Elastic and quasi-elastic electron scattering off isotopic and isotonic chains

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- Properties of exotic nuclei through electron scattering
- Relativistic model for ground-state observables
- Elastic electron scattering
- Inclusive quasi-elastic electron scattering
- Parity violating asymmetry

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Lagrangian

$$\begin{split} \mathcal{L} &= \bar{\psi}(i\gamma^{\mu}\partial_{\mu} - M)\psi + \frac{1}{2}\partial^{\mu}\sigma\partial_{\mu}\sigma - U(\sigma) - g_{\sigma}\sigma\bar{\psi}\psi \\ &- \frac{1}{4}\Omega^{\mu\nu}\Omega_{\mu\nu} + \frac{1}{2}m_{\omega}^{2}\omega^{\mu}\omega_{\mu} - g_{\omega}\omega_{\mu}\bar{\psi}\gamma^{\mu}\psi \\ &- \frac{1}{4}\vec{R}^{\mu\nu}\vec{R}_{\mu\nu} + \frac{1}{2}m_{\rho}^{2}\vec{\rho}^{\mu}\vec{\rho}_{\mu} - g_{\rho}\vec{\rho}_{\mu}\bar{\psi}\gamma^{\mu}\vec{\tau}\psi \\ &- \frac{1}{4}F^{\mu\nu}F_{\mu\nu} - eA_{\mu}\bar{\psi}\gamma^{\mu}\frac{(\mathbb{I}-\tau_{3})}{2}\psi \end{split}$$

Nonlinear sigma potential

$$U(\sigma) = \frac{1}{2}m_{\sigma}^{2} + \frac{1}{3}g_{2}\sigma^{3} + \frac{1}{4}g_{3}\sigma^{4}$$

Field tensors

$$\begin{split} \Omega^{\mu\nu} &= \partial^{\mu}\omega^{\nu} - \partial^{\nu}\omega^{\mu} \\ \vec{R}^{\mu\nu} &= \partial^{\mu}\vec{\rho}^{\nu} - \partial^{\nu}\vec{\rho}^{\mu} - \mathbf{g}_{\rho}(\vec{\rho}^{\mu}\times\vec{\rho}^{\nu}) \\ F^{\mu\nu} &= \partial^{\mu}\mathbf{A}^{\nu} - \partial^{\nu}\mathbf{A}^{\mu} \end{split}$$



Relativistic model for ground-state observables

• Gap parameter: for open shell nuclei is employed the constant gap approximation with empirical Δ

$$\Delta^{(5)}(N) = -\frac{1}{8} \left[M(N+2) - 4M(N+1) + 6M(N) - 4M(N-1) + M(N-2) \right]$$

Five different isotopes



- One photon exchange
- Differential cross section

$$\left(rac{\mathrm{d}\sigma}{\mathrm{d}\Omega}
ight)_{EL} = \sigma_M |F_{
ho}(q)|^2$$

Charge form factor

$$F_{
ho}(q) = \int \mathrm{d}r \, j_0(qr)
ho_{
ho}(r)$$

Coulomb distortion



Inclusive differential cross section

$$\left(\frac{\mathrm{d}\sigma}{\mathrm{d}\varepsilon\,\mathrm{d}\Omega}\right)_{\mathrm{QE}} = \sigma_{\mathrm{M}}\left[\mathbf{v}_{\mathrm{L}}\mathbf{R}_{\mathrm{L}} + \mathbf{v}_{\mathrm{T}}\mathbf{R}_{\mathrm{T}}\right]$$

Coefficients

$$v_L = \left(\frac{|\mathbf{Q}^2|}{|\mathbf{q}|^2}\right)^2 \qquad v_T = \tan^2 \frac{\theta}{2} - \frac{|\mathbf{Q}^2|}{2|\mathbf{q}|^2} \qquad \mathbf{Q}^2 = |\mathbf{q}|^2 - \omega^2$$

Longitudinal and transverse response functions

 $R_L(q,\omega) = W^{00}(q,\omega)$ $R_T(q,\omega) = W^{11}(q,\omega) + W^{22}(q,\omega)$

Hadron tensor

$$W^{\mu\mu}(\boldsymbol{q},\omega) = \overline{\sum_{i}} \sum_{f} \left| \langle \Psi_{f} | \hat{J}^{\mu}(\boldsymbol{q}) | \Psi_{0}
angle \right|^{2} \delta(\boldsymbol{E}_{0} + \omega - \boldsymbol{E}_{f})$$



