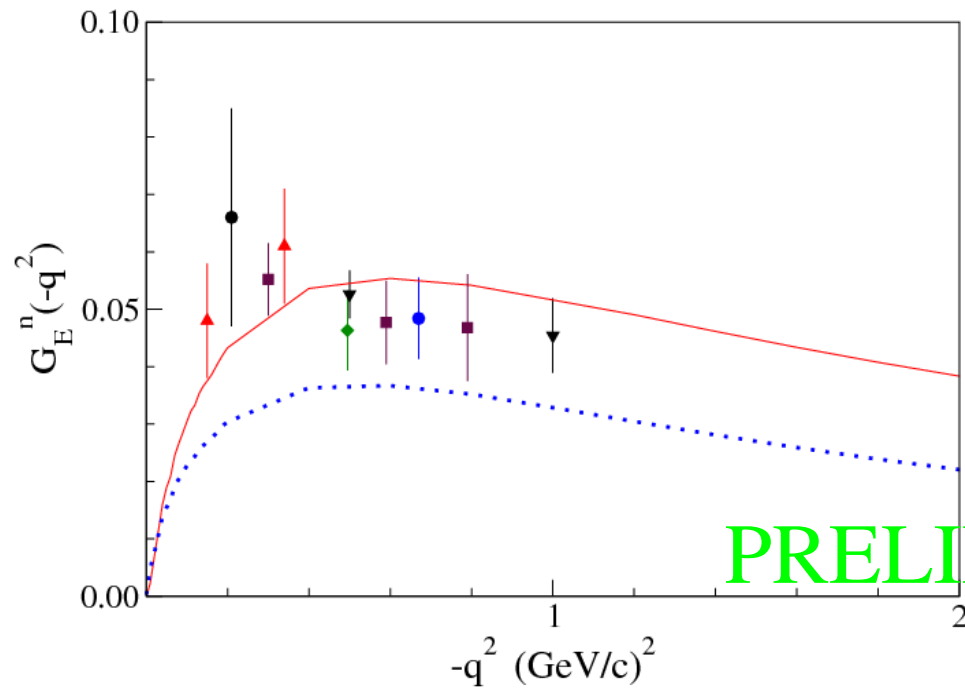
PRELIMINARY

Considerations:

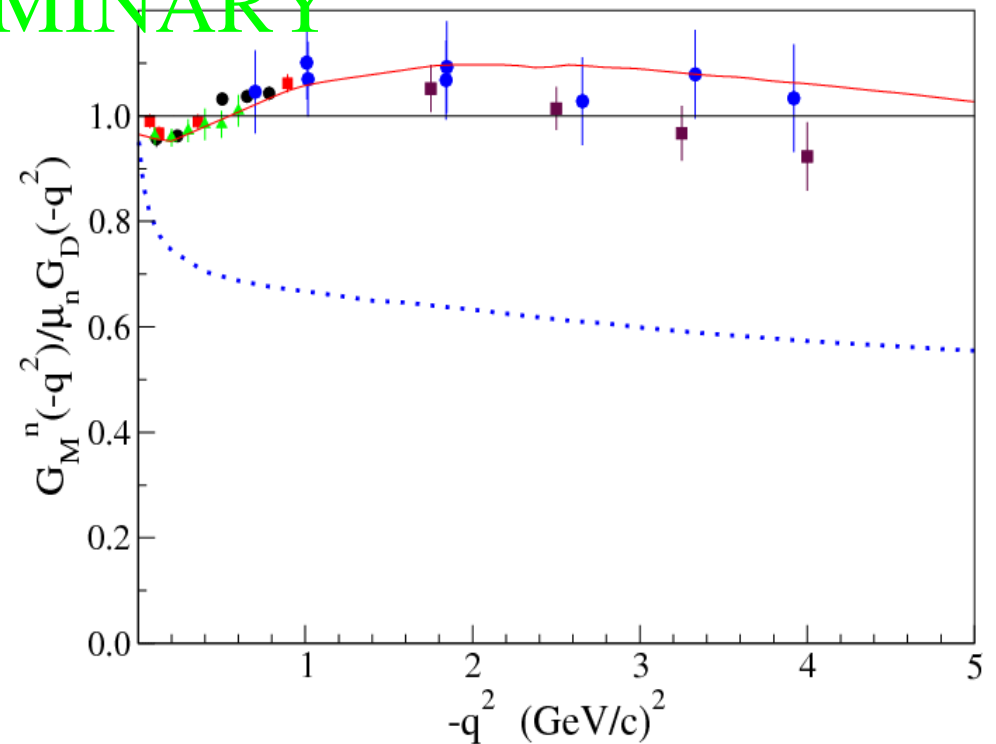
- From these results turns out to be very important the Z-diagram term
- The almost linear decrease in the ratio is well reproduced