Equivalent expression for the hadron tensor

$$W^{\mu\mu}(\boldsymbol{q},\omega) = -rac{1}{\pi} \mathrm{Im} \langle \Psi_0 | J^{\mu\dagger}(\boldsymbol{q}) G(E_f) J^{\mu}(\boldsymbol{q}) | \Psi_0
angle$$

Final expression for the hadron tensor

$$W^{\mu\mu}(\boldsymbol{q},\omega) = \sum_{n} \left[\operatorname{Re} \ T_{n}^{\mu\mu}(\boldsymbol{E}_{f} - \varepsilon_{n}, \boldsymbol{E}_{f} - \varepsilon_{n}) \right. \\ \left. - \frac{1}{\pi} \mathcal{P} \int_{M}^{\infty} \mathrm{d}\mathcal{E} \frac{1}{\boldsymbol{E}_{f} - \varepsilon_{n} - \mathcal{E}} \operatorname{Im} \ T_{n}^{\mu\mu}(\mathcal{E}, \boldsymbol{E}_{f} - \varepsilon_{n}) \right] \\ \left. T_{n}^{\mu\mu}(\mathcal{E}, \boldsymbol{E}) = \lambda_{n} \langle \varphi_{n} | j^{\mu\dagger}(\boldsymbol{q}) \sqrt{1 - \mathcal{V}'(\boldsymbol{E})} | \tilde{\chi}_{\mathcal{E}}^{(-)}(\boldsymbol{E}) \rangle \right. \\ \left. \times \left. \langle \chi_{\mathcal{E}}^{(-)}(\boldsymbol{E}) | \sqrt{1 - \mathcal{V}'(\boldsymbol{E})} j^{\mu}(\boldsymbol{q}) | \varphi_{n} \right\rangle \right.$$



Parity-violating asymmetry - Elastic scattering

- Dirac equation $[\alpha \cdot \mathbf{p} + U_{\pm}(r)]\Psi_{\pm} = E\Psi_{\pm}$
- Total potential $U(r)_{\pm} = V(r) \pm \gamma_5 A(r)$

• Axial potential
$$A(r) = \frac{G_F}{2\sqrt{2}} \rho_W(r)$$

• Weak charge density

$$\rho_{W}(r) = \int \mathrm{d}\boldsymbol{r}' \, \boldsymbol{G}_{E}\left(|\mathbf{r}-\mathbf{r}'|\right) \times \left[-\rho_{n}(r') + (1-4\sin^{2}\Theta_{W})\rho_{p}(r')\right]$$

- Electric form factor $G_E(r) \approx \frac{\Lambda^3}{8\pi} e^{-\Lambda r}$ $\Lambda = 4.27 \text{ fm}^{-1}$
- Parity-violating asymmetry in Born approximation

$$A_{pv} = \frac{\frac{d\sigma_{+}}{d\Omega} - \frac{d\sigma_{-}}{d\Omega}}{\frac{d\sigma_{+}}{d\Omega} + \frac{d\sigma_{-}}{d\Omega}} = \frac{G_F Q^2}{4\sqrt{2} \pi \alpha} \left[4 \sin^2 \Theta_W - 1 + \frac{F_n(q)}{F_p(q)} \right]$$













Elastic differential cross sections

 $\varepsilon = 496.8 \text{ MeV}$





A. Meucci, M. Vorabbi, C. Giusti, F. D. Pacati and P. Finelli, Phys. Rev. C87, (2013) 054620.

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Inclusive QE differential RPWIA cross sections







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Inclusive QE differential RGF cross sections







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Physical Review C 85, 032501(R) (2012)

Physical Review Lett. 108, 112502 (2012)



Asymmetry averaged over the acceptance

$$\langle A_{\rho\nu} \rangle = rac{\int \mathrm{d}\theta \,\sin\theta A_{\rho\nu}(\theta) rac{d\sigma}{d\Omega} \,\epsilon(\theta)}{\int \mathrm{d}\theta \,\sin\theta \,rac{d\sigma}{d\Omega} \,\epsilon(\theta)}$$

Exp. Result:

$${\cal A}_{\rm pv}^{Pb} = 0.656 \pm 0.060 ({
m stat}) \pm 0.014 ({
m syst}) ~{
m ppm}$$







Proposal to Jefferson Lab

C-REX: Parity-violating measurement of the weak charge distribution of ⁴⁸Ca at 2.2 GeV











In preparation

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- Elastic and quasi-elastic electron scattering give information on the global properties of nuclei and on the single particle aspects of the nucleus
- Isotopic and isotonic chains
- Parity-violating asymmetry parameter investigates the neutron skin
- Good agreement with PREX measurement on ²⁰⁸Pb
- Prediction for the future experiment CREX on ⁴⁸Ca