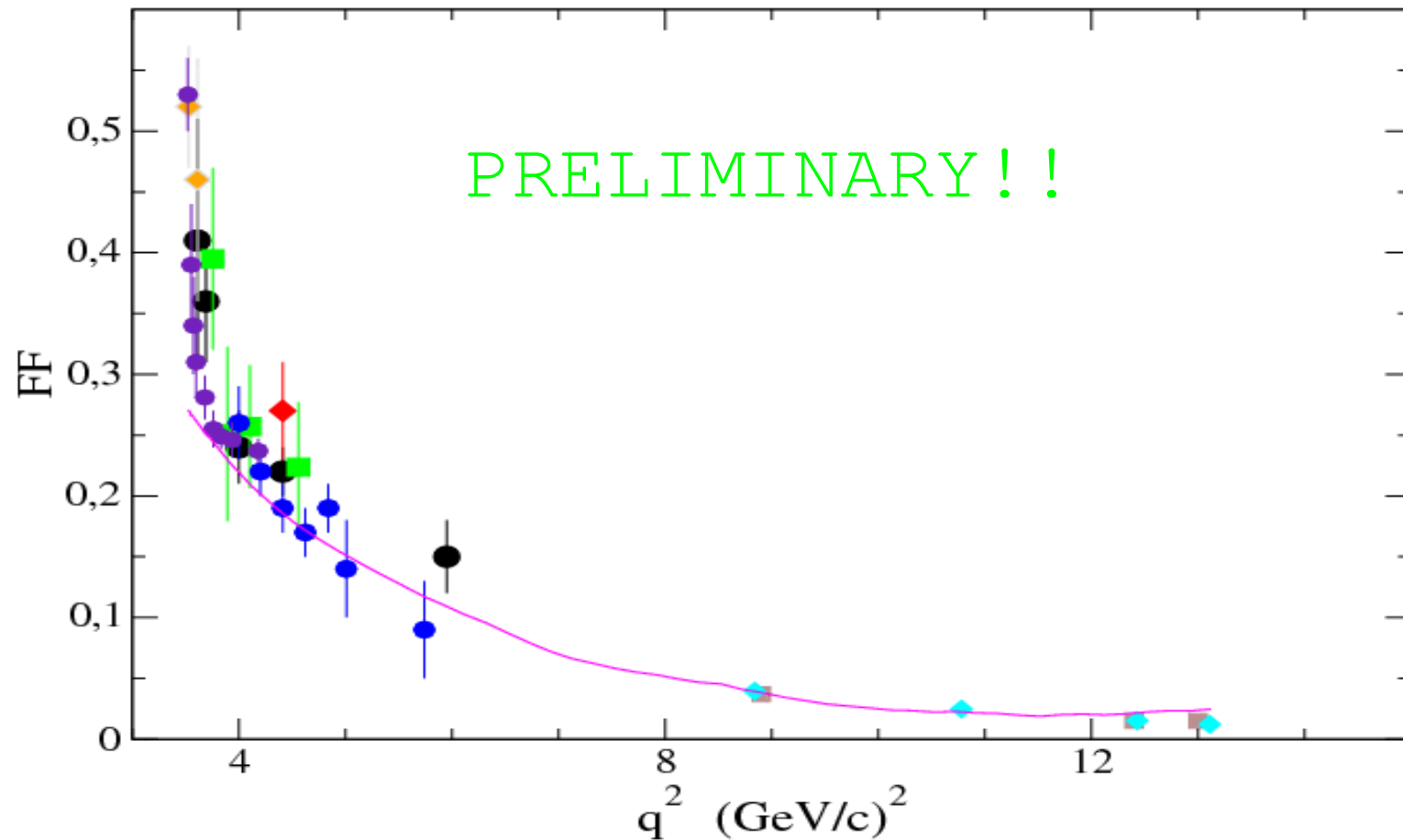
Results in the spacelike region: NEUTRON



PRELIMINARY



Proton results in the timelike region



$$FF = \sqrt{\frac{\sigma(e^+e^- \rightarrow N\bar{N})}{\sigma_{pl}}}$$

$$Z_B(TL) \sim 3 Z_B(SL)$$

Conclusions...

- A microscopical model for nucleon electromagnetic form factors in both spacelike and timelike region has been proposed
- The quark-photon vertex for the process where a virtual photon materializes in a quark-antiquark pair is approximated by a VMD model plus a bare term
- the Z-diagram are fundamental in the adopted reference frame ($q^+ > 0$)
- The preliminary results in the spacelike region are quite good, and the possible zero in the ratio $\mu p G_E p / G_M p$ seems to be strongly related to the Z-diagram contribution

... and Perspectives

- A better description of the quark-photon vertex may improve the results: a fully covariant VMD model is under investigation
- Different approximations for the nucleon wave function may be tested.